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THE TURBULENT BOUNDARY LAYER ON A ROUGH, POROUS PLATE: EXPERIMENTAL HEAT TRANSFER WITH UNIFORM BLOWING

By

J. W. Healzer, R. J. Moilut and W. M. Kays

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Thermosciences Division

Department of Mechanical Engineering

Stanford University

Stanford, California

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ABSTRACT

Stanton number measurements have been made for a transpired turbulent boundary layer on a rough surface. Tests were conducted at uniform blowing fractions from 0 to .008, with uniform surface temperature and uniform free-stream velocity. The x-Reynolds number range of these tests was from 10 to 10 and the roughness Reynolds number range from 20 to 200. The data are believed to be accurate to within ±.0001 Stanton number units over most of this range. At each test condition, several velocity profiles were taken to measure the boundary layer growth. The boundary layer momentum thickness variation along the test surface has been used to estimate rough-plate skin friction.

The data indicate the expected increase in both skin friction and heat transfer due to roughness. The data display some unusual features when plotted against boundary layer size. For any given value of F, tor all free-stream velocities, the Stanton number data fall on a single curve when plotted against enthalpy thinkness, and the skin friction data collapse (but less well) to a single curve when plotted against momentum thickness. This behavior might have been expected for a "fully rough" plate but was not expected at the lower end of the roughness Reynolds number range of these tests. The blowing data indicate that within the uncertainty of the experiment, a form of the Couette flow model used for smooth surface boundary layers can also be used for rough surfaces to predict the effects of blowing on Stanton number.

Predictions of the experimental boundary layers have been carried out using a finite-difference boundary layer prediction method which employs the Patanker-Spalding finite-difference formulation and a mean field closure with a mixing-length model employing van Driest type damping. The effects of roughness have been incorporated into this prediction program by modification of the mixing-length model used for smooth-surface turbulent boundary layers.

The apparatus constructed for these tests is a closed-loop wind tunnel using air at essentially ambient condition as both the transpired and free-stream fluid. The rough surface consists of 24 porous plates

forming an eight foot long test section. The individual plates were fabricated by brazing together 50-mil OFHC copper balls stacked in their most dense array. In the test section assembly each plate is individually controlled with its own electric heater, transpiration air supply and instrumentation. The construction of the test section is such that blowing, no blowing, and suction cases can be tested.

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NOMENCLATURE

- A van Driest damping function for a smooth surface.
- A_p van Driest damping function for a rough surface.
- B Blowing parameter, F/St.
- Cu Defined by Equation (4-9).
- c Specific heat of fluid, Btu/lb_*F.
- c_f Skin friction coefficient, $\tau_o/(\rho u_m^2/2)$
- d Pipe diameter, ft.
- E Parameter defined by the equation $y^+ = \frac{1}{\kappa} l_n(Ey^+)$
- F Blowing fraction, $\rho_o v_o / \rho_\omega u_\infty$.
- g Proportionality factor in Newton's Second Law, 1b ft/1b sec2.
- G Free-stream mass velocity = $\rho_{os} u_{os}$, $1b_{m}/sec$ ft².
- k Roughness particle size, ft.
- k Equivalent sand grain roughness size, ft.
- & mixing length, ft.
- Δl_{λ} Mixing length at wall, ft.
- m" Mass flux through the plate surface, lbm/sec ft2.
- n Exponent in equation (1-1).
- r Ball radius, ft.
- T Temperature, *F.
- u Boundary layer free-etream velocity, ft/sec.
- u_{τ} Shear velocity = $\sqrt{\tau_{o}g_{c}/\rho_{o}}$, ft/sec.
- v Velocity normal to the test surface, ft/sec.
- v Average velocity in a pipe flow, ft/sec.
- x Distance measured downstream from test section inlet, ft.
- Distance measured downstream from test section inlet to virtual origin, ft.
- y Distance normal to the surface, ft.

Nomenclature (cont.)

Greek

- a Thermal diffusivity, ft²/sec.
- V Kinematic viscosity, ft²/sec.
- $au_{_{
 m O}}$ Wall shear stress, 1b $_{_{
 m f}}$ /ft 2 .
- δ Boundary layer displacement thickness, ft.
- δ_{qq} 99% thickness of the momentum boundary layer, ft.
- λ Proportionality constant used in outer region mixing-length. formulation $(\ell = \lambda \delta_{00})$.
- κ Mixing-length constant, $\ell = \kappa y$.
- θ Boundary layer momentum thickness, ft.
- Δ Boundary layer enthalpy thickness, ft.
- ρ Density, 1b_m/ft³.
- μ Viscosity, lb_m/sec ft.
- ω Humidity, $1b_m$ water vapor/ $1b_m$ dry air

Dimensionless Groups

- Red Pipe Reynolds number = dv/v.
- Re. Roughness size Reynolds number = ku_m/v.
- Re_{τ} Roughness Reynolds number = $k_g u_{\tau}/v$.
- Re_{x} x-Reynolds number = xu_{∞}/v .
- Re_A Momentum thickness Reynolds number = $\theta u_{\infty}/v$.
- Re_{Δ} Enthalpy thickness Reynolds number = $\Delta u_{\infty}/v$.
- St Stanton number = h/Gc
- ΔSt Stanton number error, see Equation (2-2).
- St. Stanton number without blowing = h/Gc.
- Pr Prandtl number = v/α .
- Pr turbulent Prandtl number.

Nomenclature (cont.)

Subscripts

- ∞ Refers to free-stream conditions.
- w Refers to value on fluid side of wall fluid interface.
- T Rafers to condition of transpiration flow beneath porous plate.
- H Thermal (used as mixing-length subscript).
- m Momentum (used as mixing-length subscript).
- s Smooth surface value.

Superscripts

+ Refers to variables non-dimensionalized in wall coordinates; $y^+ = yu_{\tau}/v$.

CHAPTER I

INTRODUCTION AND BACKGROUND

Heat transfer between a surface and a fluid stream is affected by the surface condition, the fluid properties, the velocity field, and the thermal boundary conditions. Studies of this complex interaction are generally conducted by fixing all but one of these descriptors and varying the remaining one to determine its effect. Most prior studies of boundary layer heat transfer have been concerned with the effects of fluid properties, temperature boundary conditions, or main stream velocity distribution. Current interest in protecting surfaces by transpiration or ablation has introduced a class of problems in which surface roughness is an unavoidable feature.

A. General Background

Much of what we know about rough surface hydrodynamics is based on ideas and results of the pipe flow experiments by Nikuradse [9]. To avoid the difficulties of defining an arbitrary rough surface with an uncertain distribution of roughness elements, he investigated the flow through sand-grain roughened tubes. Roughness elements for these experiments consisted of selected sand, carefully sieved and attached in maximum density to the tube walls. The roughness in these experiments was described by a single parameter, $\mathbf{k_s}$, the size of the sand-grain elements. This 'sand-grain' measure of roughness has become a standard in skin friction studies. It is still common practice to express the effect of an arbitrary roughness in terms of an 'equivalent sand-grain roughness', $\mathbf{k_s}$.

This introductory paragraph was first used by Professor Moffat in the original proposal to the Navy for the rough surface research contract that supported these tests. This same introduction was repeated again in each of the 12 quarterly reports describing our progress on the project. It seems appropriate that it be used again to introduce the final reporting of this phase of the roughness work.

Numbers in brackets refer to references listed at the end of this report.

The pipe flow experiments showed three domains of behavior in terms of the roughness Reynolds number, $Re_{\tau} = u_{\tau}k_{s}/v$. For values of roughness Reynolds number less than 5 the flow behaved as though the surface were smooth. For values of roughness Reynolds number greater than 70, the pipe friction factor became independent of Reynolds number, $Re_{d} = vd/v$. This state was described as 'fully rough'. Values of roughness Reynolds number between 5 and 70 defined a region of 'transitonal roughness'.

Prandtl [10] and von Karman [11] used Nikuradse's results (in 1934) to predict the behavior of the rough surface boundary layer. These predictions indicated that the boundary layer would attain a 'fully rough' state such that the skin friction would be a function only of x/kg. Early rough surface boundary layer experimental studies were carried out by Moore [12] and Hamma [13]. Moore looked at air flowing over a flat plate roughened by regular arrays of square bars attached to its surface. Hamma's experiment was carried out with air flowing over surfaces roughened by screens attached to them. Several other experimenters looked at the rough surface boundary layer, but most notable among these were two studies by Perry et al. [14,15] and Liu et al. [16]. The experiments by Perry and his co-workers were carried out with lateral rectangular bars on a flat plate and with zero and adverse pressure gradients. They found that pressure gradient did not alter the roughness effects on the boundary layer. They also attempted to simulate a change in effective sandgrain roughness by changing the spacing between the array of bars used to roughen the plate. They found that changing the bar spacing did not change the effective sand-grain roughness of the surface as had been expected. This led them to define a second type of roughness whose behavior, unlike sand-grain roughness, could not be correlated in terms of an external length scale, but scaled on the 'logarithmic asymptote'. This is the distance between the top of the roughness elements and the apparent surface of the plate, the plane where the velocity profiles extrapolate to zero.

The experiment by Liu was conducted in water using an array of rectangular bars to roughen the otherwise smooth surface. Liu also varied his bar spacing to change the roughness of his surface and determined equivalent sand-grain roughnesses for all surface roughness that he tested. He was not able, however, to correlate his skin friction data from different surface roughnesses with a single x-Reynolds number expression. Liu's report also contains a very complete table of both the experimental and analytical work that had been done till then (1966) for both rough pipe and boundary layer type flows.

More recent rough surface hydrodynamic studies have been made by Grass [17], who used a hydrogen bubble technique to measure instantaneous velocity distributions in a water tunnel above a sand-grain roughness. His study documented details of the turbulent structure near the wall. Wu [18] used a floating element balance to measure skin friction in an air tunnel with a sand-grain roughened surface. These tests showed 'fully rough' behavior and agreed well with the PrandtI-Schlichting prediction. Tsuji and Iida [19] examined velocity profiles over rough surfaces and showed that they could be predicted using a modified mixinglength approach, maintaining a non-zero value of the mixing length at the wall. In a related study, Antonia and Luxton [20] investigated the effect of an abrupt change from a smooth to a rough surface. In this study the rough surface was constructed from parallel square bars like many earlier studies, but one of the bars was instrumented with static pressure taps to provide a direct measure of form drag. The method of measuring form drag on individual roughness elements was also used by Perry [14]. Townes et al. [21,22] have studied the structure of turbulent flow in sand-grain roughened pipes.

B. Previous Heat-Transfer Experimental Work

Much less has been done in the field of heat transfer. One of the first systematic rough surface experimental studies was carried out by Nunner [23]. These experiments used air flowing through rough pipes. The results were used to establish a simple empirical relationship between the increase in Nusselt number due to roughness and the increase in the skin friction. Several important heat transfer studies followed, notably by Dipprey and Sabersky [24], Owen and Thompson [25] and Gowen and Smith [26].

Dipprey and Sabersky looked at the flow of four fluids of different Prandtl numbers through one smooth pipe and three rough pipes. These studies demonstrated that rough wall heat transfer varied with the Prandtl number, even in the fully rough regime where molecular viscosity effects seem not important.

Owen and Thompson proposed a flow model near the rough surface which explains why increasing surface roughness can increase the heat transfer. Based on their ideas, they developed an expression for a sublayer Stanton number which was able to correlate a large block of the rough pipe heat transfer data.

Gowen and Smith studied flow in several rough pipes using fluids of three different Prandtl numbers. In one of their pipes they were able to make temperature profile measurements in the flow over the rough wall. Like earlier investigators, they developed an expression for rough wall Stanton number based on the rough wall skin friction, but in this case they were able to include information obtained from their temperature profiles.

There were, of course, other rough wall heat transfer studies, mostly confined to internal flow applications. Two papers which summarize this work are by Sood and Johnson [29] and more recently by Norris [30]. The survey by Norris points out that of the many forms of rough wall heat transfer correlations that have been proposed, none is simpler or has had a great deal more success in correlating the data than the simple expression by Nunner,

$$\frac{Nu}{Nu_{g}} = \left(\frac{f}{f_{g}}\right)^{n} , \qquad (1-1)$$

where the exponent n can be expressed as a function of Prandtl number. The general form of this relationship is shown in Fig. 1.1, and Norris's recommended Prandtl variation is shown in Fig. 1.2. An approximate expression for the exponent is:

$$n = .68(Pr)^{0.215}$$
 (1-2)

An interesting feature of all the pipe experiments is that there is a 'ceiling' which corresponds to a rough-to-smooth skin friction ratio of about 4.0. This ceiling seems to vary with the type of roughness, with sand-grain roughness producing lower heat transfer than surfaces roughened with wires and square ribs. This range in the heat transfer ceiling is shaded in Fig. 1.1. For increases in skin friction above this ceiling, there is not a corresponding increase in the heat transfer.

There continues to be current interest in rough wall heat transfer. Much of it is specifically related to heat transfer to nose cones and other re-entry vehicles. Unfortunately, much of the work done is not in the open literature. Two related studies by Reshotko [27] and Boldman and Graham [28] examine heat transfer to a nozzle with a rough surface.

C. Rough-Surface Heat-Transfer Models

Several authors have presented prediction methods for rough wall boundary layer heat transfer. Integral boundary layer prediction schemes which include roughness effects have been described by Dvorak [32,33] and by Chen [31]. Nestler [34] has proposed a scheme using a correlation by Owen and Thompson to relate the increase in Stanton number due to roughness to the increase in the skin friction and other boundary layer parameters. Finite-difference turbulent boundary layer prediction schemes which include roughness effects have been described by Lumsdaine et al. [35] and by McDonald and Fish [36]. Each of these studies has modified the mixing-length distribution to introduce roughness effects. The Lumsdaine paper discusses a prediction program for skin friction only. It employs a van Driest type damped mixing-length model near the wall. To account for roughness, they follow the recommendation of van Driest [39] and add an additional term to the mixing-length damping expression which includes a roughness parameter.

The McDonald and Fish method includes both skin friction and heat transfer calculations. Roughness effects are entered through the mixing-length damping by defining an incremental damping due to roughness which is added to the normal boundary layer damping. This damping effect is the sole means of including roughness effects into both the hydrodynamic and heat-transfer predictions in this method.

D. Objectives of the Present Research

The first objective of this study was to find a suitable rough surface to test. We sought a surface which was at the same time repeatable, describable, and also porous.

The second objective was to measure the heat transfer from the rough surface to a turbulent boundary layer over a wide range of free-stream velocities. As a secondary part of this objective, the overall hydrodynamic performance of the rough surface was to be determined in terms of boundary layer growth and skin friction.

The third objective was to study the effect of blowing on the rough surface heat transfer to determine whether or not transpiration changed the effective roughness.

A final objective was adopted after the apparatus had been completed and testing had begun: to study the effects of roughness and blowing on heat transfer in the transition region. At the lowest test section velocity, it was discovered that a fully developed, laminar boundary layer existed over the first quarter of the test section. It was decided that instead of tripping the boundary layer at the test section inlet to immediately establish a turbulent layer, that we would examine the effects of blowing and test section velocity in the transition region as these data would be of considerable interest and could be obtained at minimal cost.

E. General Features of the Progr. m

Spherica) roughness elements were selected for this study. The test surface was constructed from 0.050 inch diameter 0.F.H.C. copper balls arranged such that the surface formed a regular array of hemispherical caps. This surface was more regular than sand grains, being geometrically describable, but more like sand-grain roughness than the transverse-bar roughness elements used in other tests.

High-conductivity copper balls were used in fabrication of the surface to ensure a high thermal conductivity of the porous plates. With a high-conductivity surface, heaters can be embedded into the plate and yield a uniform surface temperature, independent of whether the transpiration is "on" or "off". This allows us to use the same surface to test

both the unblown as well as the blown boundary layers. This is a necessary requirement if we are to determine if blowing changes the effects of surface roughness.

The apparatus constructed to test this rough surface has the capability of operating over a wide range of free-stream conditions selected to optimize the usefulness of the data. Fig. 1.3 shows an operating map of the rough surface heat transfer apparatus. The coordinates used for this map describe the performances in terms of a roughness particle-size Reynolds number and an x-Reynolds number. The range of test section velocities of the rig was selected such that at a given particle-size Reynolds number, the x-Reynolds number could be varied by a factor of 10^3 . In these coordinates, it's possible to overlap operating conditions with the existing smooth surface boundary layer apparatus. The rig operating parameters were also restricted to ensure an essentially constant property boundary layer to minimize the effects of variable fluid properties.

The following chapters describe the heat transfer apparatus and its qualification tests, the data which have been obtained in these tests, and, finally, the attempts to predict the data with a finite-difference boundary layer prediction program.

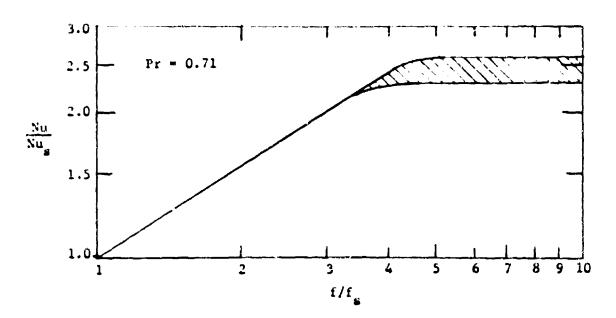


Fig. 1.1 Heat transfer increase ratio, rough-to-smooth Nusselt number versus rough-to-smooth friction factor

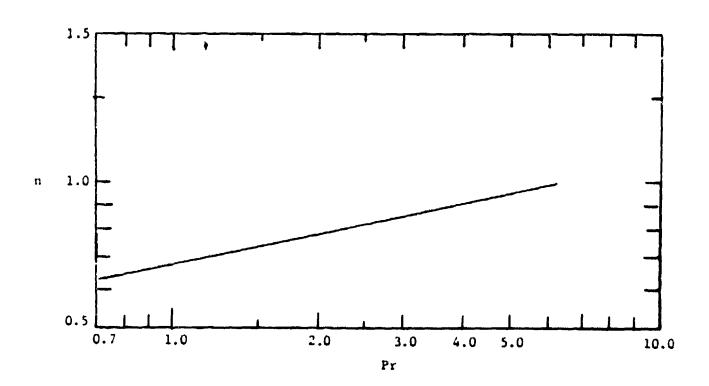


Fig. 1.2 Friction factor ratio exponent

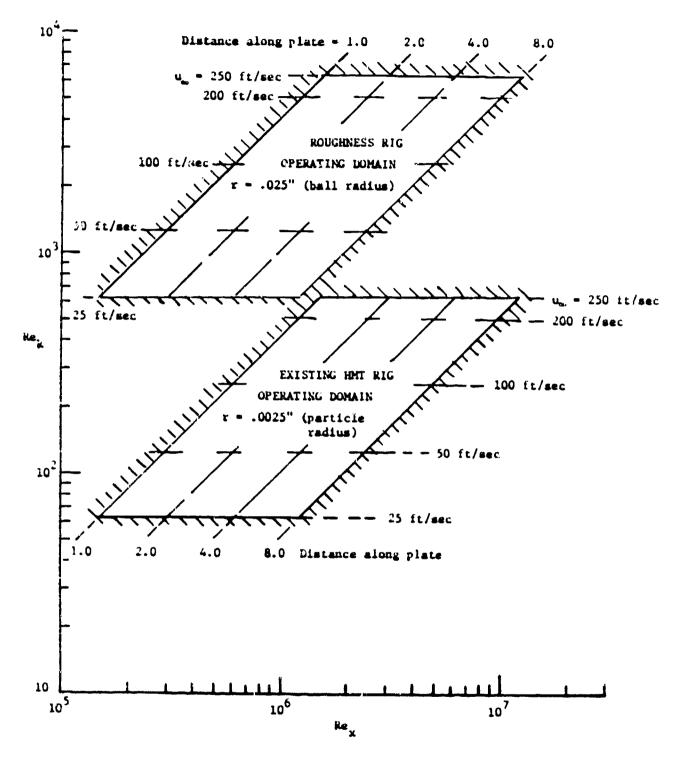


Fig. 1.3 Roughness Rig operating map, x-Reynolds number versus roughness size Reynolds number

CHAPTER II

THE EXPERIMENTAL APPARATUS

The apparatus used in these experiments was constructed especially for this study and in the discussion that follows, it will be referred to as the Roughness Rig. The Roughness Rig is located in the Thermosciences Laboratory on the second floor of the Mechanical Engineering Buil ing at Stanford. Its basic design was copied after an existing heat transfer facility which the Heat and Mass Transfer Group at Stanford has used over the past six years to investigate the transpired turbulent boundary layer on a smooth surface. This heat and mass transfer apparatus was described first by Moffat [1]. References[2 - 8] describe modifications that have been made since Moffat's original experiments. This existing HMT Rig is still very much in use and continues to play an important role in the Heat and Mass Transfer Group's research activities.

A. Description of the Apparatus

The Roughness Rig is a closed loop wind tunnel using air at essentially ambient conditions. Its test surface consists of a 24-segment porous plate, 18 inches wide and 8 feet long. Figure 2.1 shows a flow schematic of the four main rig systems: main air supply, transpiration air supply, plate heater electrical power system, and the heat exchanger cooling water system. A photograph of the Roughness Rig is shown in Figure 2.2. The following is a description of the four main rig systems.

A.1 The Main Air Supply System

The flow path of the main air system is as follows: (1) main air blower, (2) overhead ducting to an oblique header, (3) main-stream heat exchanger, (4) screen box, (5) nozzle to test section inlet, (6) 8 foot long test section and (7) a multistage diffuser which returns the main-stream air back to the blower.

The Roughness Rig main air supply blower is a 445-BL Class 3 Buffalo Blower with a 20 horsepower belt drive. At rated conditions, this blower delivers 8300 cfm and develops a head of 12 inches of water. The blower and drive are mounted on a 900 lb seismic base to minimize vibration. Flexible boots connect the blower to the remainder of the tunnel. The main air stream velocity in the test section is varied by changing the pulleys and belts on the blower and drive. The range of blower speeds attainable with this drive arrangement is from the blower rated speed of 2400 rpm down to 372 rpm.

From the main blower discharge, air is delivered through a 24" diameter overhead duct to an oblique inlet header on the main-stream heat exchanger. Ducting construction is from galvanized sheet metal with gasketed joints and sealed internally at all seams with a silastic gasket sealer material. The oblique inlet header to the heat exchanger was designed based on recommendations by Wolf [40]. The header shape is specified for uniform flow distribution and minimum pressure loss.

The header supplies air to a 5 row, 33" x 48" heat exchanger used for main air stream temperature control. The heat exchanger cooling water is continuously pumped from a holding tank through both the main-stream heat exchanger and the transpiration air heat exchanger. In operation, the cooling water temperature is adjusted until the main-stream air at the test section inlet is at smoient temperature. This avoids the formation of a thermal boundary layer prior to the test section inlet.

The heat exchanger is followed by a screen box containing four stainless steel, #40 mesh, 0.0065" dia. wire screens. The screens were sized to minimize the mean field velocity disturbances as well as reduce the main air stream turbulence level. Based on the work of Schubauer et al. [41], this screen pack should reduce mean velocity disturbances by a factor of 600 and turbulent fluctuations by a factor of 10. The screen pack combined with the large area contraction nozzle provides a uniform and well damped inlet velocity field to the test section. Details of the inlet velocity field measurements are reported in Section F.

The nozzle accelerates the flow from the screen box to the test section inlet, a 19.8 to 1 area contraction. The nozzle wall design is based on a polynominal shape as recommended by Rouse and Hassan [42]. This conforms with experience in designing similar nozzles at Stanford. The wall shape, which is a two-dimensional contraction, has been chosen such that both the first and second derivatives of flow area are zero at the nozzle exit. There is a slight acceleration in the inlet flow to avoid possible separation in this region. This was accomplished by designing the nozzle length for 40" and cutting off the first inch when the nozzle was actually built. A teledeltos model of each of the nozzle walls was used to insure that the nozzle smoothly accelerated the flow and there was not separation at inlet or exit.

The test section consists of the test plate assembly, two side walls and a movable upper surface. At the inlet, the test section is 20 inches wide and 4 inches high. The upper surface is pivoted at the inlet so that it can be adjusted to give either an increasing or decreasing flow area in the flow direction. The top and side walls are constructed from continuous sheets of 1/2 inch thick plexiglass. The top is sufficiently flexible so that it can be warped to accomodate the nonlinear growth of the test section boundary layers. The side walls contain two sets of static pressure taps, on 2 and 12 inch centers in the flow direction. The pressure taps on 12 inch centers in the flow direction are used to position the top to achieve the desired free stream velocity distribution. The tests described here were conducted with uniform free stream velocity. To obtain this test condition, the top wall is set, experimentally, for each run by adjusting it until there is no measurable change in static pressure along the test section centerline. In practice, local deviations of +.001 inches of water were accepted. A probe sled, which spans the test section, locks onto the side walls in fixed positions over the center of each of the 24 individual test plates. Probes, supported from this sled extend down through access holes in the movable top. In addition to the access holes along the test section centerline the second, eleventh and twenty-third (of twenty-four) of the individual test plates are provided with a full set of access holes extending across the width of the test section on one inch centers. These are primarily to examine boundary layer uniformity in the transverse direction.

Flow from the test section exits into a multistage vaned diffuser. The diffuser inlet has a movable top to allow alignment with the test section top. The inlet section is followed by three separate vaned two-dimensional diffusion sections finally emptying into a plenum box. The diffuser area ratio between test section and plenum box is in excess of 7:1. The plenum box is connected to the main blower inlet through a flexible boot. Each of the individual diffuser sections were designed based on the experimental work of Cocharan and Kline [43] and the total diffuser relovers approximately 40% of the kinetic energy head in the flow leaving the test section. A small charging blower is attached to the plenum box to make up air leakage from the tunnel. This blower is used to ensure that the test section static pressure is equal to the ambient pressure in the laboratory. This avoids any leakage of air into or out of the test section which might disturb the two-dimensionality of the flow.

Appendix A contains a brief description of the construction details of the screen box, nozzle and diffuser.

A.2 The Transpiration Air Supply System

The flow path of the transpiration air system is as follows: (1) inlet air filter box, (2) transpiration blower, (3) transpiration heac exchanger, (4) transpiration flow header box, (5) individual delivery tubes to each of the porous plate sections.

The purpose of the inlet filter box is to prevent clogging and contamination of the porous plates. It is constructed using 5 micron retention filter felt material and is sized to have 60 square feet of filter area.

The filter box is connected to a Buffalo type V, size 25 blower. This blower is driven by a 15 horsepower, 3600 RPM motor, connected directly to the blower. In this configuration, the blower runs at one speed and transpiration flow is controlled by ball valves in the individ-

ual delivery lines. This particular blower was designed for a higher volumetric flow than is required at low transpiration rates. In order to avoid off-design operation, the blower was equipped with dump valves so that a part of its discharge air could be dumped back into the lab. It was found, however, that the blower performance was stable over the entire flow range of interest in these experiments without venting its discharge.

Transpiration air from the blower is delivered through a 10 inch flexible duct to a large box containing the transpiration air heat exchanger and bypass system which allows partial or total heat exchanger by-pass. A 5 row, 18 x 24 inch heat exchanger is used in the transpiration system. The heat exchanger receives its cooling water from the same recirculating water system which supplies the main air heat exchanger. From the heat exchanger enclosure, the transpiration air is passed through several mixing turns, then dumped into a header box. In operation, the heat exchanger and its by-pass system are used to insure that the transpiration air is delivered to the header at ambient temperature. This avoids the formation of temperature gradients in the delivery system which acts itself like a heat exchanger between the transpiration air and the laboratory.

The transpiration header box is physically located under the test section. Individual supply lines, one for each porous plate, are connected to the side of the header box. The first section of each supply line is a 3 foot long riser reaching from the header box to a ball valve used to control the flow to each plate. In each riser section, upstream of the valve, is a hot-wire type flowmeter. Details concerning the fabrication and calibration of these flowmeters are described in Appendix B. The hot-wire flowmeter was chosen because of its wide sensitivity range. The range of flows which must be accurately metered to cover the operating range of the Roughness Rig is from less than 1 cfm to over 50 cfm. To cover this range with a loss-based metering system such as an orifice or rotameter would require two or more parallel meters in each line. To complete the supply line from the control valve to the test plate assembly, 1 inch flexible tubing was used. This final link in the supply line makes the Roughness Rig transpiration system compatible

with the existing HMT Rig plate assemblies. It would be possible to install the existing smooth plate assemblies on the Roughness Rig and they would be completely compatible with the transpiration system. To minimize the interaction between the transpiration air and the surroundings, the header box and supply lines have been insulated. The header box is covered with a layer of aluminum foil backed, rock wool insulation and the supply lines have been encased in a cardboard and masonite zone box to insure a thermally uniform environment for the system. Thermocouples in the individual supply lines indicate temperature differences on the order of one degree Fahrenheit from end to end of the 8 foot header box, the hotter end being nearest the transpiration blower.

When a blown boundary layer is established in the test section, the transpiration system operates in an open loop mode. Transpiration air is drawn in through the filter box and delivered to the test section. It is then dumped, downstream at the diffuser plenum box. The small charging blower which is used to control test section pressure is disconnected and air is bled from the diffuser plenum box through a control slide valve. Test section static pressure can easily be balanced to match the ambient pressure by controlling the flow of air through the control valve.

When boundary layer suction is required the flexible duct connection is moved to the suction side of the blower and the blower discharge is used to charge the diffuser plenum to provide test section static pressure control.

A.3 The Plate Header Electrical Power System

The plate heater power supply is a 750 amp, 24 kilowatt Lincoln Arc Welder. The field resistance of the welder has been modified to fix its output at 22 volts D.C. Power connections are made to the ground and 'high point' taps on the welder to minimize line voltage droop with increasing load. Power is delivered to a bus bar box mounted on the side of the Roughness Rig through overhead copper bus bar system. Each individual plate has its own heater which consists of a single piece of #26 AWG stranded copper wire with irradiated PVC insulation. This wire

is laced back and forth in eight grooves in the back of the plate. Ends of the heater wire are connected to #12 AWG stranded copper leads which pass through the side of the aluminum support casting. One heater lead is connected to a precision ammeter shunt, one for each plate, and then to the ground bus bar in the box. The other lead is connected to a power transistor mounted on the other bus bar in the box. The power transistor is part of the power control circuit for each plate heater. Plate power is controlled by individual amplifier circuits, one for each plate, which can adjust heater voltage. All of the control circuit elements except the power transistors on the bus bar have been reduced to a printed circuit assembly, 6 channels per board and are mounted in a card box behind a control panel in the instrument console located on the left of the rig, see Fig. 2.2. On this control panel are potentiometers with digital read out faces and small edge voltmeters which indicate heater voltage for each channel. The plate voltage and therefore the power is proportional to the potentiometer setting for that channel. By recording pot settings from a particular run, test conditions can be very nearly repeated by resetting the pots to their former netting. A detailed description of the power control circuits is given in Appendix C. The bus bar on which the power transistors are mounted acts as a heat sink for the power transistors. Under certain conditions these transistors reject a substantial heat load and it is necessary to provide water cooling of the bus bar. This was accomplished by soft soldering 1/2 x 1 inch rectangular copper wave guide into two parallel grooves milled lengthwise in the surface of the 8 foot long, 3 x 3/8 inch copper bus bar. The wave guide is used as a cooling water passage to reduce transistor temperatures. A magnetic valve wired into the starter switch circuit for the welder power supply automatically turns the bus bar cooling water on when the welder is started. A time delay relay set for a one minute delay insures that the bus bar cooling continues until after the welder has coasted down once it has been stopped.

Heater power measurements are made by measuring the voltage drop across the heater and across a precision shunt in the heater circuit.

Heater voltage and shunt voltage leads are carried in shielded pairs to

selector switch read-out stations. These selector switches are 'zone boxed' to prevent any thermally induced, stray signals which could introduce errors in power readings.

These read-out stations are also equipped with VIDAR plug connections which will allow automatic data acquisition at a later date. The ammeter shunts used for power read-out were individually calibrated, in place, against a high precision shunt from the Thermosciences Measurements Center. In addition to individual shunt calibration, the shunt connections were potted with RTV compound to insulate them against possible external thermal effects and each of the 24 shunts were provided with an air cooling system to insure that there would be no stray thermally induced EMF signals.

A.3 The Heat Exchanger Cooling Water System

Cooling water for both the main air and transpiration air heat exchangers is supplied from a single loop which continuously circulates cooling water through both heat exchangers from an insulated holding tank. By maintaining a high water circulation rate, in excess of 20 GPM, temperature gradients across the heat exchangers are minimized. This ensures uniform temperature in the air being cooled. The main air heat exchanger is sized to remove the heat load from the main blower as well as the heater power added in the test section. The transpiration air heat exchanger removes the heat load due to the transpiration blower. Temperature control of the cooling water is achieved by dumping a portion of the heat exchanger return water and making it back up from the water supply main. The make-up water is mixed in the holding tank to provide a buffer against temperature fluctuations in the supply water. It was found that the cooling water temperature dominated the exit air temperature from both the main and transpiration air heat exchangers. This was not unexpected for an air-water system with high effectiveness heat exchangers such as are used here. By holding the circulating water temperature just below the ambient temperature, the exit air temperature from both heat exchangers can be held at the laboratory ambient temperature.

B. The Test Plate Assembly

The test plate assembly forms the bottom surface of the Roughness Rig test section. Its design has been copied after the test plate assembly used in the existing HMT Rig. It is 22 inches wide, 2 inches wider than the inside dimension of the test section, and 96 inches long in the flow direction. It is made up of four subassemblies, bolted to a common support structure. Each subassembly consists of an aluminum casting in which are mounted 6 individual porous plates, each 4 inches wide and 18 inches long. The porous region of the test surface thus is 96 inches long and 18 inches wide. A photograph of a machined casting is shown in Fig. 2.3. A cross section through one of the plates and the compartment in the aluminum casting beneath it is shown in Fig. 2.4.

Transpiration air enters through the air delivery tube from the flow control valve. Like the smooth plate HMT Rig, a one inch delivery tube is used for each compartment. The air jet entering the lower compartment is diverted by a baffle plate to avoid direct impingement on the pre-plate. The air flow in this first inlet plenum may be poorly distributed and may have significant eddies. To protect the working plate, the upper surface of the inlet plenum is a porous bronze pre-plate which provides a pressure drop type damping of the maldistribution. Air passing through the pre-plate is distributed no worse than the variation in permeability of the commercially available pre-plates.

Directly above the pre-plate is the thermocouple location for measuring the temperature of the air delivered to the porous plates. A layer of honeycomb material is fastened to the bottom surface of the porous plates. This honeycomb has hexagonal cells, 3/16 inch in diameter and is 3/8 inch thick. The honeycomb acts as a 'flow straightener' for the flow from the pre-plates. It also eliminates eddies or jets which could cause both flow and temperature maldistributions. The honeycomb not only isolates the plate from contact with eddies in the underbody region, it significantly reduces the radiation view factor between the back of the plate and the underbody. Essentially all that the back of the plate can 'see' through the honeycomb is the pre-plate directly below it.

Thermocouples for measuring the plate temperatures are set into the porous plate with their junctions located .068 inches from the ball crests which form the rough surface. The surface and air temperature thermocouples are 0.010 inch diameter iron-constant thermocouples.

The electric heater wires are stranded, 26 AWG copper wire. Irradiated polyvinyl chloride insulation for the wire was chosen for the temperature tolerance of this material and a series of glue bond tests. The heater wires are glued into groves spaced 0.43 inches apart on the back of the porous plates with a thin bead of Armstrong A-31 epoxy. Each plate is supported along its long edges by a 1/32 inch thick linenreinforced phenolic stand-off. These strips serve to thermally isolate the plates from the castings and minimize heat losses. The phenolic strips are glued to the sides of adjoining plates and inserted into slots milled into the web dividing two adjacent compartments in the casting. Care was taken to provide an air tight seal entirely around each compartment to prevent leakage between adjacent compartments.

Assembly of the plates into the casting is a hand operation since adjacent plates must fit together with a minimum disturbance to the surface roughness pattern. Small irregularities between plates were corrected by hand-finishing and matching each group of six plates. The first step in the assembly was to glue the hand fitted group of six plates together to form a continuous sheet, 18 x 24 inches. Between each pair of plates, the phenolic stand-offs are also glued in place. This gluing assembly is made by clamping the plates together on a precision surface table, under a high clamping load. The next step was installation of the heater wires in each of the plates. The casting and plates were next assembled together, again on a surface table, clamped together. Figure 2.5 shows a photograph of the clamping arrangement used for gluing the casting and plate assembly together. The final step was the installation of the honeycomb, thermocouples, pre-plates and the bottom plates on the castings. On the top surface, balsa filler strips were used to fill the gap between the plate edges and the top surfaces of the casting side rails. Figure 2.6 is a photograph of the back of the plate assembly showing the honeycomb, thermocouples and preplaces. A photograph of the top surface of the casting assembly in place on the support structure is shown in Fig. 2.7.

Five surface temperature thermocouples are installed in each plate. One is in the geometric center of the plate. The remaining four are arranged 1 1/3 inch upstream and downstream and three inches to the right and left of the central thermocouple. This array can sense temperature gradients both in the flow direction and in the transverse direction. The area enclosed by this thermocouple pattern is considered to be the measuring area, the rest of the plate acts as a 'guard' to minimize thermal and hydrodynamic edge effects.

Casting temperatures are measured by thermocouples installed in every other web of each casting at the test section center line. Temperature control of the casting is provided by cooling water tubes in the casting webs. To maintain casting temperature as close as possible to the transpiration air temperature, the casting cooling water is first circulated through a copper tube heat exchanger inside the transpiration header box. The purpose of holding the casting temperature at the transpiration air temperature is to minimize heat transfer between the transpiration air and the casting. Any heat exchange after the temperature of the gas stream has been measured in the casting underbody would introduce an energy balance error. Casting temperature control is also useful to evaluate conduction losses from the plates as will be discussed in Section E of this chapter.

C. Porous Plates

When construction was began on the Roughness Rig, considerably effort went into the problem of plate fabrication. The porous plates used in the smooth plate HMT rig were made by sintering carefully screened bronze particles whose average diameter was .005 inches in a stainless steel mold, under Stanford supervision. This was necessary because commercially sintered plates exhibited permeability variations of 20 to 30%. The decision to use 0.050 spherical particles on the Roughness Rig virtually eliminated sintering as a possibility.

C.1 Plate Fabrication

When construction of the Roughness Rig began, several west coast vendors were contacted and asked to submit samples of sintered porous plates for porosity testing. Only one vendor, the Commercial Filter Corporation of Los Angeles, submitted a sample. This was constructed using bits of a .030 inch bronze wire sintered into an 8 inch square, 1/2 inch thick plate. Tests revealed that the porceity of this sample was highly non-uniform and it was not acceptable. At this same time, the decision was made to fabricate the plates by assembling particles. rather than by a bulk process, and O.F.H.C. copper balls were chosen as ideal elements. The problem then became one of joining the balls together. Brazing seemed most likely to succeed, providing the braze material could be uniformly applied. Discussions with brazing specialists revealed that a nickel-phosphorus brazing alloy could be deposited on the copper balls. This established the technique which was used to fabricate the porous plates used for the Roughness Rig. The plates are constructed from O.F.H.C. copper balls, .050 inches in diameter. The braze material was provided by plating each ball with .0005 inches of electroless nickel. The balls were then arranged by hand into rows and layers inside copper molds in their most dense array and fired in an inert atmosphere furnace to just above the melting temperature of the nickel plating. This resulted in a brazing together of the ball pack into a uniformly porous plate. The plating material, as it melted, formed small fillets between adjacent balls at each contact point. A close-up photograph of the surface of the porous plate is shown in Fig. 2.8. The final dimensions of each plate were 18.0 x 4.0 x 0.5 inches. This method of plate fabrication, although tedious, provided a well defined surface roughness pattern and uniformly porous plate for the transpiration experiments. A more detailed description of the plate fabrication technique is given in Appendix D.

C.2 Plate Heaters

Eight heater grooves were formed in the back surface of each plate by omitting eight of the rows of balls. The allowable heater wire spacing

is governed by the requirement that the temperature of the top surface of the plate remain uniform. The uniformity of temperature across the top plate surface is determined by the wire spacing, the thermal conductivity, the plate thickness and the transpiration flow rate. An analysis of this problem was carried out by Moffat [1] when the original HMT Rig was designed. Using parameters from his analysis, the temperature non-uniformity in the Roughness Rig plates have been estimated to be on the order of 0.01°F. A key parameter in this analysis is the thermal conductivity of the porous plate. In order to estimate the thermal conductivity of the Roughness Rig plates, the ball pack was thermally modeled by a series of resistances representing the ball layers and nickel connecting spots. Several brazed samples were examined under a microscope to estimate the size of the braze connections between adjacent balls. These varied somewhat but a conservative size for the connecting 'spot' diameter was .008 inches. Using this as an average spot size, the conductivity of the ball pack becomes 30 BTU/hr ft *F. This is the value used for the estimate of surface temperature uniformity. It can be argued that this is a somewhat conservative estimate since the arrangement of the balls in the pack are such that a given ball is in contact with 6 surrounding balls on the same layer, but with only 3 balls each in the balls layers above and below. Since the conductivity of the ball pack is controlled by the size and number of contact spots, the ball pack used for the Roughness Rig has a higher apparent thermal conductive across a given ball layer than it does between ball layers.

C.3 Thermocouple Installation

Plate thermocouples were imbedded to a depth that located their junctions at the center of the ball layer below the surface ball layer. Several braze samples were drilled for thermocouple holes and then sectioned to examine the condition of the ball pack around the hole. The brazing process apparently removes all the temper from the OHFC copper balls and they become very soft. It proved difficult to drill the ball pack and the technicians had to carefully prepare the drill bit used for this operation to achieve clean holes. The soft copper

actually wiped itself at the hole surface. The holes appeared to be through solid material when examined in cross section. Upon closer examination, it appeared as if the interstitial regions between ball layers were covered with a very thin layer of copper at the hole surface so the hole appeared to have a solid wall. Further, by drilling the holes on a milling machine, it was possible to locate the bottom of the hole at a ball center in the ball layer below the surface. In this location, the thermocouple junction was .068 inches below the top of the surface layer of balls. Because of the high conductivity of the ball pack, the temperatures recorded by these thermocouples were taken to be the plate surface temperatures. The thermocouples were installed by partially filling the holes with apoxy using a hypodermic needle, then inserting the thermocouple into the hole until it was felt to bottom out. This resulted in a small collar of epoxy squeezing up around the thermomouple at the plate back surface. Every installation was examined from the front surface and there were no cases where bleeding through of the epoxy around the thermocouple holes could be detected.

C.4 Plate Permeability

Uniformity of plate permeability is an important requirement for any transpired boundary layer experiment. Non-uniformities can result in variations in both surface temperature as well as local transpiration rate. Permeability measurements on the plates used in the smooth plate MMT rig were made before their installation. An instrument was developed which sealed off both faces of the plate except for a 3/4 inch diameter spot, and then forced a meterod flow of air through that spot. The pressure drop across the plate provide the measure of the resistance to flow through the test area. With this device, porosity maps were constructed by making measurements on one inch centers over the surface of each plate.

Porosity variation measurements on the Roughness Rig plates were made after the plates were in place in the sub-assembly castings, with thermocouples and heater wires installed, using the actual rig transpiration air supply system. After assembly of each casting was complete, it

was installed on the rig support structure, with the transpiration air supply connected and test section sidewalls but not the test section top in place. A constant current hot-wire anemometer was supported on a traversing mechanism mounted on the sidewalls. It was adjusted to a height of about 1/8 inch above the plate surface. Output from the anemometer was run directly into an X-Y plotter. A resistance bridge built into the traversing mechanism drove the plotter such that the pen movement across the graph corresponding approximately to probe movement over the plate. This set-up was used to map the surface velocity of the air delivered through each plate.

Initial results from these tests yielded velocity traces which appeared very lumpy at high blowing rates. The non-uniformity in velocity was several percent and the spanwise period of the velocity disturbances appeared to be several ball diameters. A literature survey into the subject of flow through porous media uncovered several interesting papers which seem to explain the phenomena that had been observed. Two of particular interest are by Bradshaw [44] and Morgan [45], both of whom studied the flow through screens. Both of these studies showed the existence of a critical velocity in the flow through the porous material. Above this velocity, the flow emerges from the porous media as a pattern of jets. These jets coalesce into random groups because they can only entrain fluid from each other. The velocity field investigated by Morgan was behind a two dimensional grid. It showed the same onset of a lumpy structure with a period that was two or three times the grid spacing. To confirm that the "lumpy" surface velocity measurements were the result of the coalescence of the jets leaving the plate surface, the tests were repeated over a wide range of surface velocities. Results from these traces are shown in Fig. 2.9. Here the surface velocity traverses have been superimposed in a single graph. The top graph shows the results obtained by starting at a high surface velocity and reducing it. The top curve is for a surface velocity of 0.39 ft/sec. It shows the lumpy appearance which was observed over all the plates. The magnitude of the non-uniformities is reduced as the surface velocity is reduced but they are still clearly present at 0.21 ft/sec. The velocity

is much smoother at 0.18 ft/sec and at lower surface velocities. This indicates there is a critical velocity at or near 0.18 ft/sec. The lower graph in the figure shows results obtained when the surface velocities are increased beginning from a low value. Here smooth traces are obtained up to 0.24 ft/sec. Above this, the lumpy traces reappear and persist to higher velocities. Apparently the onset of the flow instabilities have a hysteresis-like response to velocity changes. The jets continue to coalesce to much lower surface velocities as the surface velocity is reduced than the surface velocity at which coalescence begins when the surface velocity is being increased.

As a final step to confirm the mechanism, these same tests were repeated over the porous plates in the existing HMT rig. Again the flow showed the same lumpy structure, but the velocity traces appeared much finer grained than had been observed over the rough surface. This was expected stace the jet pattern over the smooth surface porous plates should be much finer to start with. It was also noted that the onset of flow instabilities occur at surface velocities above the normal operating range of the rig.

The next step was to investigate the possible effects of this jetting action on the heat transfer experiments. It was feared that the jetting action might artificially enhance boundary layer heat transfer. To evaluate this possible effect, a series of eight Stanton number runs with blowing were made at a test section velocity of 90 ft/sec. At this velocity, it was possible to run several blowing fractions above and below the critical surface velocity observed in the permeability mapping. If there was an effect on heat transfer due to the onset of a surface jetting, it should clearly be seen in the heat transfer data. The data was taken and reduced without corrections for radiation or conduction losses. Although it showed some scatter, it clearly indicated that there was no abrupt change in the surface heat transfer. Apparently, the presence of the boundary flow on the porous surface stabilized the transpiration flow and there was no onset of instability or any similar effect that can be seen in the heat transfer data.

The porosity maps for the plates were completed but the flow striations could be eliminated only by reducing the transpiration velocity to low values. At these low levels, the width of the recorded line was equal to about 5% of the velocity signal. This prevented the careful resolution of non-uniformities that had been hoped for. With that reservation, the plate porosity appeared entirely uniform: the recorded traces contained no discernible evidence of non-uniformities.

D. Rig Instrumentation

D.1 Temperature Instrumentation

All temperature measurements on the Roughness Rig are made using iron-constantan (ISA type J) thermocouples. Many precautions have been taken to reduce the spurious EMF's frequently encountered in thermocouple circuits. The thermocouples are all brought together at a common test console zone box where they are connected to rotary thermocouple switches for read-out. A diagram of the thermocouple circuit is shown in Fig. 2.10. The thermocouples were made sufficiently long so that the thermocouple wire itself could be used for the test console lead wires. To avoid introducing sharp temperature gradients in any of the thermocouples. all plate, transpiration air, and casting thermocouples were thermally guarded with Polyflo tubing. The test console zone box is constructed from 1/2 inch plywood, lined inside with 1/16 inch aluminum plate and insulated outside with aluminum foil backed, rock wool insulation. All entry ports are gasketed and all busing of wire connections within the zone box is done with thermocouple wire. These precautions were taken to reduce temperature stratification within the zone box and to minimize its effect. Thermocouples are mounted at opposite ends of the zone box to provide a direct measurement of the zone box temperature gradient. In its present configuration, temperature gradients greater than 1/10°F have never been observed. The entire thermocouple circuit uses a single ice-bath reference junction which is also thermally guarded in Polyflo tubing to stretch its thermal gradients. Thermocouple output is measured with a Hewlett-Packard Integrating Digital Voltmeter, Model 2401C.

The most difficult temperature measurement problem in this system is the plate surface temperature. These thermocouples are glued into

close-fitting holes in the bottom of the plate. The wire extending from the bottom surface is exposed to the transpiration air which is at a different temperature than the plate. This tends to introduce an error in the measurement of the plate temperature. To estimate this error on the smooth plate HMT Rig, Moffat [1] developed a thermal model of the thermocouple treating it like a cylindrical fin, partly exposed to the plate and partly exposed to the transpiration air. The problem is complicated by the fact that the thermocouple may not be perfectly bonded in the hole. This problem is modeled by introducing a 'bond factor' between the plate and thermocouple. Because of the small thermocouple wire diameter and relatively deep immersion into the plate, bonding factors can be taken as low as 0.5 and still have conduction errors less than 0.1°F for a 20°F temperature difference across the plate. Subsequent qualification tests on the plate indicated that the thermocouple error due to imperfect bonding and to other sources are acceptably small. These tests are discussed in Section E of this chapter.

None of the temperature measuring locations on the Roughness Rig are so restrictive in space that conduction error becomes a problem. Ample immersion depth is available everywhere.

As noted earlier, each plate contains five thermocouples arranged in a cross pattern. In the tests performed here, these thermocouples were wired in parallel in the instrument console zone box so that an average of their collective readings were taken. Care was taken to size all the thermocouple lengths the same for each casting so if they were garged together, they would give a true average reading.

D.2 Pressure

Pressures were measured on a variety of manometers and transducers. A 3 inch inclined Merriam Manometer was used to measure tunnel static pressures and to set the top for zero pressure gradient conditions. This same manometer was used to record mainstream total pressure when test section velocities were below 100 ft/sec. At higher mainstream velocities, total pressures were read out with a transducer. Two

Statham unbonded strain gauge differential pressure transducers were used for velocity profile measurements; a PM5 (pressure range 0 to 0.5 psi) and a PM97 (pressure range 0 to 0.05 psi). Both units were equipped with zeroing bridges and individually calibrated in the Thermosciences Measurements Center against a precision 30" Merriam Micromanometer. The H-P 2401 IDVM equipped with an external quartz crystal oscillator clock was used to read the pressure transducers.

D.3 Flow Rate

Transpiration flow rates for each of the porous plates are measured using hot-wire type flow meters which are discussed in Appendix B. The flow signal from these meters is from a differential thermocouple which is read out with the H-P 2401 IDVM using the external clock to extend signal integration time. Before each flow meter reading is taken, the flow meter heater circuit current is checked to insure it is set exactly at the calibration value. This is done by reading the voltage drop across a precision Weston 1 amp shunt, again with the IDVM.

In addition to the flow meter signal, transpiration air temperature for each supply line is measured so the appropriate temperature corrections can be made.

D.4 Electric Power Measurement

One advantage of the D.C. power supply system used in the Roughness Rig is that power measurements are relatively simple. The heater voltage is measured directly with leads which are attached to the heater wives just as they leave the castings. The heater ground connection is made through individual shunts which are mounted beneath the bus bar box. Measurement leads from the shunt and heater are taken in shielded pairs to the same instrument console where the thermocouples are read, but are connected in a separate 'zoned' selector switch station for read out. The heater and shunt voltage are read separately, using the IDVM, and plate power calculations are made in the data reduction program using the individual shunt resistances obtained from an in-place calibration of each shunt. The plate heater power supply voltages have been

carefully scoped to insure that a constant voltage level is applied to each plate over the entire power range. This insures that the selective sampling of the IDVM will yield a truly representative average measurement for the voltage across the heater and shunt.

D.5 Mainstream Conditions

For all runs, mainstream temperature and total-to-static pressure were measured. In addition, mean velocity profiles were taken at several positions along the test section length. Mainstream temperature was measured through the circuitry already described with a probe constructed using .004 inch iron-constantan thermocouple wire. This probe is a fixed position version of the traversing probe described by Kearney [6]. The mainstream total pressures were measured with a Kiel-type probe located in the center of the potential flow region and the static pressures were taken from the adjacent wall tap. All static wall taps were .040 diameter at the wall plane with 0.125 inch diameter tubing connections outside the test section.

Mean velocity profiles were taken using a small diameter boundary layer probe mounted in a micrometer driven traversing probe holder. The probe holder is supported from an instrument sled which is aligned with the instrument ports in the test section top using locating pins and is attached with hold down bolts to the side walls. A photograph of the probe itself is shown in Fig. 2.11. The probe stem is made from 1/8 inch diameter brass tubing. Into the brass stem is soldered a .030 inch stainless steel hypodermic needle which has been bent into a C-shape. The probe mouth has been flattened down to a thickness of .022 inches.

In use, the probe was lowered until it was in visible contact with the rough surface. This could be confirmed with a resistance reading from a VOM attached to the probe and to a thermocouple in the plate the probe was centered over. As the probe was traversed away from the surface, the circuit was broken when the probe lifted from the surface. This method was used to locate the first reading with respect to the crest of the surface balls. The traversing probe was checked at least

once during each run by comparing it with a Kiel probe to insure it was not clogged with dirt particles which would result in erroneous readings. Care was also taken to align the probe before every set of readings to avoid error introduced by probe yaw.

E. Rig Qualification Tests

The Roughness Rig was tested in detail for reliability before being approved for use. There were three types of qualification tests performed: tests of the mainstream condition, instrument system qualification tests and energy balance tests to determine validity of the data reduction program.

E.1 Mainstream Conditions

Extensive uniformity and stability tests were conducted on the test section inlet flow by M. Crawford, another member of the HMT group. The purpose of these tests was to measure mainstream uniformity and local free stream turbulence intensity. The importance of a uniform free stream inlet condition is to insure a two-dimensional flow in the test section. Boundary layer skin friction and enthalpy thickness are inferred from momentum thickness and Stanton Number measurements and the integral momentum and energy equations. A two-dimensional boundary layer is necessary if these techniques are to be used in reducing the experimental data.

To test mainstream uniformity, velocity profiles across the test section inlet were taken using a differential Kiel probe technique. The differential traversing employs two identical probes connected to a differential measuring instrument. This technique allows separation of space-wise and time-wise variation in flow, to a first order at least. If the velocity field can be represented as a product of two functions; one varying slightly in space, the other varying in time, then differential traversing (using probes of the same time constant) will display only space-wise variations on the flow field. Two test section velocities were examined; the lowest velocity, 32 ft/sec and maximum test section velocity, 242 ft/sec. Messurements were taken at 135 positions in the

test section inlet plane at each of the two velocities. The study showed velocity defects no greater $\pm 0.16\%$ at the low test section velocity and $\pm 0.10\%$ at the higher test section velocity.

Free stream turbulence intensity measurement was also made with a DISA 55D05 constant temperature anemometer read through a DISA 55D15 linearizer. The velocity fluctuations about the mean velocity were measured on a Thermosystems RMS voltmeter, Model 1060. Measurements were made at the same two test section velocities and at a total of 120 positions in the test section inlet plane. These results showed a very uniform turbulence intensity over the entire inlet plane with an intensity level of about 0.4% at both test section velocities. This was higher than expected but still below levels measured in the existing HMT Rig. The design objective for the screen pack was a free stream turbulence intensity level of 0.1% to 0.2%. This higher than expected level does not limit the rig's performance.

An additional check on the two-dimensionality of the test section boundary layer was made by measuring its uniformity in the lateral direction at the end of the test section. Lateral holes are provided for this purpose in the test section top over plate ?3. Velocity profiles were taken at the test section center and at two positions on either side of the center for a tunnel velocity at 90 ft/sec. Momentum thicknesses from these profiles are plotted in Fig. 2.12, the variation can be seen to be no more than $\pm 3\%$.

E.2 Instrument System Qualification Tests

To qualify the plate temperature measuring system a scheme was used that had been developed for use on the smooth plate HMT Rig. Rather than calibrate the individual thermocouples, they were calibrated in place for each plate. This was accomplished using a plenum box which could be placed over each test plate in such a way as to collect all the air passing through the center 6-inch span. The plenum box was surrounded by four similar guard chambers and was equipped with a guard heater system so that the walls of the box could be maintained at the same temperature as the air. The entire unit was lined with

aluminized Mylar and insulated with 1/2 inch felt. A mixing section was installed in the plenum box so that the air entering from the test plate would be mixed thoroughly before its temperature was measured with two independent thermocouples. The guard chamber surrounding the plenum box was equipped with adjustable area vent holes so all chambers had the same static pressure to prevent cross flows. The bottom surfaces of the box, where it contacts the plates are covered with closed cell fosm rubber to further prevent cross flow between chambers. In operation, the guard heater system power is adjusted until thermocouples in the box walls indicate that the box is at the same temperature as the mean mixed air temperature, as measured by the mixing chamber thermocouples. This temperature is then compared to the average of the five thermocouples in the plate beneath the plenum box. These tests were conducted over a range of blowing rates from 5 to 20 CFM and over a range of air temperature differences from zero (no plate power) to 20°F. Over this wide range of temperature and flow conditions none of the plate average thermocouple readings differed from the mixing chamber thermocouple readings by more than 0.2°F (6 microvolts). From these tests it was concluded that the plate temperature measurements as recorded by the five plate thermocouples was satisfactory and errors due to imperfect bonding of the thermocouples and other sources are acceptably small.

While these tests were in progress, thermocouple traverses were made under the plates, just below the honeycomb using a special probe installed through fittings in the side of the castings in the center of each plate. These traverses showed a nearly uniform temperature distribution beneath the plate. Typically these profiles showed higher readings at the plate center by 2 or 3 microvolts than at 3 inches to either side of center.

A final series of tests were performed using the transpiration air system with the top removed, but without plate power. These tests were run with both hot and cold transpiration air, by either by-passing or using maximum cooling in the transpiration air heat exchanger. In addition, either maximum or no casting cooling water was used. The

purpose of these tests was to provide a maximum amount of thermal distortion, without plate power, to insure that the temperature measuring system could properly record these unusual conditions. Operating in this mode, the plate thermocouples can be used to check the single. centrally located air thermocouple located just below the plate. Without plate power, the five plate thermocouples read the average temperature of the transpiration air. In cases where there is a temperature mismatch between the casting and transpiration air, the air thermocouple is subject to error. It was discovered that the casting temperature had a much more dominant effect on the transpiration air temperature than had been thought. When the casting inlet air was colder than the casting, the five plate thermocouples indicated a higher reading than the single air thermocouple behind the plate. The opposite was true when air entering the casting was hotter than the casting. Apparently the air next to the casting is heated or cooled by the casting and enters the plates with a temperature gradient. This gradient was not observed in the traverses taken beneath the plates since these were taken in the longitudinal direction. The measurement error due to the inlet temperature gradient is related to the mismatch between the casting temperature and the temperature of the transpiration air as it enters the casting.

The temperature of the transpiration air is measured at the header box, just upstream of the flow meters to provide the temperature corrections for the flow meter readings but not at the inlet to the casting. To estimate the air temperature at the casting inlet, the transpiration system was modeled as a heat exchanger between the air and the laboratory. Thermocouples were installed at the casting inlet of four plates to check this model and good agreement was obtained. Using the casting inlet temperatures estimated with this model, a correction term was formulated for the air thermocouples. The error was assumed to vary linearly with the inlet air-to-casting temperature difference. This correction, while relatively small aided in closure of system energy balance.

E.3 Energy Balance Tests

Energy balance tests were conducted to establish the validity of

the experimental system. The system includes not only the Roughness Rig and its measuring system but the procedure used to reduce data from the rig. Experimental Stanton numbers are obtained by taking plate power, subtracting off the losses and the energy carried away by the transpiration air, and dividing by the plate-to-free stream temperature difference. In equation form this is given by:

St =
$$\frac{(\text{plate power}) - \dot{m}''c (T_w - T_t) - (\text{losses})}{G c (T_w - T_{\infty})}$$
 (2-1)

A key part of the data reduction program is the modeling used to evaluate energy losses. Depending on blowing fractions and plate temperatures, these losses represent from 1% to 10% of the total power supplied. The loss mechanisms modeled are: radiation from the top and bottom surfaces of the plates, conduction between the plates and the casting through the phenolic support webs and conduction through the stagnant air beneath the plates when there is no transpiration. The purpose of the energy balance tests was to evaluate the modeling used for energy losses.

Two kinds of energy balance tests were conducted on the Roughness Rig. Each of these tests provided a direct measure of the ability of the rig's measuring system to produce zero values of measured heat transfer when test conditions were such that there was zero heat transfer.

In the first energy balance mode, the wind tunnel was operated without transpiration and without main stream cooling. In this mode, the
free stream temperature equilibrated near 100°F. Plate power was adjusted
to match the plate temperature to the free-stream temperature. In this
operating condition, the plate power exactly equaled plate losses. These
losses were due to radiation from the back of the plate and conduction
to the casting and to the air beneath the plate. These tests were run
with no casting cooling, to minimize conduction losses, and then again
with maximum casting cooling to magnify conduction losses. Since the
plate and free-stream temperatures were the same, there was no heat
transfer between the plate and the free-stream. The difference between
measured plate power and calculated losses for the observed conditions
represent a potential Stanton number error. Data taken from this energy

balance run were reduced with the data reduction program and the net Stanton number error was evaluated using the energy balance difference and a typical plate-to-free stream temperature difference:

St =
$$\frac{\text{(plate power)} - \text{(calculated losses)}}{\text{G c}(T_w - T_w)_{\text{typ}}}$$
 (2-2)

The typical temperature difference used was 25°F which is representative of the Roughness Rig test conditions. Figure 2.13 shows a plot of Stanton number error versus distance along the test section. There was some difficulty in closing the energy balance in the first few plates, attributed to an inlet thermal boundary layer. With the tunnel containing hotter than ambient, a thin thermal boundary layer of cool air next to the duct walls was unavoidable. Since the plate and free-stream temperatures were equal, the first few plates removed the inlet thermal boundary layer and acted as guard heaters for the rest of the test section.

In the second energy balance mode, the tunnel was operated with transpiration flow only and with the test section top removed so that the transpiration flow moved directly upwards. In this operating condition, there was, again, no surface heat transfer. All the plate power was taken up either in losses or by the transpiration flow. These blowing energy balances were conducted over a wide range of transpiration rates and air-to-plate temperature differences, again to demonstrate the adequacy of the data reduction program to predict the losses and achieve energy balance closure at all operating conditions. For this case the potential Stanton number error was given by:

$$\Delta St = \frac{\text{(plate power)} - \hat{m}''c (T_w - T_t) - \text{(calculated losses)}}{G_{typ} c (T_w - T_w)_{typ}}$$
(2-3)

Blowing energy balance runs were made at 2, 8, 15 and 22 CFM which covers the range over which the Roughness Rig experiments were conducted. As before, the typical temperature difference, $(T_w - T_w)_{typ}$, was taken at 25°F. The typical mass flow, G_{typ} was chosen so that the combination of it with the transpiration rate would give a blowing fraction of .008,

the maximum tested. This was done for all but the lowest transpiration rate. The value of $G_{\rm typ}$ used for both 2 CFM and 8 CFM corresponded to the minimum test section velocity. A plot of the Stanton number error from these tests versus distance along the test section is shown in Fig. 2.13. As can be noted, the Stanton number error was, in almost every case, within $\pm .0001$ Stanton number units. In those few cases where the error exceeds this value, subsequent energy balance tests did close to within $\pm .0001$ level. This has led us to believe that the Stanton number data are reliable to within $\pm .0001$ Stanton number units over the range of conditions tested.

E.4 Verification of the Data from Plate 24

Despite the excellent closure that was obtained in the energy balance tests, an anomaly was noticed in the heat transfer data from plate 24. At conditions of high blowing and high test section velocity, plate 24 exhibited a much lower Stanton number than its neighbor, plste 23. This behavior was entirely repeatable and in fact was visible in earlier data taken when the rig was first being shaken down, but was not as apparent then. After carefully confirming the flow meter calibration and the temperature measuring system it was concluded that the reduced Stanton number was real and not the result of trouble in the measuring system. A review of the possible sources of this problem led to the conclusion that it was related to the diffuser, directly downstream of plate 24. It seemed likely that at high blowing rates, when there was injection of low momentum fluid into the boundary layer, the exit diffuser experienced stall. This stall probably began right at the diffuser inlet but may have propagated upstream to plate 24 itself. The presence of the 'stall bubble' could disturb the flow in the region near plate 24 enough to reduce the plate heat transfer. The exit diffuser was designed to operate without stall, however it is known that the inlet velocity profile has an important effect on diffuser performance. A thickened boundary layer due to blowing could result in stall that would otherwise not have been there. To confirm this possibility a temporary extension of the test section surface was installed in the first stage of the diffuser. This

extension produced a step in the flow path at the inlet to the second stage of the diffuser and reduced the overall performance of the diffuser, but did remove the influence of the diffuser from plate 24. The tunnel was then run at 240 ft/sec with a blowing fraction of .002, without the step. After reaching equilibrium, data were recorded and then the diffuser insert was installed and the run repeated. The insert reduced diffuser performance so that the test section velocity only reached 226 ft/sec. Since the flow meter settings were unchanged, a slightly higher blowing fraction resulted. This reduced the Stanton number slightly below the previous data, but the plate 24 Stanton number was significantly increased. It was still slightly low (.0001 Stanton number units) but much improved over the Stanton number without the insert. Figure 2.14 shows the Stanton number data from these two tests.

Because of this difficulty, the plate 24 data has not been shown for data runs at high velocity and blowing fraction but is listed with the tabulation of experimental results in Appendix E, for completeness.

F. Data Reduction Program

In reducing data taken from the Roughness Rig experiments, three separate data reduction programs were used, two for the hydrodynamic data and one for the heat transfer data.

The first data reduction program was used for the boundary layer mean velocity profiles and converted the pitot tube pressure readings to velocity. It then executed a simple trapezoidal rule numerical integration to get momentum and displacement thicknesses and finally shape factor. The second data reduction program used a least square curve fit to the momentum thickness data for the entire plate length to determine the skin friction, using the integral momentum equation for the two-dimensional boundary layer. It was assumed that over the relatively small momentum thickness range of a single run, the momentum thickness varied as a simple power law with distance along the test section. The equation which was curve fit was:

$$\theta = \mathbf{a}(\mathbf{x} - \mathbf{x}_0)^{\mathbf{b}} \tag{2-4}$$

where x_{λ} is the boundary layer virtual origin and a, b and x are determined by the curve fit. The procedure used was that of 'weighted residuals' as described in Chapter 18 of Scarborough [46]. Since no trip was used in the Roughness Rig experiments both laminar and turbulent regions existed inside the test section. Only momentum thickness measurements taken in the fully turbulent boundary layer were used in the fit. The Stanton number data was then examined to ensure that each of the velocity profiles used had been in the fully turbulent boundary layer. To qualify this procedure for predicting skin friction, the momentum thickness measurements taken by Andersen [8] and by Simpson [2] were used as a trial data set. The skin friction data reported by Simpson were based on his 'best estimate' using several evaluation methods. Skin friction predictions by Andersen were made using turbulent shear measurements taken in the boundary layer, then extrapolated to the wall. Figure 2.15 compares skin friction predictions made by the present curve fitting method to the values originally reported. For both data sets the curve fitting method yields a reasonably good check of the skin friction values that Simpson and Andersen found by independent means. Since both of these data reduction programs are straightforward applications of well known principals, listings of the programs are not given.

The data reduction program for the heat transfer data is more complicated since it must convert data from the Roughness Rig into Stanton numbers. A flow diagram for the data reduction program is shown in Fig. 2.16 and a Fortran listing is given in Appendix F. The program has several features. It contains the individual calibrations for each flow meter and performs a logarithmic interpolation to obtain the flow from each flow meter reading, based on the calibration data points. Thermocouple readings are converted into temperatures by interpolating in tables taken from National Bureau of Standards Circular #561. Plate voltage and shunt voltages are converted into power using calibrated resistance values for each shunt and correcting for excess heat release due to the turnaround wrappings of the heater wires at the plate ends. It also estimates conduction and radiation losses from each plate based on plate, casting and other rig temperatures. The program also converts manometer readings

to pressure levels and the free-stream thermocouple reading into the free-stream total temperature using a 0.36 recovery factor for the thermocouple probe. The program has two output options: the summary output shown in Appendix E and a more detailed output with additional diagnostic information.

G. Experimental Uncertainties

To investigate the possible effects of measurement uncertainty on the experimental Stanton number, the following procedure was used. The data reduction program was used to calculate the derivative of each experimental variable with respect to the Stanton number. This was done by sequentially varying each variable by the amount of its uncertainty as input to the data reduction program and calculating the change in Stanton number which resulted. Assuming that these uncertainties are all independent they can be combined by the procedure of Kline and McClintock [47].

The results of this uncertainty analysis is shown in Fig. 2.17 for all test section velocities, over the range of blowing in these experiments. As can be seen, the lowest test section velocity, at high blowing is the most affected by the propagation of uncertainties in these experiments. The following besic uncertainty intervals were assumed for the input values:

Plate Voltage	l millivolt
Shunt Voltage	10 microvolts
All Temperatures	10 microvolts
Stagnation Pressure	.002 inches of water
Flow Meter Readings	25 microvolts

It can be concluded from this study that except at high blowing the uncertainty in Stanton numbers due to random errors is less than the \pm .0001 Stanton number units cited earlier based on the energy balance tests.

Anothe area of experimental uncertainty associated with these experiments is associated with the velocity profile measurements. A problem which is common to all rough surface boundary layers is the location of the origin of the velocity profile. Profile measurements begin with the probe resting on top of the rough surface. The location of this data

point with respect to the surface is known, but examination of the profile data indicates that the origin of the profile is somewhere below the ball crests. Several authors have discussed systematic procedures for determining 'origin error'. What has been done in the present work is to examine each profile on an individual basis and extrapolate : to what appears to be a reasonable choice for its velocity origin. Generally this turned out to be .010 inches below the ball crests. The bounds on this number would be from zero (the profile origin at the ball crests) to a maximum of .025 inches, the profile origin at half of the ball diameter. With these bounds, it seems unlikely that the experimental uncertainty of the profile origin could exceed +.010 inches. To investigate the sensitivity of the integral boundary layer parameters to an uncertainty in the profile origin of $\pm .010$ inches, the velocity profile data reduction program was used. A portion of the profile data was reduced again with a .010 inch shift in the profile origin. It was not surprising to find that the displacement thickness was very sensitive to the shift while the momentum thickness was not sensitive. There was no detectable effect of test section velocity: this is not surprising since the boundary layers were all very similar at all test section velocities. A more important parameter was the thickness of the boundary layer. Near the front of the test section where the probe could not penetrate as deeply into the thin boundary layer, an origin error of .010 inches producted errors on the order of 10% in the boundary layer displacement thickness and 1% in the momentum thickness. At the end of the test section, where the boundary layer was thicker, the uncertainty in profile origin produced errors of 7% to 8% in displacement thickness and again, 1% in momentum thickness. From these results, we have concluded that the experimental uncertainty in the profile virtual origin results in uncertainties on the order of 19% in the boundary layer displacement thickness and 1% in the momentum thickness.

Some additional work has been done to translate the momentum thickness uncertainty into uncertainty in the skin friction prediction. This can be done by varying the inputs to the program which uses a fit to the commentum thickness data to predict skin friction. The results from

study have been inconclusive. The effect of the momentum thickness uncertainty depends on how many other points used by the curve fit remain unchanged and even on which point is perturbed. Another very important parameter is the blowing fraction, since it is subtracted from the momentum thickness before the skin friction fit is made. A general result which was obtained is that a 1% variation in momentum thickness, for the non-blown boundary layer could produce a varation of 8% to 10% in skin friction in a boundary layer with 5 or more velocity profiles as basis for its least square fit. In many cases, the variation was less and in cases with high blowing the variation was greater. Because of the many difficulties inherent in taking the derivative of experimental data points, it is felt that no further claims about the accuracy of the skin friction data generated by fitting the momentum thickness measurements was justified. Experimental work now in progress on the Roughness Rig should produce much better skin friction data, with much less experimental uncertainty.

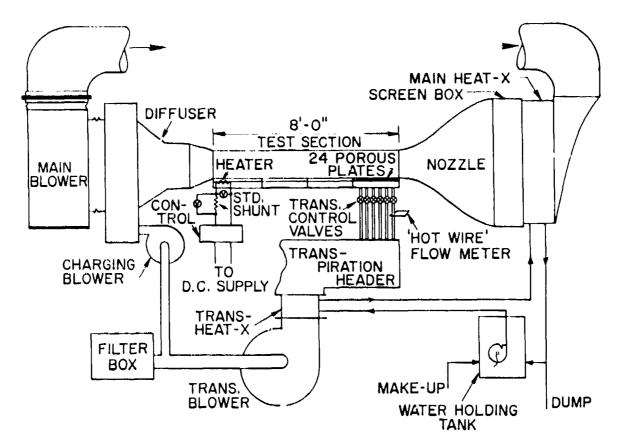


Fig. 2.1 Rough surface wind tunnel flow schematic

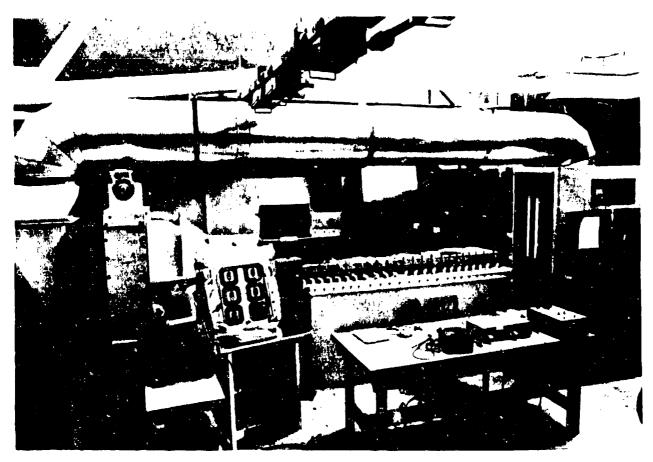


Fig. 2.2 Photograph of Roughness Rig

1. POROUS PLATE
2. HEATER WIRES
3. THERMOCOUPLES
4. SUPPORT WEB
5. TRANS. AIR
THERMOCOUPLE
6. HONEY COMB
7. CASTING WATER PASSAGE
8. PRE-PLATE
9. TRANS. AIR DEFLECTOR
10. CASTING THERMOCOUPLE
11. CASTING THERMOCOUPLE
12. TUBE
13. THERMOCOUPLE
14. THERMOCOUPLE
15. THERMOCOUPLE
16. THERMOCOUPLE
16. THERMOCOUPLE
17. THERMOCOUPLE
1

Fig. 2.4 Cross section view of typical compartment

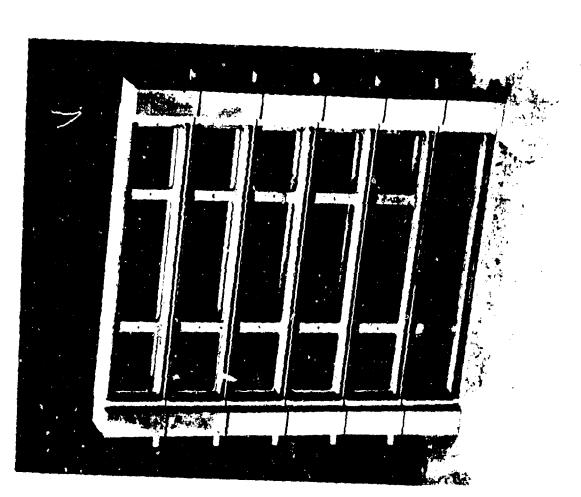


Fig. 2.3 Photograph of aluminum casting



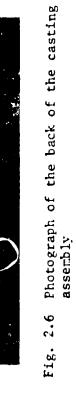




Fig. 2.5 Photograph of clamping arrangement used for gluing the plate and casting assembly

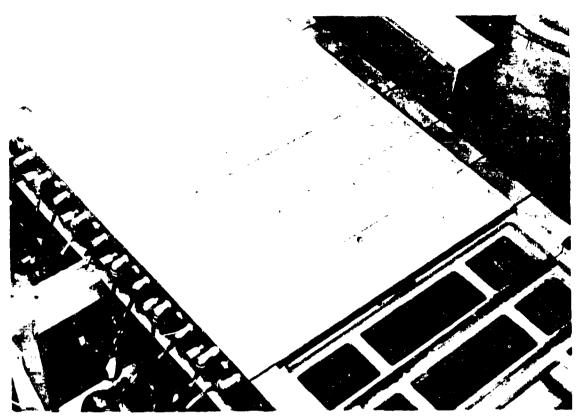


Fig. 2.7 Photograph of the top of the casting assembly

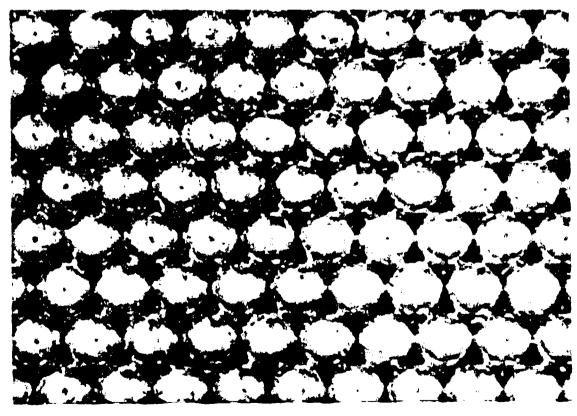
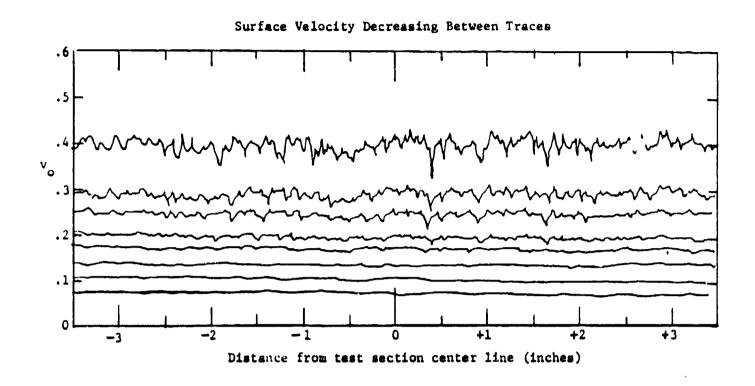


Fig. 2.8 Photograph of the rough surface, close-up



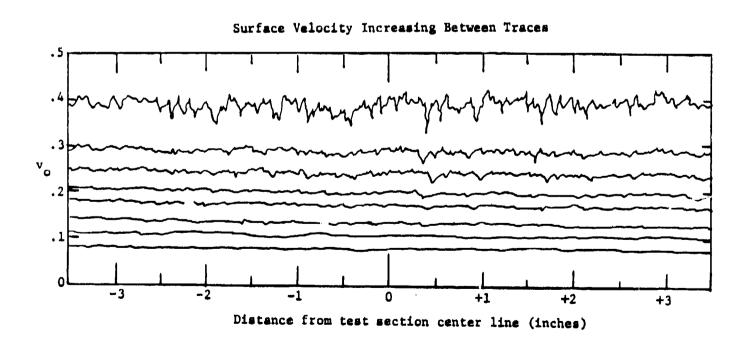


Fig. 2.9 Velocity distribution above the Roughness Rig surface

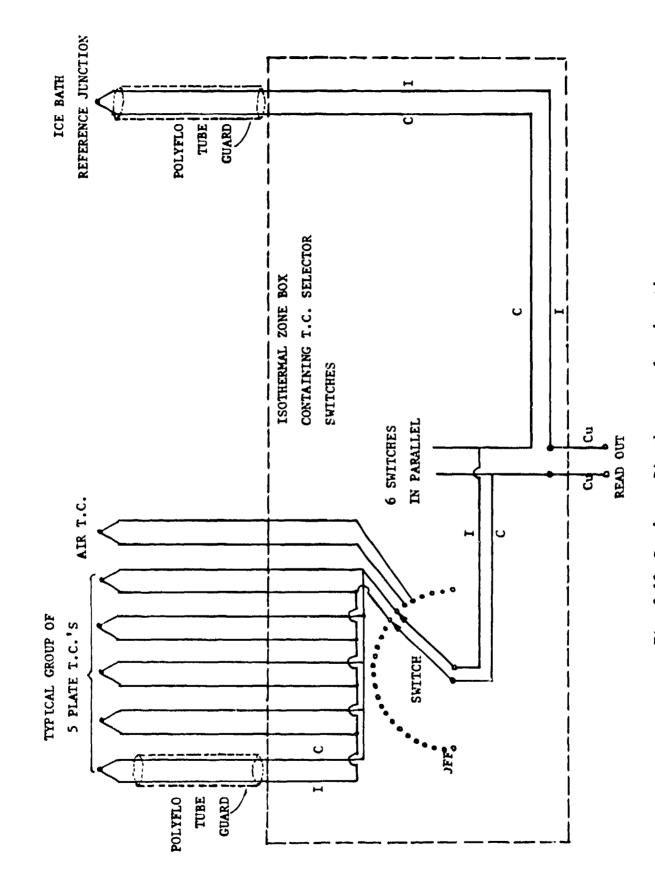


Fig. 2.10 Roughness Rig thermocouple schematic



Fig. 2.11 Photograph of the Pitot probe

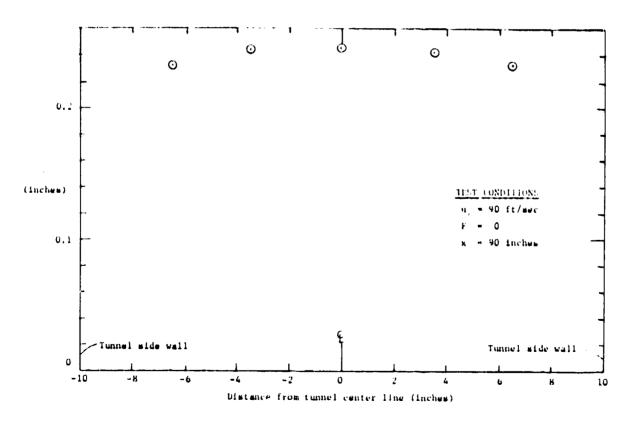


Fig. 2.12 Momentum thickness distribution across the tunnel at plate 23

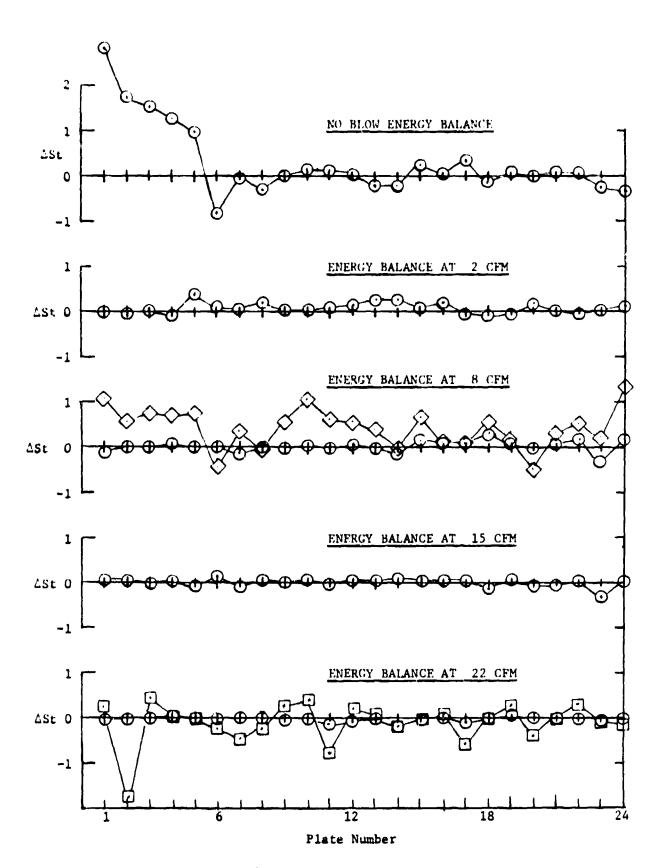


Fig. 2.13 Non-blown and blowing energy balances

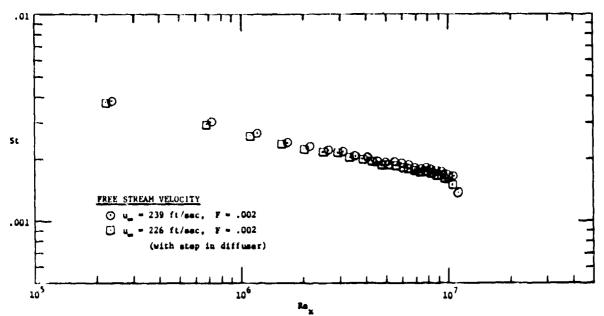


Fig. 2.14 Stanton number versus x-Reynolds number for $u_{\infty} = 239$ fps, F = 0.002 (standard) and for $u_{\infty} = 226$ fps, F = 0.002 (with an insert in the diffuser)

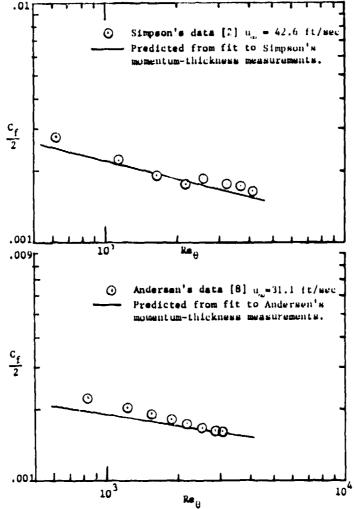


Fig. 2.15 Comparison of skin friction predictions to data by Andersen [8] and Simpson [2]

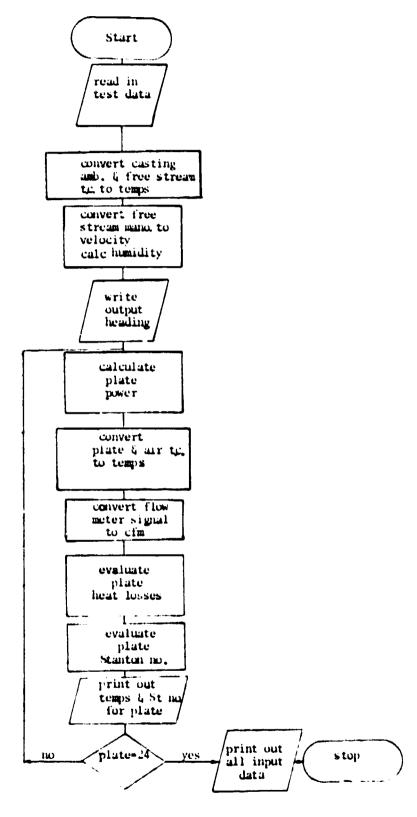


Fig. 2.16 Data reduction program flow chart

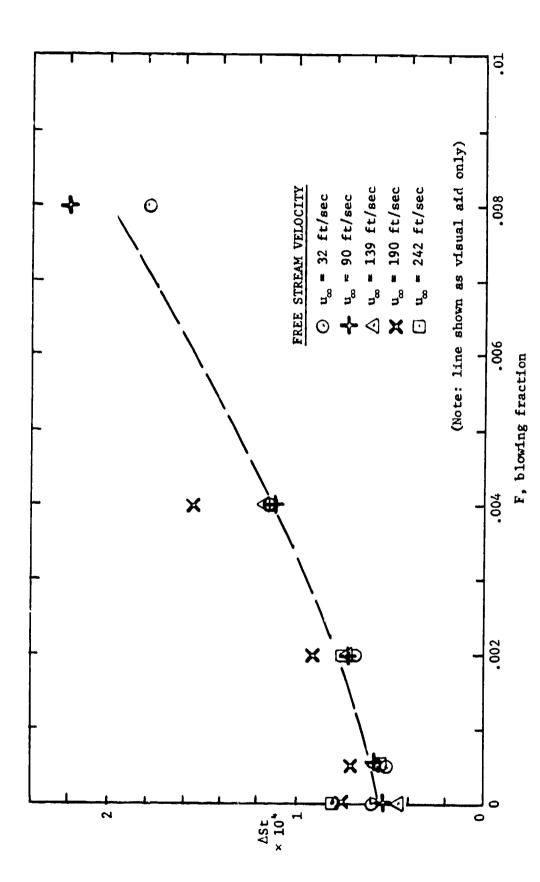


Fig. 2.17 Stanton number error due to experimental uncertainties

CHAPTER III

EXPERIMENTAL RESULTS

Measurements taken in these experiments can be divided into two types -- Stanton number data and mean-velocity profile data which were used to infer boundary layer momentum thickness and skin friction.

The range of conditions for the Roughness Rig tests can be summarized as in Table 3-1. The data taken can be organized into the following blocks: fully developed turbulent boundary layer skin friction and Stanton number without blowing, the effects of blowing on rough surface Stanton number and friction factor, and rough surface transition with and without blowing. The presentation of the data will be organized in the same manner.

A. Rough Surface Boundary Layer Data without Blowing

Rough surface boundary layer friction factors and Stanton numbers are plotted against x-Reynolds number in Fig. 3.1 for each of the five test-section velocities used. It is apparent that there is a roughness effect and that, in these coordinates, the friction factor and the Stanton number increase with test section velocity. This is an important difference from smooth plate data, which do not show a velocity effect in these coordinates. Smooth plate data for air in this range of x-Reynolds number would correlate, for all test-section velocities, with equations of the form:

Smooth
$$c_f/2 = .0295 \text{ Re}_x^{-0.2}$$
, (3-1)

For comparison, these two correlations are shown in Fig. 3.1. Only the 32 fps data seem to be approaching the smooth plate correlations.

With transition occurring at different locations for each velocity, it might be argued that the failure to correlate is due to different locations of the boundary layer 'virtual origin'. To avoid this problem,

Table 3-1

Roughness Rig Test Conditions

Nominal Test Section Velocity	32	06	135	190	242
Blowing Range (m'/G)	0 to .008	0 to .008	0 to .004	0 to .004	0 to .002
x-Reynolds Number × 10 ⁻⁶ (measured from test section inlet)	.03 to 1.63	.03 to 1.63 .09 to 4.23 .19 to 6.60 .44 to 9.11	.19 to 6.60	.44 to 9.11	.49 to 11.23
Momentum Thickness Reynolds Number × 10 ⁻³ (for no blowing)	.24 to 3.1	.41 to 10.6	.84 to 17.0	.41 to 10.6 .84 to 17.0 1.41 to 23.9 1.52 to 30.5	1.52 to 30.5
Roughness Reynolds Number (for no blowing)	24 to 29	64 to 79	101 to 124	142 to 160	178 to 200

boundary layer data are frequently presented in local coordinates.

Fig. 3.2 shows friction factor and Stanton number as functions of local momentum thickness Reynolds number and enthalpy thickness Reynolds number, respectively. In these coordinates, too, the effect of the rough surface has been to separate the data taken at different test section velocities. Also shown are typical correlations for a smooth plate in local coordinates:

Smooth
$$c_f/2 = .0128 Re_{\theta}^{-.25}$$
 (3-3)

Smooth St = .0128 Pr^{-0.5} Re_{$$\Lambda$$} = .0153 Re _{Λ} -.25 (3-4)

It is apparent that boundary layer thickness Reynolds number is not a sufficient descriptor for the rough surface data. At any boundary layer thickness Reynolds number, higher velocities yield higher values of friction factor and Stanton number. While the data at 32 fps lie close to the smooth correlations, at the highest test section velocity the friction factor is nearly twice the expected smooth plate value. At that velocity the Stanton number is up by about one and one-half times. This behavior is similar to that reported by Nunner [23] for rough pipes. His results, given by Eqn. (1-1), show the increase in heat transfer due to roughness approximately equal to the square root of the increase in skin friction.

The present data are clearly affected by the surface roughness. As already noted in Chapter I, the roughness Reynolds number, $\mathrm{Re}_{\tau} = \mathrm{u}_{\tau} k_{\mathrm{g}}/\mu_{\mathrm{s}}$ is often used to categorize the behavior of rough surfaces. Values less than 5 to 10 define 'smooth' behavior, while between 10 and 70 the surface is described as 'transitionally rough', and above 70 it is referred to as 'fully rough'. The roughness Reynolds number range for these tests is shown in Table 3-1. The 32 fps data lie between 24 and 29, in the 'transitionally rough' range, while all other velocities have 'fully rough' roughness Reynolds numbers. For these calculations, k_{g} , the equivalent sand-grain roughness was taken as 0.625 times the ball diameter, yielding an equivalent sand grain roughness of .031 inches. This multiplier is the value recommended by Schlichting [51] for densely packed spheres.

The effects of roughness are also evident in the rough surface velocity profiles. Fig. 3.3 shows velocity profiles from these experiments at different test section velocities. In Fig. 3.4, velocity profiles are shown for one velocity and soveral different stations along the test section, along with a typical velocity profile for a smooth surface. The rough surface velocity profiles show the characteristic offset or depression when plotted in wall coordinates. The amount of the depression, Δu^{\dagger} , can be used to determine the equivalent sand grain roughness of the surface, as described by Schlichting. Applying this method to the non-blown velocity profiles from these experiments, the ratio of the equivalent sand grain roughness to the ball diameter is found to be very nearly the same for all the velocity profiles lying between 0.60 and 0.63. This confirms the value recommended by Schlichting for a similar surface, but tested in duct flow instead of boundary layer type flow.

One difficulty in interpreting rough surface velocity profiles arises in determining the profile origin. Typically, the apparent surface of the rough plate lies below the tops of the roughness elements. Different investigators have handled this problem in different ways. Tsuji and Iida [19] measured velocity profiles from the crests of the roughness elements. Others, such as Moore [12], Perry [14,15] and Liu [16], place the profile origin below the crests of the roughness elements. Perry has argued that the distance below the rough element crests should be adjusted until the velocity profile plotted in semi-log coordinates exhibits the familiar 'log' region. This approach establishes both the all shear and the 'logarithmic asymptote', which is the distance below the roughness element crests to the profile origin. All of the methods which have been suggested require some judgment by the individual examining the data. In the present work, still another method was used: origins for the velocity profiles taken in these experiments were determined by extrapolating the data to zero in cartesian coordinates, by eye. Typically, the origins for the turbulent velocity profiles fell about .010 inches below the crests of the balls. Laminar and transitional profiles were depressed less. Using the origins determined

in this manner, the profiles were integrated to obtain momentum thicknesses which were then used to evaluate skin friction. This approach provies a convenient basis for determining origins for all the profiles and yield results similar to what would have been obtained using Perry's method for the fully turbulent profiles.

One interesting result obtained from the velocity profile data is that the apparent roughness of the plate seems independent of both distance along the plate (boundary layer thickness) and of free-stream velocity (viscosity effects). Thus the roughness effects seem to scale on roughness particle size, not on 'relative particle size'. This is different from rough pipe data, where roughness effects scale on $(k_{\rm g}/d)$, the ratio of the roughness size to the pipe diameter.

Early rough surface boundary layer experiments demonstrated that the skin friction becomes a function of x/k_g alone. It follows from this that the momentum thickness also must be a function of x/k_g alone for the zero-pressure gradient flow if the surface roughness is not changing. Fig. 3.5 shows measured values of θ versus $(x-x_0)$ for each of the test section velocities, x_0 being the appropriate virtual origin. The growth curves of all the boundary layers are substantially the same. Also shown is the expected behavior of a smooth plate at the same test-section velocities. Although the 32 fps data are somewhat low and the 190 and 242 fps data show a slight shift at the downstream end, the coherence is good compared to the smooth plate curves. The two-dimensional integral boundary layer equation for flat plate unblown flows is:

$$\frac{c_{f}}{2} = \frac{d\theta}{dx} . \qquad (3-5)$$

From Fig. 3.5 it can be seen that $d\theta/dx$ is substantially independent of velocity, above 32 fps, being the same function of θ for all velocities. From this observation $c_f/2$ is mainly a function of θ alone. This might have been anticipated from early data which showed $c_f/2$ was only a function of x/k_s and independent of velocity.

 S_{kin} friction, deduced by differentiating the individual curves of θ versus x, is shown in Fig. 3.6 plotted against the ratio of momentum thickness to ball radius. The 32 fps data are clearly low, and there

remains a small velocity dependence in the other data. At the two highest velocities, the skin friction lies 10% to 15% above the other data at the downstream end of the plate. It's difficult to assess the significance of this, based on the present data. An error in measurement of θ of only 3% to 5% could produce the differences observed in the friction factors. Further, smooth surface data, if plotted in these same coordinates would place the low-velocity skin friction above high-velocity values. This suggests that the velocity dependence observed here is probably not an approach to the 'fully rough state'. Finally, the Stanton number data show no such tendence, and they are less sensitive to errors which may be present in the skin friction data.

Figure 3.6 also shows the Stanton number data for all five velocities used, plotted versus Δ/r , where Δ is the enthalpy thickness of the boundary layer. The data are coherent for all velocities in the turbulent region, including the data for 32 fps. After an initial overshoot of a few percent, the 32 fps Stanton number data fall on top of the data for all higher velocities. The turbulent data show no velocity dependence; the Stanton number is a function only of Δ/r , for all velocities. It is interesting to compare the values of Stanton number and friction factor at the same boundary layer thicknesses (i.e., $\Delta/r = \theta/r$). Friction factors for 90 and 140 fps tend to lie about 5% above the Stanton number values all along the plate, rather than 15% low, as in the smooth plate case. The agreement between the 32 fps data and those for higher velocities was surprising, in view of the friction factor results. Furthermore, the roughness Reynolds number for that velocity lies between 22 and 30, hardly 'fully rough'.

To further investigate the Stanton number behavior at low roughness Reynclds numbers, a special test was conducted at a free-stream velocity of 18.8 fps. The boundary layer was augmented by blowing at F = .004 over the first two feet of the test section. The boundary layer was then allowed to relax into its natural state. Stanton number data for the unblown plates are shown in Fig. 3.7, along with the mean of the 90 fps data. It is notable that the low velocity Stanton number data form a natural extension of the unblown 90 fps data. The roughness Reynolds

number at the end of the test section for the 18.8 fps run was only 14, as determined from hot wire measurements of the shear stress near the surface. This is far below the 'fully rough' condition and nearly 'smooth' by the usual criteria. It seems clear that 'fully rough' behavior of the heat transfer data persists to much lower roughness Reynolds numbers, for the present surface geometry, than had been expected. The classification of 'rough' or 'smooth' based on roughness Reynolds number does not seem to be a reliable indicator for the heat transfer performance of rough surface tested here.

In the turbulent range, the present data fit the following correlation:

Rough St =
$$.0043 \left(\frac{\Delta}{r}\right)^{-0.25}$$
 (3-6)

The usual smooth-plate correlation in terms of enthalpy thickness Reynolds number can be cast in a similar form:

Smooth St =
$$\frac{0.0153}{\left(\frac{\mathbf{u}_{\infty}\mathbf{r}}{\mathbf{v}}\right)^{-.25}}$$
 (3-7)

In this form the Stanton number for both rough and smooth behavior is substantially the same at 32 fps. At 18.8 fps, the smooth plate correlation predicts a value nearly 50% higher than the rough. This is because of the inverse dependence of the smooth correlation on free-stream velocity. Clearly, the smooth plate correlation in this form does not provide a suitable limit for the rough surface behavior.

In summary, the unblown data in Figs. 3.5, 3.6 and 3.7 seem nearly independent of free-stream velocity. Stanton number appears entirely independent of velocity, being only a function of enthalpy thickness. Skin friction may be independent and at most has a small dependence. Skin friction, when θ/r is fixed, is only about 5% higher than Stanton number at the same value of Δ/r . The effective value of the Reynolds analogy factor is about 0.95, instead of 1.15, which is typical of a smooth plate. The boundary layer seems to be completely turbulent, with no discernible molecular effect. The effective value of the turbulent

Prandtl number seems near unity. Roughness effects on heat transfer persists to very low values of the roughness Reynolds number, with no evidence of the 'transitional roughness range' in the Stanton number data. Momentum thickness and friction factor data do, however, seem to show a transitional effect.

B. The Effects of Blowing

Blowing diminishes the Stanton number and skin friction for a rough plate, just as it does for a smooth plate. Roughness Rig skin friction and Stanton number data at each of the five test section velocities are shown in Fig. 3.8 through 3.17. There data exhibit very similar trends to the smooth plate skin friction and heat transfer data presented by Moffat [1] and Simpson [2].

An important feature of the unblown turbulent boundary layer Stanton-number data is its lack of dependence on free-stream velocity. This independence is preserved with blowing. Fig. 3.18 shows Stanton number data for all the blown boundary layers at all velocities plotted versus Δ/r parametric in F. Although the scatter is increased, the fully rough blown turbulent boundary layer exhibits the same independence of velocity as does the unblown boundary layer.

In view of this agreement, it is of interest to examine the roughness Reynolds number range of the data, since this is so often used as the measure of the rough surface behavior. At 32 fps and F=0.002, the roughness Reynolds number is about 17 at the point farthest downstream. For all higher velocities or lower blowing fractions, Re_{τ} is greater. Higher blowing reduces the wall shear still further, and, for 32 fps, takes the roughness Reynolds number down to 10 (F=0.004 at the 20th plate). No reliable friction data are available for F=0.008 at 32 fps, but Stanton number data for F=0.008 are shown in Fig. 3.18. The 32 fps Stanton number data at both F=0.004 and 0.008 are higher than the remaining data, and it is felt that they may represent a different state of the boundary layer than the higher velocities. The scatter in these data sometimes (though not often) exceeds \pm 0.0001 Stanton-number units. No cause has been assigned.

Values of Stanton number with blowing are predictable from the unblown values by a Couette relationship used for smooth plate flows, with proper interpretation:

$$\frac{\operatorname{St}}{\operatorname{St}_o} = \left[\frac{\ln(1+B)}{B}\right]^{1.25} (1+B)^{.25} . \tag{3-8}$$

The fully rough prediction is to be made at constant Δ , whereas the smooth plate prediction is made at constant Re_{Δ} . Predictions based on this equation are shown as solid lines in Figs. 3.19 through 3.23. The agreement is excellent for all data more than about 30 boundary layer thicknesses past the peak of the transition hump. Data points more than 30 boundary layer thicknesses past transition are accented in Figs. 3.19 through 3.23. The application of the Couette flow estimator based on enthalpy thickness instead of enthalpy thickness Reynolds number is very reasonable in view of the success of this coordinate in organizing the Stanton number data.

The skin friction data with blowing show a high degree of scatter and like the unblown data it is difficult to draw any definitive conclusions from them. It is interesting, however, to examine the momentum thicknesses of the blown boundary layer. These measurements are less prone to uncertainties and provide a better basis for interpreting the rough surface behavior. Fig. 3.24 shows all of the momentum thicknesses, plotted against distance from the virtual origin, parametric in F. It is clear that in these coordinates, the boundary layers for all velocities are substantially the same. This leads one to suspect that in terms of local thickness coordinates the skin friction should be the same. There is so much scatter in the blown skin friction data that this cannot be confirmed. It does seem likely, however, that the skin friction has only a small dependence on free-stream velocity, since that was observed in the unblown data. More accurate skin friction measurements will be taken in the continuing experimental program on the rough surface boundary layer, and this question will undoubtedly be answered in the future.

The curve fitting technique used to determine skin friction did not work well at blowing fractions greats than .004. The skin friction is

so small at these high blowing rates that even small uncertainties in the momentum thickness measurement result in differences in skin friction that are percentage—wise large. Because of this, skin friction predictions for the .008 blowing fraction runs have not been included in the figures. The skin friction at 90 fps and F = .002 was accorded some special treatment. The curve fit technique predicted a virtual origin for the boundary layer that was inconsistent with the other data at that velocity. In this one particular data set, the virtual origin was selected by extrapolating the boundary layer momentum thicknesses to zero in cartesian coordinates. The skin friction was then evaluated with the same curve fit routine but with fixed virtual origin. In all other data sets, the curve fitting routine was allowed to select both the virtual origin and the curve from which the skin friction data were determined.

The effects of blowing on heat transfer and skin friction to a 'fully rough' turbulent boundary layer can be summarized as follows. Stanton number is a function of Δ and F only, and the blown value can be predicted from the customary predictive equation, but evaluated at constant Δ instead of constant Re_{Δ} . The effects of blowing on skin friction are less well defined, but boundary layer thickness measurements indicate that the growth of the momentum boundary layer is substantially independent of velocity and a function only of the blowing fraction. Based on this and the behavior of the unblown skin friction data, it seems likely that the rough surface blown skin friction has at most small dependence of free-stream velocity. In view of the data scatter, more definitive conclusions are not possible.

C. Roughness and Blowing Effects on Transition

All boundary layers in these experiments were allowed natural transitions without a boundary layer trip. There are two advantages of 'tripping' the boundary layer and having the virtual origin at the test-section inlet: the momentum boundary layer origin is matched to the thermal boundary layer, and it makes maximum use of the test section length. However, when the well-defined laminar boundary layer and transition region were found at the 32 fps test section velocity, it was decided

that the opportunity to observe the effects of roughness and blowing on transition more than offset any advantages offered by tripping the boundary layer. There is very little experimental information in the literature about transition of the blown boundary layer on a rough wall. The present experimental program was not specifically designed to study transition, but some interesting observations could be made from the data concerning rough surface boundary layer transition.

The effect of the laminar boundary layer on Stanton number can clearly be seen in the non-blowing data in Fig. 3.1. Increasing the free-stream velocity moves the transition upstream. It is interesting to note, however, from Fig. 3.2 that the transition is well correlated by momentum thickness Reynolds number. Transition occurs at a momentum thickness (and enthalpy thickness) Reynolds number of about 400.

The effects of blowing on transition can be seen using the data for 32 fps in Fig. 3.10. Shown on this figure are momentum thickness Reynolds numbers at the onset of transition. The data clearly show that blowing moves the transition upstream but does not much affect the momentum thickness Reynolds number (about 400) at which the transition occurs. It is important to note that the thermal protection offered by blowing does not offset the effect of an early transition. Blowing through a region which would otherwise have remained laminar can result in a much higher heat load than no blowing at all.

The present data indicate that, like pipe flow data, below a certain critical Reynolds number the rough surface boundary layer remains laminar. For these experiments, transition occurs at or near the same Reynolds number that would be expected for a smooth surface transition. There are signs of transition in all of Sranton number data, even at the higher test section velocities.

A final remark that can be made about the laminar-to-turbulent transition observed in these tests is that once initiated, it occurs in a relatively short distance. Smooth plate boundary layer transition is often spread over a larger Reynolds number range. Fig. 3.25 compares smooth plate transition data by Reynolds [53] to the present experiment. The smooth plate Stanton number after transition approaches the asymptotic turbulent value slower than the rough surface data from these

experiments. A possible explanation of this is that the surface roughness elements may aid the transition process once initiated and help bring it to completion over a shorter section of the plate.

D. Comparison with Other Rough Surface Experiments

Direct comparison of skin friction and heat transfer data to similar rough surface experiments is useful because it demonstrates that the pressent results are similar to what has been found in other rough surface experiments. Several investigators have compared their rough surface skin friction to the 'fully rough' skin friction correlation suggested by Prandtl and Schlichting [10].

$$c_f = (2.87 + 1.58 \log x/k_g)^{-2.5}$$
 (3-9)

Moore's [12] data compare reasonably well with this equation, but he found different size roughness elements were not organized in these coordinates. The data by Liu [16] were reasonably well organized, but the predicted skin friction was about 20% higher than the data. Wu [18] measured average skin friction over a 10 × 21 inch test plate which agrees well with the Prandtl-Schlichting equation for the average skin friction. Skin friction measurements from these experiments are shown in Fig. 3.26. Like Liu's, the present data Lie about 20% lower than the Prandtl-Schlichting curve.

In a more recent and extensive study, Lakshman and Jayatilleke [52] examine nearly every available set of rough surface skin friction data from both pipes and plates. They correlate the data in E and Re_{τ} coordinates, where E is defined by the equation:

$$u^{+} = \frac{1}{\kappa} \ln(Ey^{+})$$
 (3-10)

Data from these experiments are shown in Fig. 3.27 in E versus Ret coordinates. Also shown is the line for sand-grain roughness. Although the data lie above the sand-grain roughness line in the transition roughness Reynolds number range, it is reasonably well organized in these coordinates and agree well at 'fully rough' roughness Reynolds numbers.

A direct comparison of the heat transfer results with other rough surface experiments is more difficult because previous rough surface heat transfer experiments have been performed in pipe geometries. The data correlations of these tests typically include pipe dimensions or have their roughness scaled by pipe diameters. The heat transfer data by Dipprey and Sabersky [24] were correlated by the following expression:

$$\frac{1}{\sqrt{c_f/2}} \left[\frac{c_f/2}{St} - 1 \right] = g(Re_{\tau}, Pr) - A(Re_{\tau}) . \qquad (3-11)$$

For fully rough flow, the g and A functions are simplified to give:

$$\frac{1}{\sqrt{c_f/2}} \left[\frac{c_f/2}{St} - 1 \right] + 8.48 = 5.19 \text{ Re}_{\tau}^{0.2} \text{ Pr}^{0.44} = g_{fr}(\text{Re}_{\tau}, \text{Pr}) . (3-12)$$

Dipprey and Sabersky were able to correlate all of their fully rough data in g_{fr} and Re_{τ} coordinates. Owen and Thompson [25] recommend the following expression for heat transfer from a flat plate:

$$\frac{1}{St} = \frac{u_{\infty}}{u_{\tau}} \left(\frac{u_{\infty}}{u_{\tau}} + \frac{1}{B} \right) , \qquad (3-13)$$

where B^* is the sublayer Stanton number. Based on their experiments and data by others, the following correlation was recommended for the sublayer Stanton number:

$$B^* = \frac{1}{\alpha} Re_{\tau}^{-.45} Pr^{0.8}$$
, (3-14)

with a lying between 0.45 and 0.7 but represented best by 0.52. The Owen-Thompson expression can be cast into a form similar to the Dipprey-Sabersky correlation:

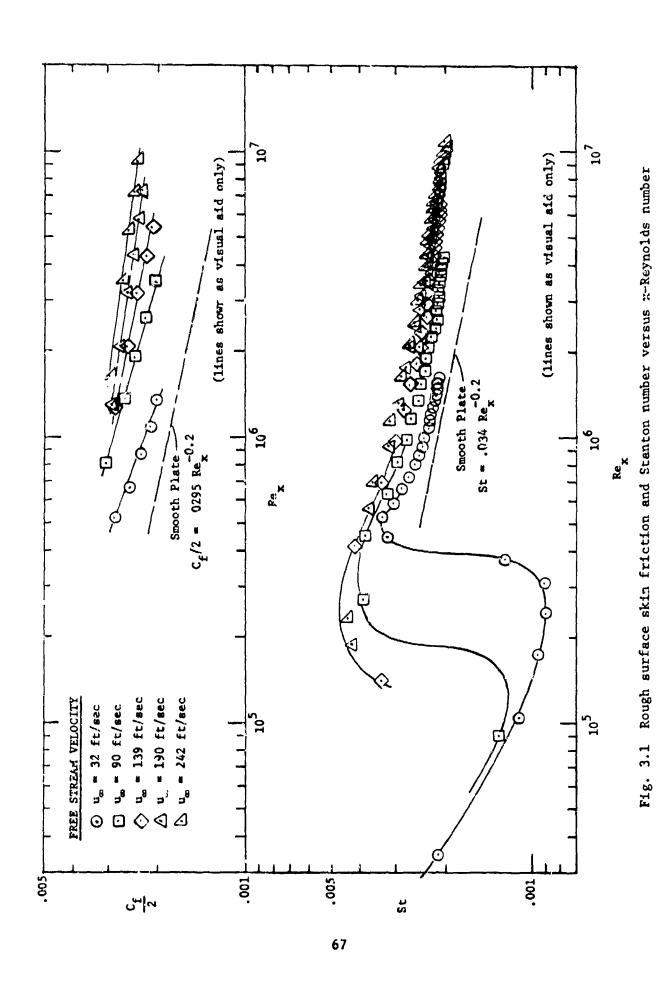
$$\frac{1}{\sqrt{c_f/2}} \left[\frac{c_f/2}{St} - 1 \right] + 8.48 = 0.52 \text{ Re}_{\tau}^{0.45} \text{ Pr}^{0.8} + 8.48 . \quad (3-15)$$

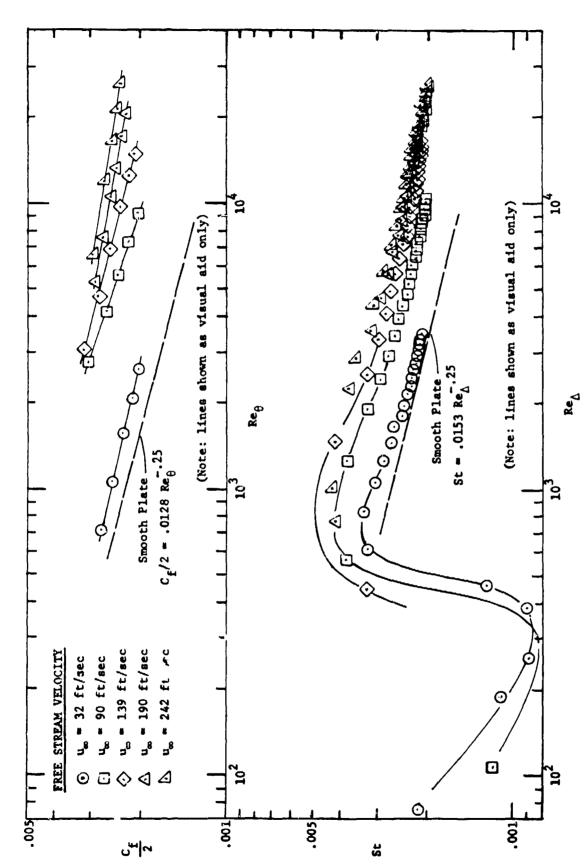
These two models for rough surface heat transfer are plotted in Fig. 3.28 along with data from the present experiments. Both correlations agree reasonably well for roughness Reynolds numbers greater than 70. The data from the present experiments fall below these predictions.

Both of these correlations employ a skin friction-to-Stanton number ratio, minus one. The correlation in this form emphasizes experimental uncertainties of the present data by looking at differences between nearly equal experimentally determined values.

Another useful data comparison can be made using the correlation suggested by Nunner [23]. This correlation compares the increase in rough surface skin friction to the increase in heat transfer. Data from these experiments is shown in Figure 3.29. Also shown is Nunner's correlation and the more recent correlation recommended by Norris [30]. The data from these experiments is best organized by Nunner's original correlation.

Based on these comparisons, it appears that the rough surface skin friction and heat transfer data from the present experiments are comparable with other similar experiments.





Rough surface skin friction and Stanton number versus momentum and enthalpy thickness Reyvolds number F1g. 3.2

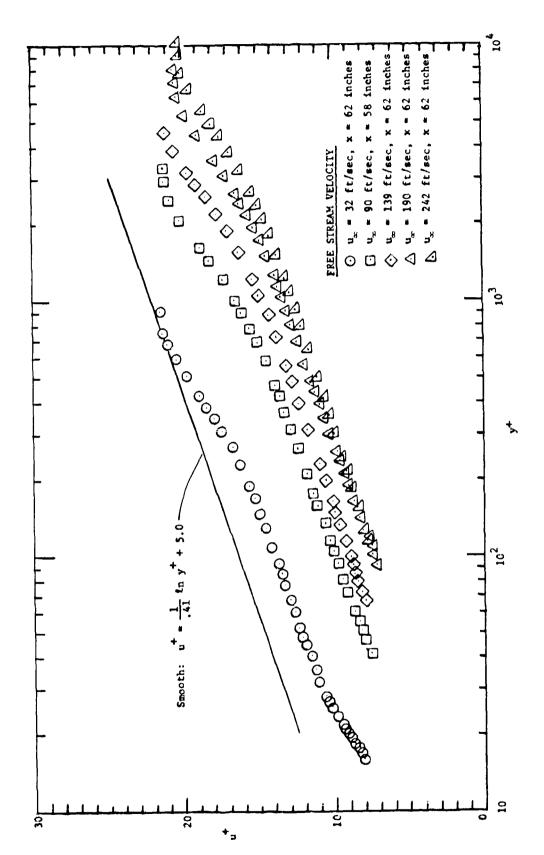
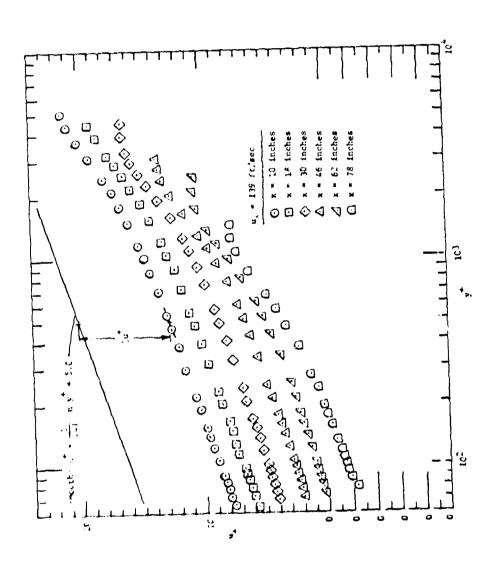
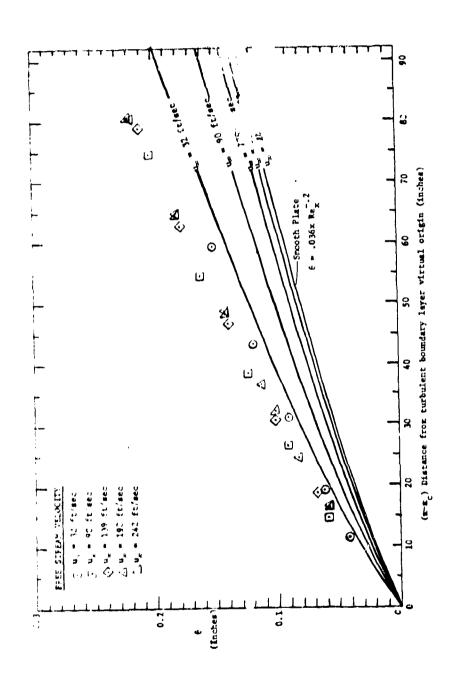


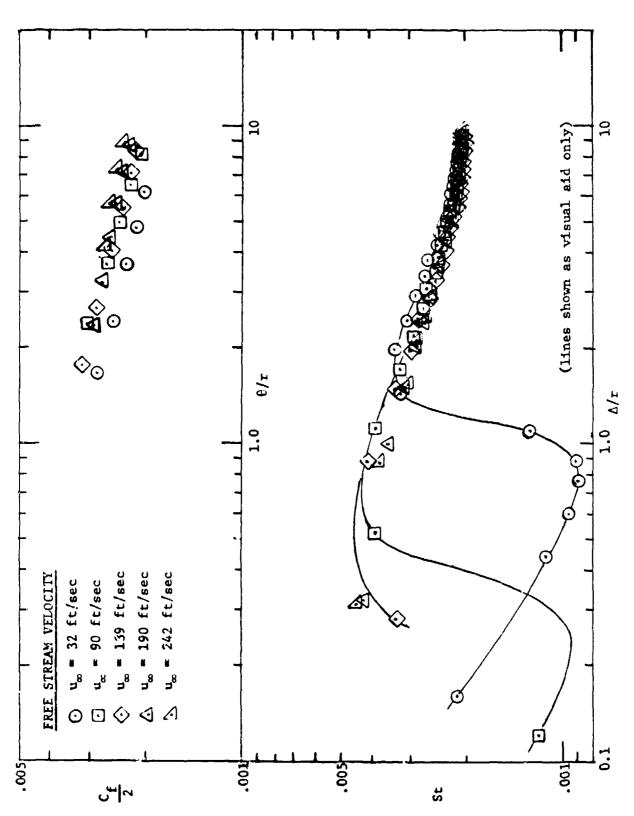
Fig. 3.3 Velocity profiles at different test section velocities



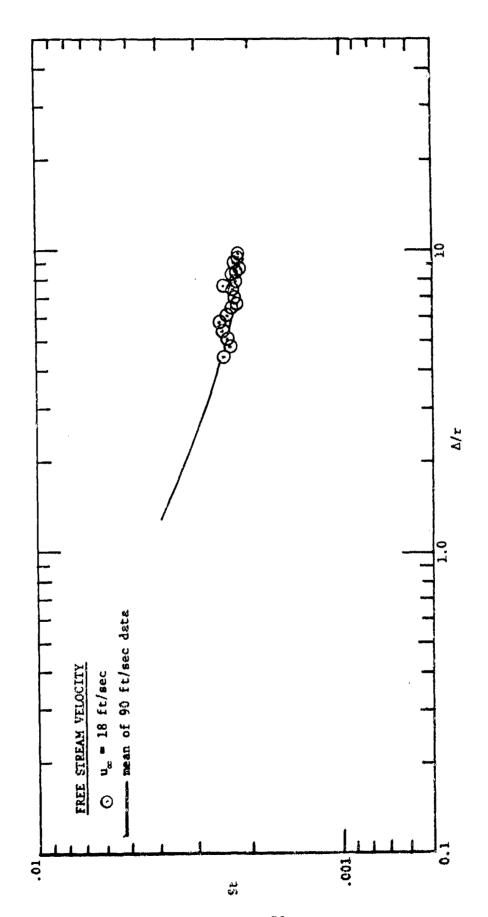
= 139 fps at different stations along the test section Fig. 3.4 Velocity profiles at



Rough surface momentum thickness versus distance from the boundary layer virtual origin Fig. 3.5



and Fig. 3.6 Rough surface skin friction versus (momentum thickness)/(ball radius) rough surface Stanton number versus (enthalp; thickness)/(ball radius)



Rough surface Stanton number versus (enthalpy thickness)/(ball radius) u_{∞} = 18.8 fps and u_{∞} = 90 fps Fig. 3.7

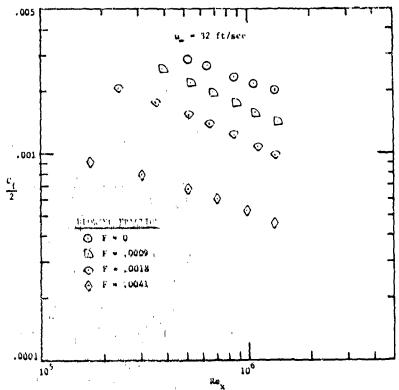


Fig. 3.8 Rough surface skin friction versus x-Reynolds number at u = 32 fps

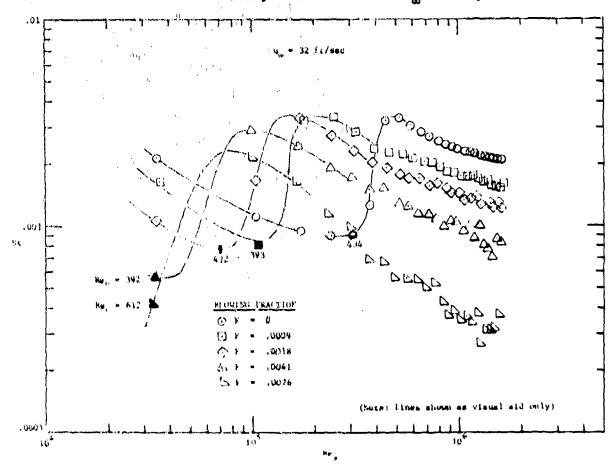


Fig. 3.9 Rough surface Stanton number versus x-Reynolds number at $u_{\omega} = 32$ fps 74

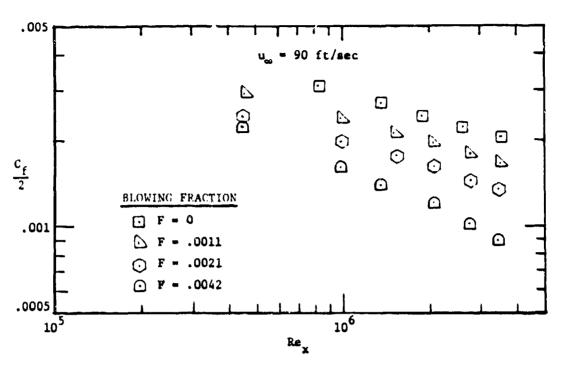


Fig. 3.10 Rough surface skin friction versus x-Reynolds number at u = 90 fps

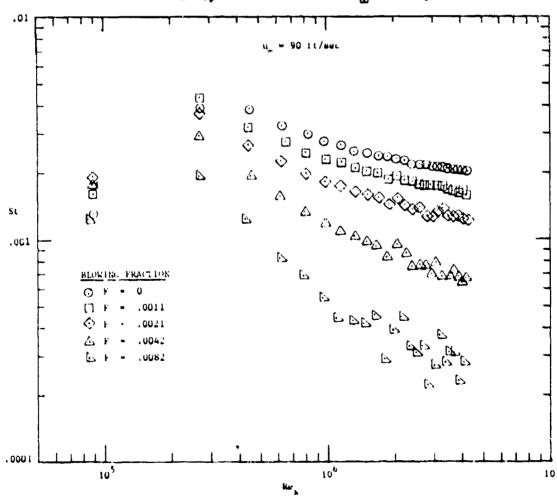


Fig. 3.11 Rough surface Stanton number versus x-Raynolds number at u_ = 90 fps

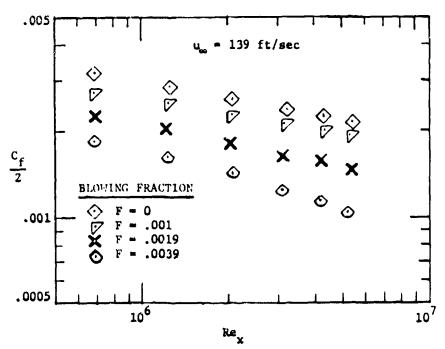


Fig. 3.12 Rough surface skin friction versus x-Reynolds number at $u_{\infty} = 139$ fps

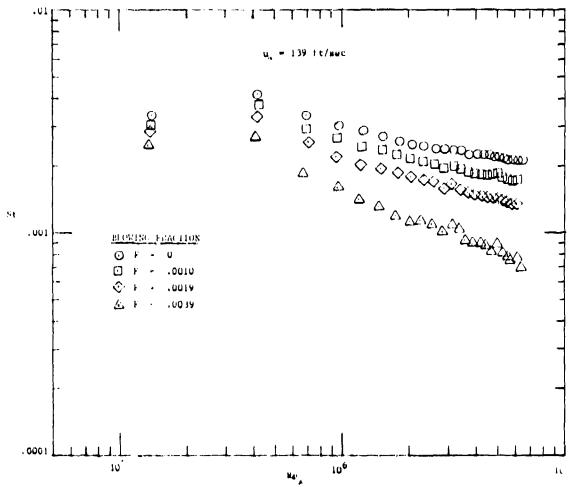


Fig. 3.13 Rough surface Stanton number versus x-Reynolds number at $u_{\rm m} \approx 139~{\rm fps}$

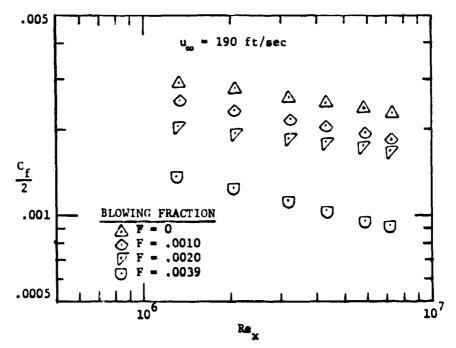


Fig. 3.14 Rough surface skin friction versus x-Reynolds number at u_m = 190 fps

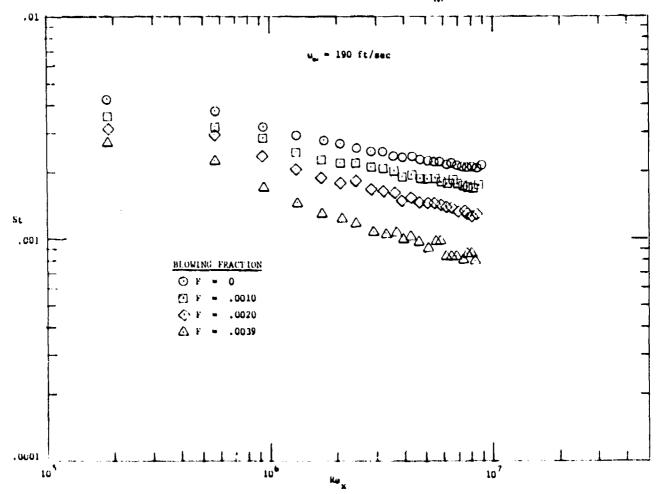


Fig. 3.15 Rough surface Stanton number versus x-Reynolds number at $u_{\rm eq} = 190$ fps

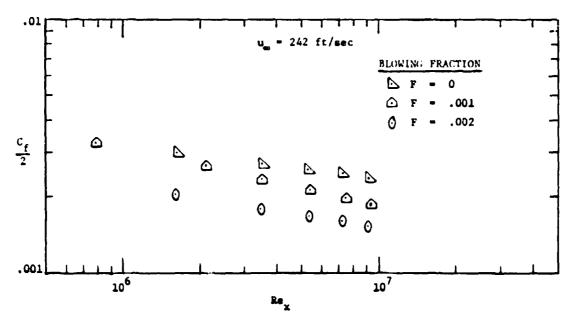


Fig. 3.16 Rough surface skin friction versus x-Reynolds number at $u_{\infty} = 242$ fps

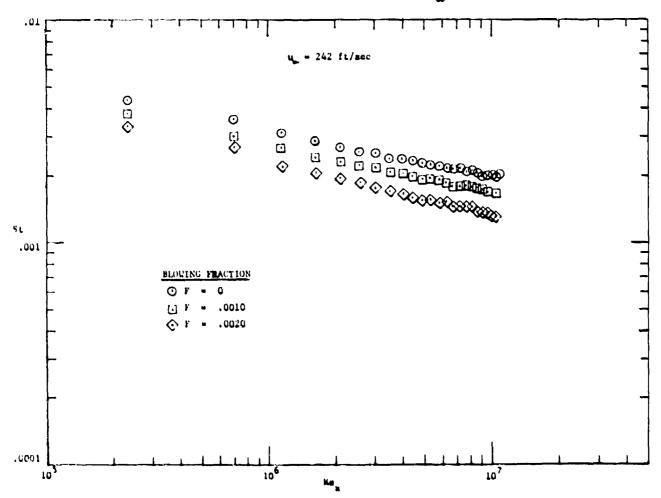


Fig. 3.17 Rough surface Stanton number versus x-Raynolds number at $u_{\rm ru} = 242$ fps

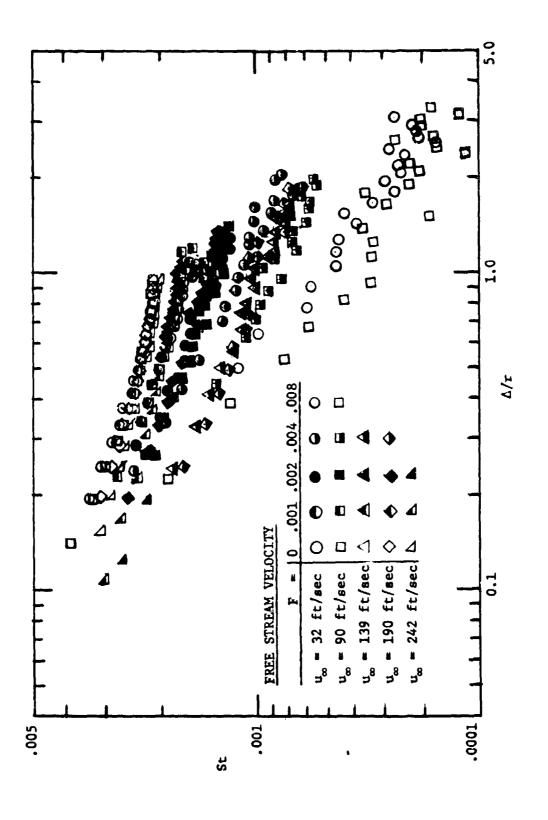
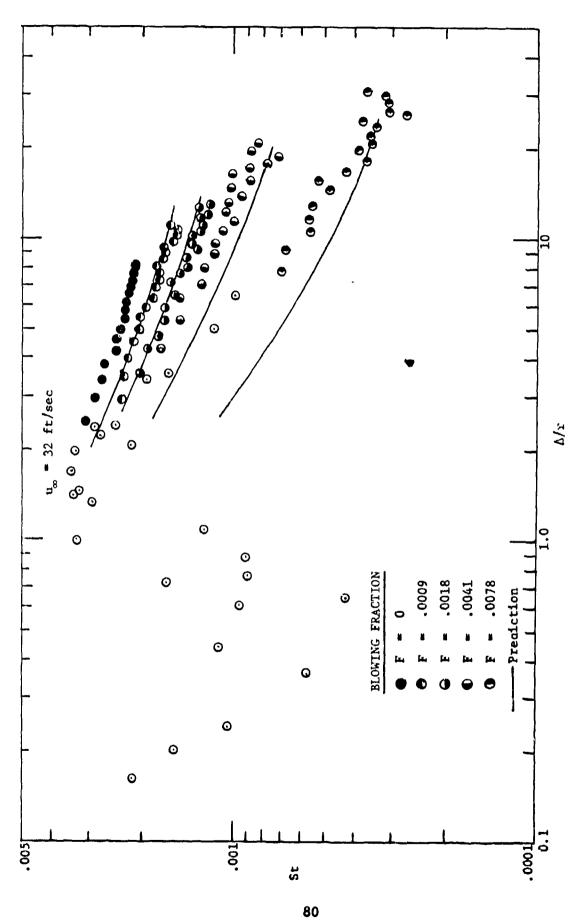
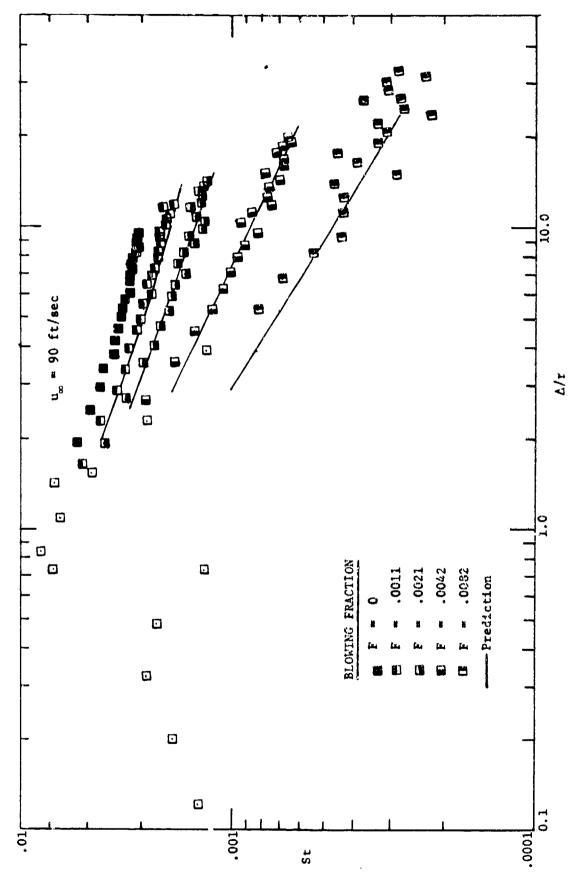


Fig. 3.18 Stanton number versus (enthalpy thickness)/(ball radius) for all boundary layers



 $u_{\infty} = 32 \text{ fps}$ Fig. 3.19 Stanton number versus (enthalpy thickness)/(ball radius) at



u = 90 fps Fig. 3.20 Stanton number versus (enthalpy thickness)/(ball radius) at

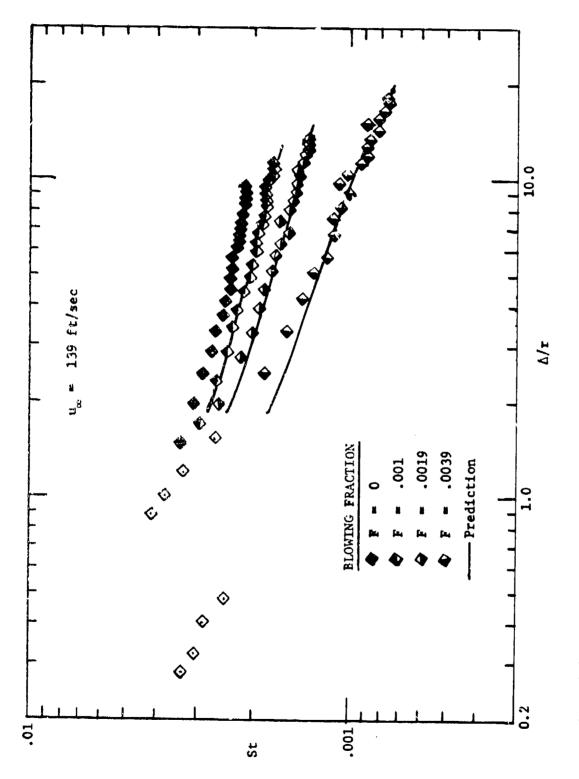
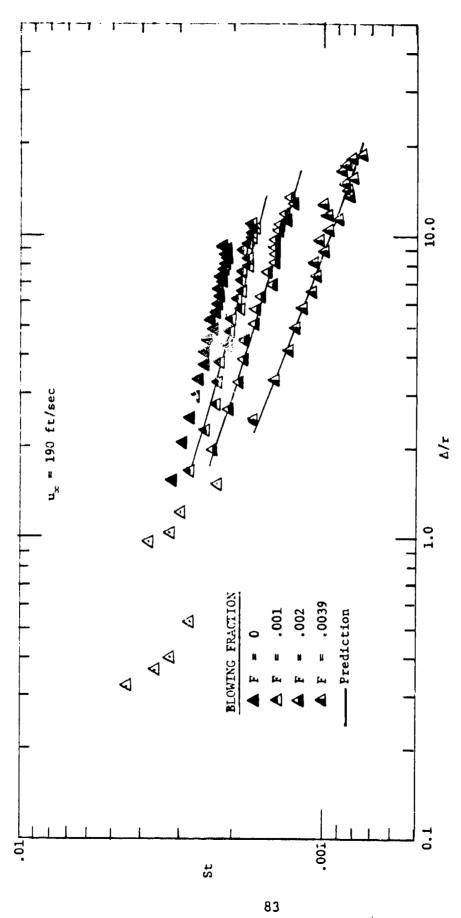


Fig. 3.21 Stanton number versus (enthalpy thickness)/(ball radius) at u_{∞} = 139 fps



Stanton number versus (enthalpy thickness)/(ball radius) at Fig. 3.22

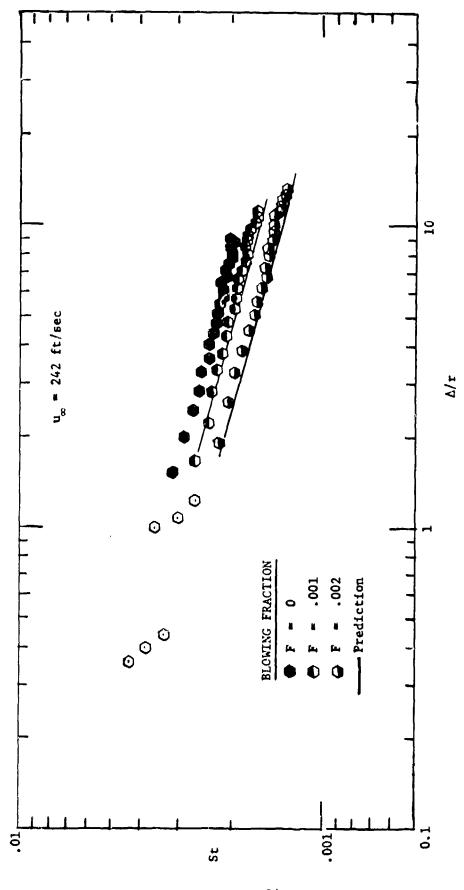


Fig. 3.23 Stanton number versus (enthalpy thickness)/(ball radius) at u_{∞} = 242 fps

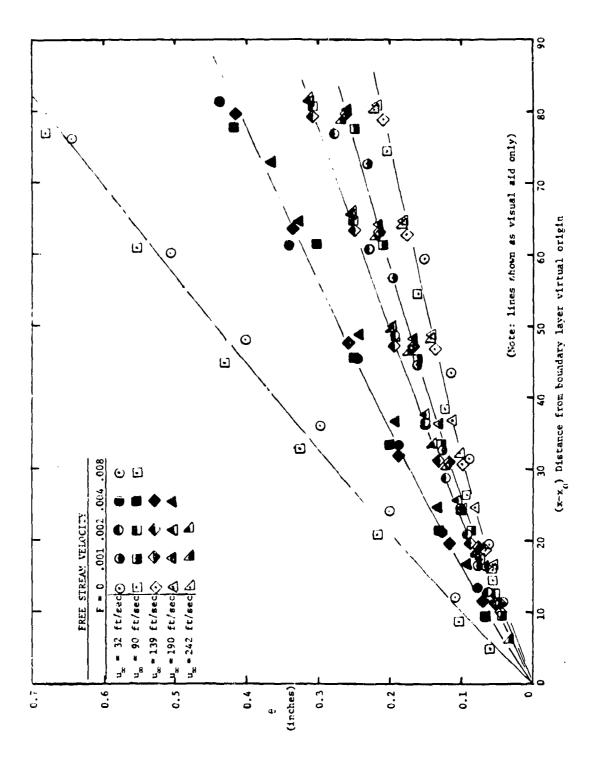


Fig. 3.24 Momentum thickness versus distance from virtual origin for all boundary layers

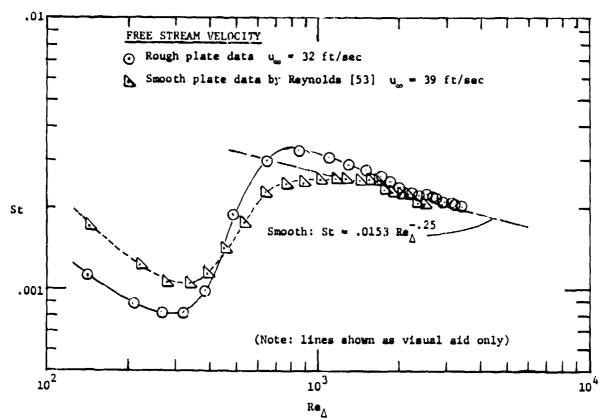


Fig. 3.25 Comparison of rough surface transition Stanton number to smooth plate transition data by Reynolds [53]

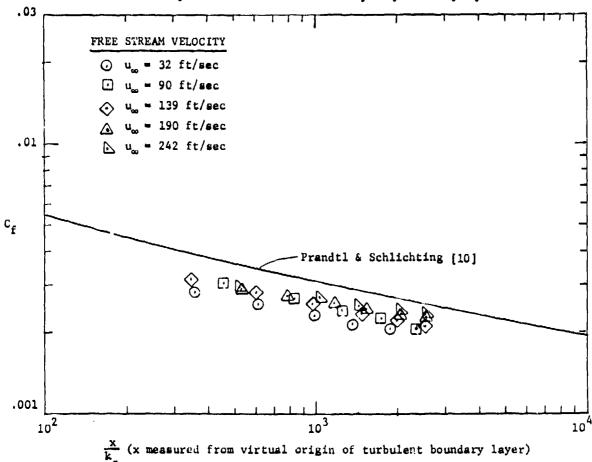


Fig. 3.26 Comparison of skin friction data with Prandtl-Schlichting correlation 86

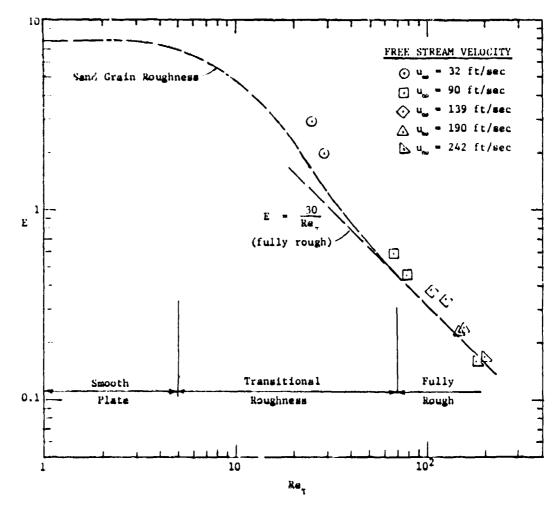


Fig. 3.27 Comparison of skin friction data in E versus $Re_{_{\mathrm{T}}}$ coordinates

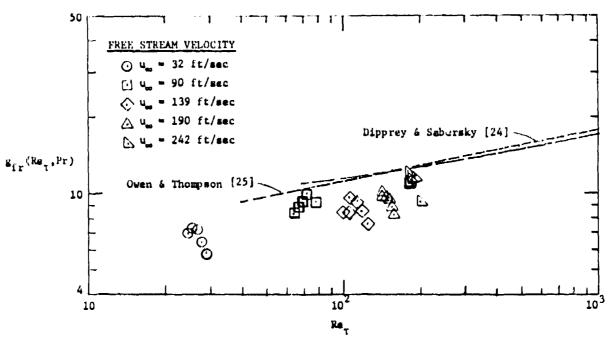


Fig. 3.28 Comparison of heat transfer data to correlations by Dipprey and Sabersky [24] and Owen and Thompson [25]

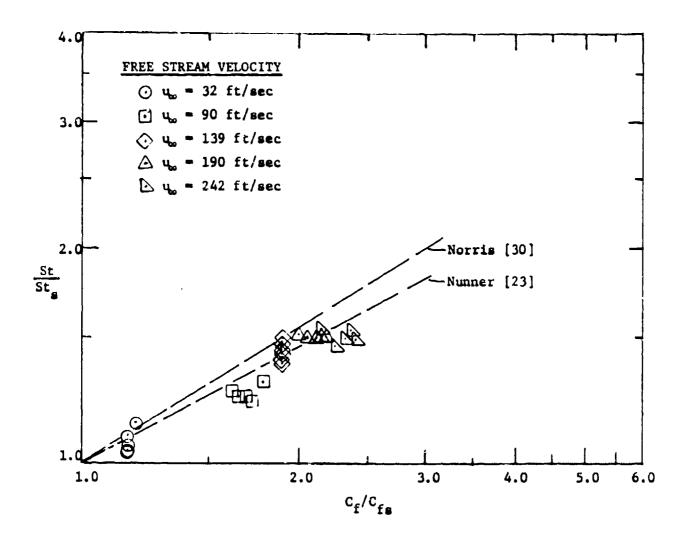


Fig. 3.29 Comparison of heat transfer data to correlations by Nunner [23] and Norris [30]

CHAPTER IV

PREDICTION OF THE ROUGH SURFACE DATA

Predictions of the rough surface boundary layers have been made using a finite-difference prediction scheme based on the computer program originated by Patankar and Spalding [37]. The turbulence model used in this program has undergone substantial development at Stanford in order to equip it to predict smooth surface, transpired turbulent boundary layers with pressure gradients. A brief description of the program is given by Kays [38], who also gives examples of its use.

A. Hydrodynamic Predictions of the Rough Surface Boundary Layer

Closure of the boundary layer equations is provided by a mixinglength scheme using a modified van Driest model in the inner region. This model has been extensively and successfully used by several similar prediction programs. In wall coordinates, it is given by:

$$\ell_{\rm m}^+ = \kappa y^+ (1 - e^{-y^+/A^+})$$
 (4-1)

Near a smooth wall, the mixing-length is 'damped' to allow viscous effects to become important, since molecular viscosity dominates the flow in that region. Farther away from the wall, the mixing-length becomes a linear function of distance from the wall, following Prandtl's assumption for mixing-length behavior in the logarithmic region.

The van Driest parameter A^{\dagger} may be interpreted as a measure of the sublayer thickness. Based on this interpretation, several authors have represented A^{\dagger} as a function of pressure gradient and blowing. These parameters are known to affect the thickness of the wall region of the boundary layer. Andersen [8] gives a brief discussion of the role of A^{\dagger} and reviews some interpretations which have been given to it.

In the outer region of the boundary layer, the mixing length is held at a constant fraction of the 99% boundary layer thickness.

Values of λ = 0.08 to 0.09 have been demonstrated to be satisfactory for smooth wall boundary layers by several authors [35,38,50]. In these predictions, λ = 0.08 was used.

It was suggested by van Driest [39] that roughness effects could be simulated by changing the damping expression, since the presence of roughness elements in the flow seemed to reduce the sublayer thickness. This can be simulated in the van Driest mixing-length model by reducing A. In the present program, the relationship between A and roughness was investigated by systematically predicting pipe flows over a series of Reynolds numbers. It was felt that the relative simplicity of the pipe geometry would provide results which could then be generalized to the boundary layer case. Pipe flow predictions using different damping constants produced results which verified that roughness effects could indeed be simulated by making A a function of the roughness Reynolds number. Fig. 4.1 shows the functional forms of A which were determined for sand-grain roughness, commercial roughness, and the rough surface used in the present experiments. The A correlation for sandgrain roughness was obtained using Nikuradse's expression [9] relating friction factor to the roughness Reynolds number. The commercial roughness A correlation was based on relative roughness read directly from a Moody chart at friction factors calculated for various pipe Reynolds numbers. The A correlation for the rough surface used in these experiments was based on the friction factor measurements made at the 32 ft/sec tunnel velocity. This was the only tunnel velocity at which the roughness Reynolds number, based on an equivalent sand-grain roughness, was below 55. It is not surprising that the A behavior of the uniformball roughness used in these experiments is different than either commercial or sand-grain roughness. Roughness elements in a commercially rough surface are irregular and may include some very large elements which can become 'active' and begin to produce roughness effects at relatively low values of the nominal roughness Reynolds number. The roughness elements in the sand-grain roughened surface are more uniform, and the onset of

'transitional roughness' is delayed to higher roughness Reynolds numbers. The extremely uniform nature of the roughness used in these experiments still further delayed the onset of transitional roughness. This is the behavior shown in the A⁺ correlations for the different surfaces. As would be expected, the different A⁺ correlations come together as the roughness Reynolds number approaches the 'fully rough' value and at a roughness Reynolds number of 55, A⁺ has reached zero for all surfaces. This is consistent with the idea that at the onset of 'fully rough' behavior, the roughness elements protrude outside the sublayer. The sublayer thickness has effectively gone to zero when A⁺ has gone to zero. Without a sublayer, form drag replaces viscous shear and the flow resistance becomes independent of viscosity -- an important characteristic of 'fully rough' flow.

In the prediction program, the turbulent shear transmitted to the wall is calculated from an eddy diffusivity which is based on the mixing-length model. One way to represent the 'fully rough' behavior at the wall is to make the mixing-length and, therefore, the eddy diffusivity non-zero at the wall. There is evidence that the rough wall affects mainly the flow very near the wall. Clauser [48] has shown that velocity profiles taken over both smooth and rough surfaces are similar when plotted in defect coordinates. His conclusion was that the outer flow (which is some 90% of the boundary layer) is not affected by the surface roughness. In these coordinates both smooth and rough wall velocity profiles are similar. The prediction scheme of the present work is based on these observations.

The mixing-length was made non-zero at the wall, becoming asymptotic to the smooth wall behavior away from the wall by expressing the mixing-length in the following manner:

$$\ell_{\rm m}^+ = \sqrt{(\kappa y^+)^2 + (\Delta \ell_{\rm o}^+)^2}$$
 (4-3)

By again predicting pipe flows, a correlation of the value of the mixing length at the wall, Δl_0^+ , versus roughness Reynolds number was obtained.

$$\Delta l_0^+ = \sqrt{\frac{Re_{\tau}^{-46}}{39}^2 - .05325}$$
 (4-4)

Figure 4.2 shows the variation of Δl_0^+ with roughness Reynolds number along with the data predictions used to obtain Eqn. (4-4). It is interesting to note that for large values of roughness Reynolds number (Re_{τ} > 100) the functional relationship reduces to a simple linear variation between mixing-length offset, Δl_0 , and equivalent sand-grain roughness. The effect of roughness on mixing-length is related only to the size of the roughness elements, an important characteristic of 'fully rough' behavior. With this correlation, it is possible to predict the behavior of all of the surface roughnesses which require different A^+ correlations. This was expected, since in fully developed roughness the behavior of many different surfaces can be expressed in terms of an equivalent sand-grain roughness.

The mixing-length variation as a function of roughness Reynolds number is shown in Fig. 4.3. Fig. 4.4 shows predictions of Nikuradse's [9] rough pipe data. These predictions were made by modeling the Nikuradse experiment with the prediction program using the mixing-length model just described. The sand-grain roughness A^+ correlation was used for roughness Reynolds numbers below 55 and the correlation for mixing-length value at the wall, Δk_0^+ , above 55. The rough pipe behavior was well predicted over a wide range of roughnesses and pipe Reynolds numbers.

Predictions were made of the rough surface boundary layer from the present experiments using these same correlations. For roughness Reynolds numbers below 55, it was necessary to use an A⁺ correlation selected for this particular type of roughness. For roughness Reynolds numbers greater than 55, the same mixing-length offset correlation already described was used. Figs. 4.5 and 4.6 show a comparison of measured and predicted velocity profiles at velocities of 32 fps and 139 fps. In general, measured and predicted profiles in wall coordinates compare favorably as long as the measured and predicted skin frictions are comparable. This is generally the case over the last two-thirds of the test section. Fig. 4.7 compares the measured and predicted skin friction for

three test section velocities. In rearly every case the predicted values are greater than measured near the front of the test section. This same trend has also been observed in the smooth wall predictions and probably results because initially the boundary layer is still recoveying from transition and attempting to reach an equilibrium state.

In making these predictions, it was found that blowing decreased the predicted skin friction. In order to keep this reduced skin friction from reducing the effect of the roughness, a first-order blowing correction has been added to the correlation for the mixing-length at the wall, $\Delta \ell_0^+$. With the blowing correction, it becomes:

$$\Delta \ell_{o}^{+} = \sqrt{\frac{\text{Re}_{\tau} \left\{ 1 + f(v_{o}^{+}) \right\} - 46}{39}}^{2} - .05325$$

$$f = \begin{cases} \text{Re}_{\tau}/5 & \text{for } \text{Re}_{\tau} \leq 80 , \\ \\ 16 & \text{for } \text{Re}_{\tau} > 80 . \end{cases}$$
(4-5)

One would expect that blowing should play the same role in the mixing-length formulation no matter if the wall is rough or smooth. The expression for roughness Reynolds number with its blowing correction can be rewritten as the sum of two roughness related Reynolds numbers:

$$\operatorname{Re}_{\tau}\left\{1+f(v_{o}^{+})\right\} = \left(\frac{v_{\tau}k_{s}}{v}\right)+f\left(\frac{v_{o}k_{s}}{v}\right) . \tag{4-6}$$

It is reasonable that if there is blowing at the wall this should affect the mixing-length correlation and that the effect be similar to the correction that other investigators [8,35,49,50] have applied to A^{+} to ensure correct prediction of smooth surface skin friction with blowing.

B. Heat Transfer Predictions for the Rough Surface Boundary Layer

Heat transfer predictions were made by solving the boundary layer energy equation using a turbulent Prandtl number model for closure. A constant turbulent Prandtl number is often used for smooth surface boundary layer predictions. There is not complete agreement among experimenters about the behavior of the turbulent Prandtl number with distance from the wall, but most agree that it increases near the wall. One turbulent Prandtl number expression which has been used successfully for air [38] and which reflects this trend is given by:

$$Pr_{t,smooth} = (1.43 - 0.17 y^{+1/4})$$
 (4-7)

for $Pr_{t,smooth} \ge 0.86$. If $Pr_{t,smooth} < 0.86$, then

A plot of this model is shown in Fig. 4.8, and it can be noted that near the wall the Prandtl number increases to 1.43, which is approximately the reciprocal of the molecular Prandtl number for air. Also shown is the constant turbulent Prandtl number model often used for air as well as other, higher Prandtl number fluids.

Figure 4.9 illustrates the results obtained using the smooth surface turbulent Prandtl number model for rough pipe heat transfer to air. The predictions are plotted in the coordinates used by Nunner [23] to organize his rough pipe heat transfer data; rough-to-smooth Stanton number versus rough-to-smooth skin friction. While the predictions are reasonably well organized in these coordinates, probably due to the correct modeling of the skin friction, the general trend of the data is high. Use of this smooth wall turbulent Prandtl number over-predicts the heat transfer. This was not unexpected. The smooth surface turbulent Prandtl number model forces the momentum and thermal eddy diffusivities, and therefore the transport process, themselves, to be similar. In solving the momentum problem, the mixing-length has been increased near the wall to account for form drag and increased mixing at and near the surface due

to the presence of roughness elements. In the heat-transfer problem, the increased mixing will increase the thermal transport across the boundary layer near the wall, but at the wall there is no heat transfer mechanism analogous to the pressure force transmitted to the wall as form drag. At the wall, heat can be transferred only by conduction. To account for this difference using only the parameters available to the prediction program, the turbulent Prandtl number must be increased both at and near the wall. To accomplish this, a thermal mixing-length has been defined. The rough surface turbulent Prandtl number is then expressed as:

$$Pr_{t,rough} = Pr_{t,smooth} \left(\frac{\ell m}{\ell_H}\right)^2$$
 (4-8)

The thermal mixing-length is defined with the same modified van Driest model used for the momentum mixing-length and uses the same roughness correlations. However, when evaluating the thermal mixing-length, a reduced roughness Reynolds number is used. Since the roughness effects are built into the models only near the wall, the thermal mixing-length is less than the momentum mixing-length near the wall, but both approach smooth surface behavior away from the wall. This produces an increased turbulent Prandtl number near the wall, but approaching the smooth surface correlation away from the wall. It was decided to first apply this model to rough pipe heat transfer data and then generalize these results to the boundary layer case. The Dipprey and Sabersky [24] pipe data were chosen because of the wide range of roughness and fluid Prandtl numbers. It soon became evident when predicting these data that the thermal mixing-length correlation, to be successful, must be a function of Prandtl number as well as roughness Reynolds number. At high roughness Reynolds numbers, the increased mixing-length used for the hydrodynamic predictions almost entirely masked any molecular Prandtl number effects on heat transfer. To correctly predict the heat transfer results, the amount by which the roughness Reynolds number is reduced when evaluating the thermal mixing-length must be a function of molecular Prandtl number. Predictions of the Dipprey and Sabersky data are shown

in Figs. 4.10 and 4.11. It's apparent that the Prandtl number effects persist even in the very rough data. In making these predictions, the roughness Reynolds numbers were reduced by the factor $C_{\rm H}$, which was given by:

$$C_{H} = 0.5$$
 for $Re_{\tau} \le 40$, (4-9)
 $C_{H} = 0.5 - .0454(1.9+Pr)(\log Re_{\tau}-1.602)$.

In using this correlation, C_H is never allowed to become less than 0.1. A plot of C_H is shown in Fig. 4.12. It should be pointed out that C_H simply provides a convenient means by which to reduce the effects of roughness on the thermal mixing-length at and near the wall and therefore produce a higher turbulent Prandtl number at and near the wall.

To predict the heat transfer data from the present tests, it is sufficient to simply reduce the roughness Reynolds number by half. As can be seen from Fig. 4.12, over the limited roughness Reynolds number range of these experiments, this is essentially the same $C_{\rm H}$ which would be produced by the correlation used for prediction of the Dipprey and Sabersky data. Fig. 4.13 compares the Stanton number predictions for the present experiments with the data. As can be seen, the agreement of the predictions with the data is very good.

Finally, it should be pointed out that the models used here are based primarily on the observation that the roughness effect is confined to a region very near the wall, with the outer region of the boundary layer being relatively unaffected. In the present heat transfer model, if the reduced roughness correction were not used, the increased turbulent transport at the wall would overshadow any molecular effects. This has been shown experimentally not to be the case — the Prandtl number of the fluid plays an important role in rough surface heat transfer. As a result, a Prandtl number effect must be included in the correlations used for the Stanton number predictions. A possible reason why molecular effects persist in rough surface heat transfer lies in the fact that conduction is the only mechanism available for the exchange at the wall.

As the height of the roughness elements increases it is likely that the inner regions can be modeled as a nearly stationary fluid film. Heat must be transferred through this film by conduction, hence the molecular Prandtl number becomes important. This fluid film is not part of the present finite-difference prediction program. An improved modeling might include a thermal conduction layer in the energy equation to better represent the molecular effect that is present in the rough surface heat transfer data.

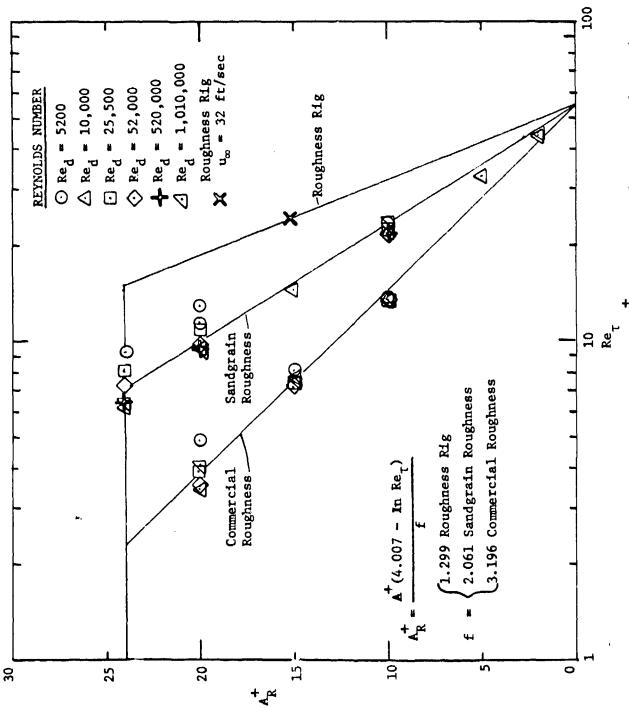
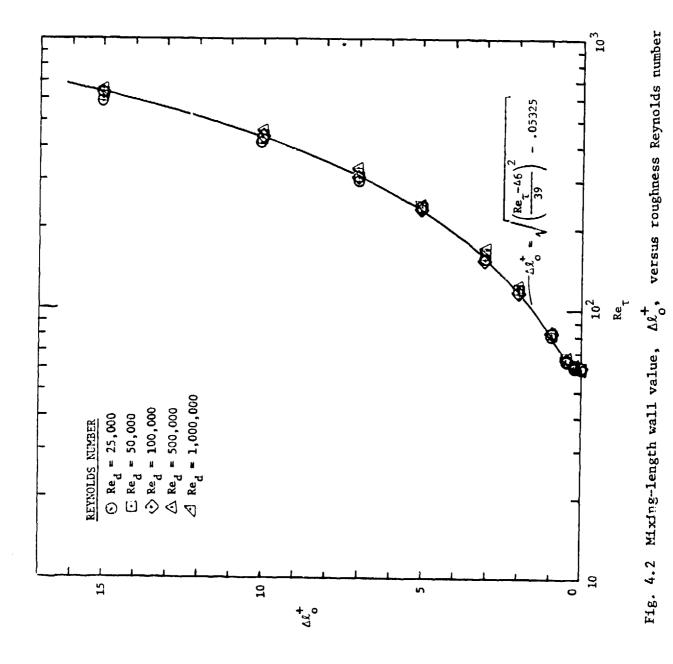
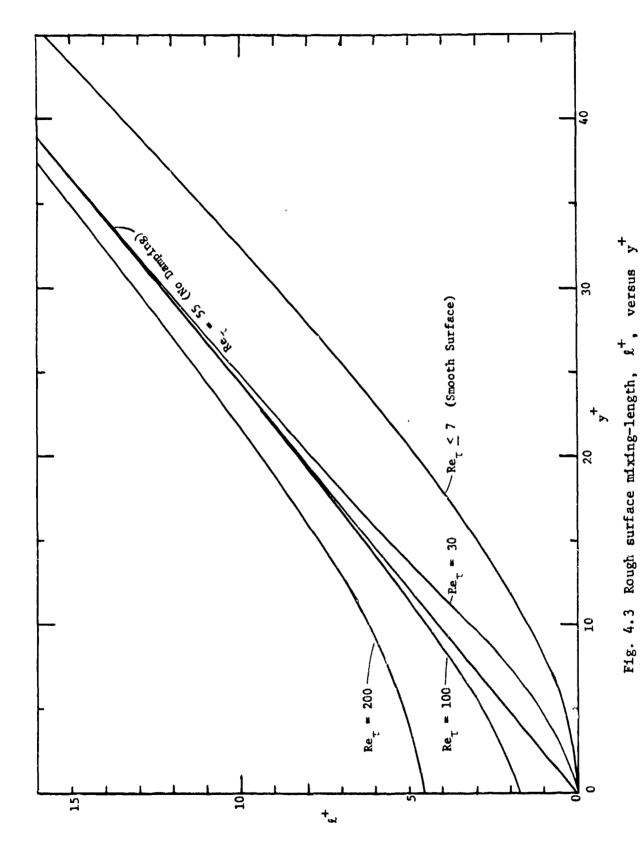


Fig. 4.1 Rough surface damping constant, A, versus roughness Reynolds number





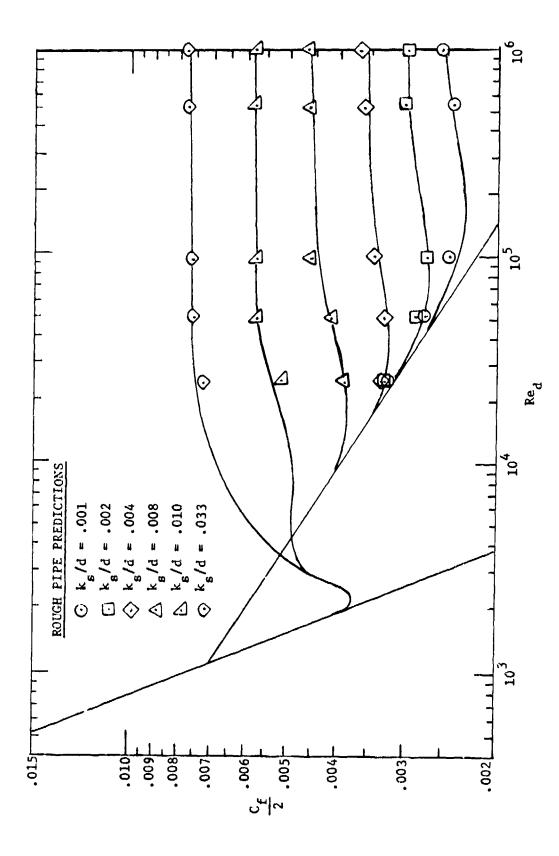
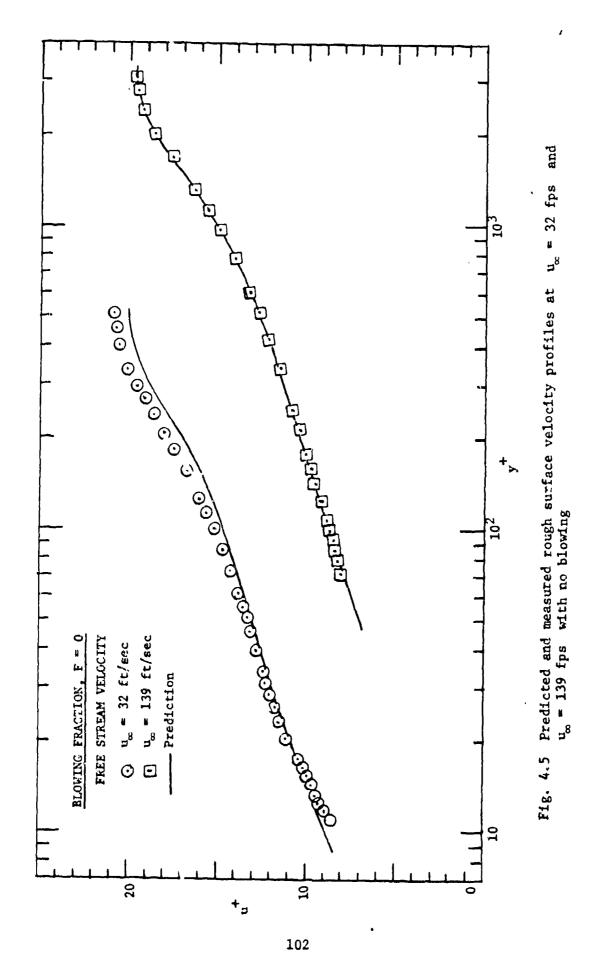
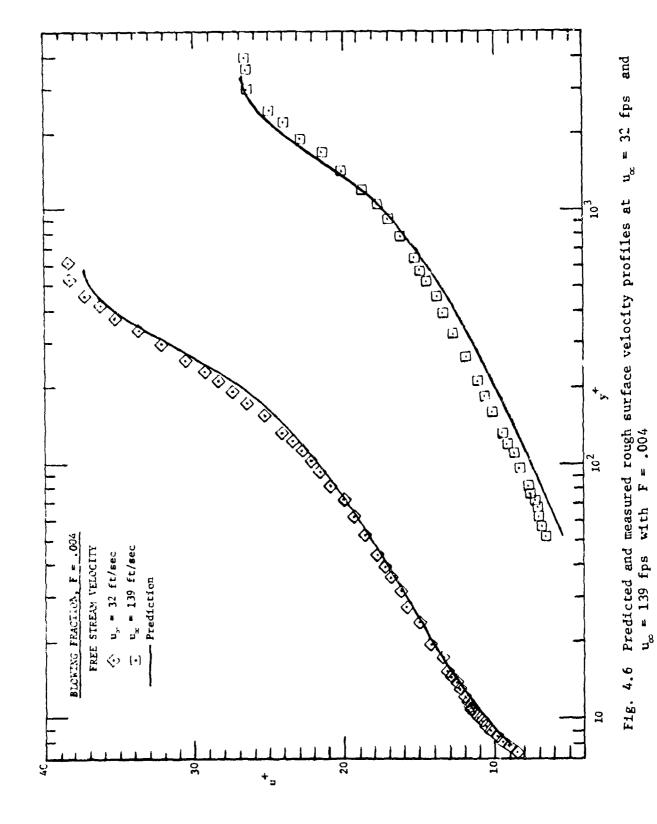


Fig. 4.4 Prediction of rough pipe friction factor data





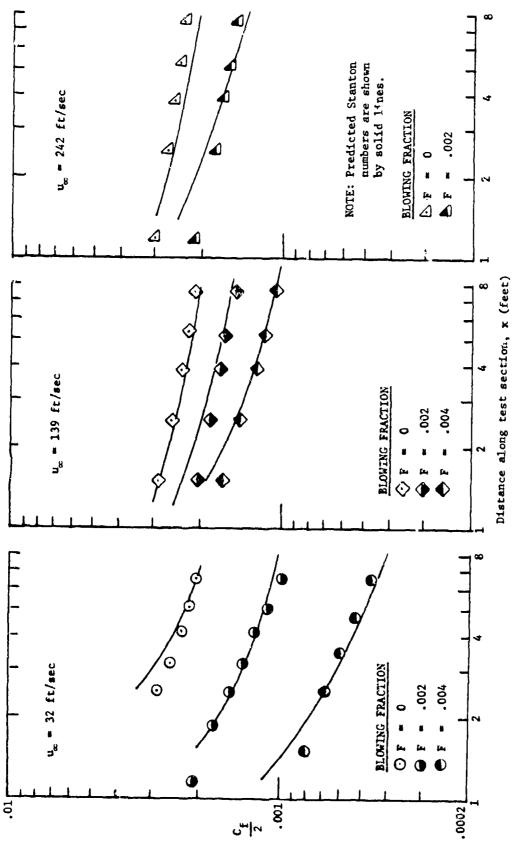


Fig. 4.7 Prediction of rough surface skin friction data

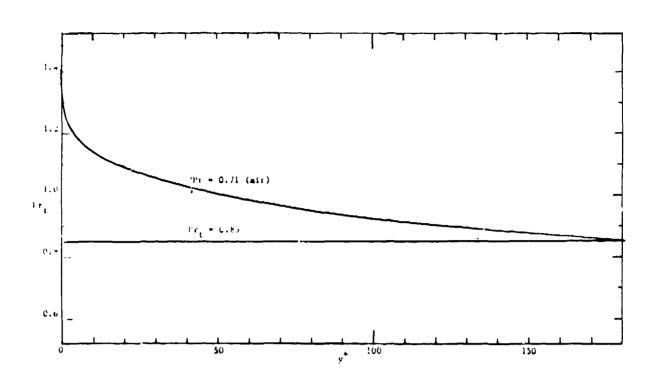


Fig. 4.8 Smooth surface turbulent Prandtl number model

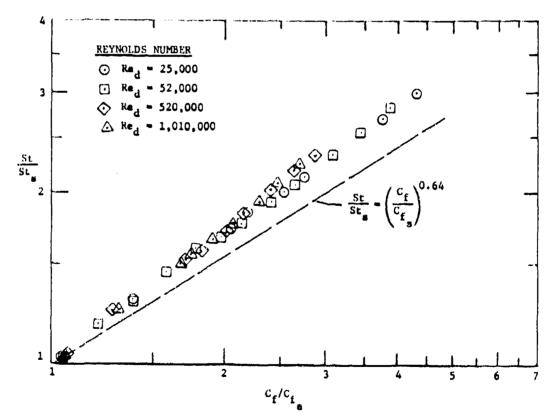


Fig. 4.9 Rough pipe heat transfer predictions with smooth surface turbulent Prandtl number model

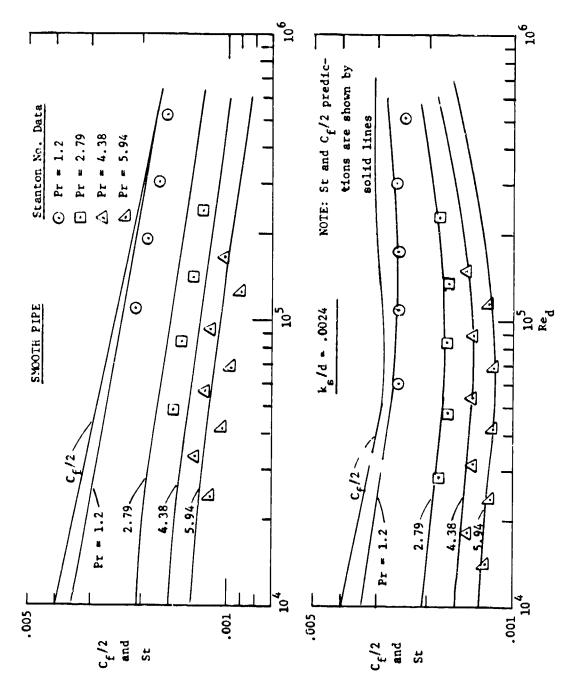


Fig. 4.10 Prediction of Dipprey and Sabersky [24] data for smooth tube and $k_{\rm s}/{\rm d}$ = .0024

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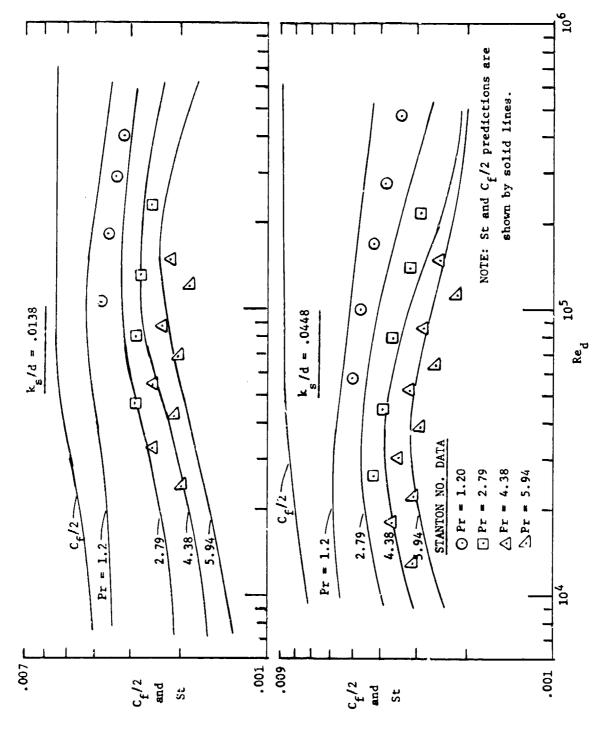


Fig. 4.11 Prediction of Dipprey and Sabersky [24] data for k_s/d = .0138 and k_s/d = .0448

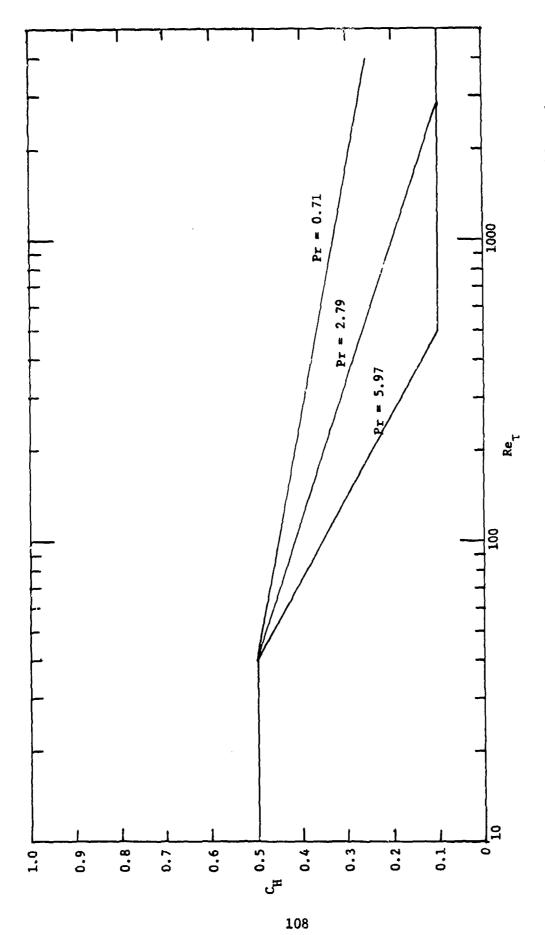


Fig. 4.12 Rough surface thermal mixing-length factor, C_{H} , versus roughness Reynolds number

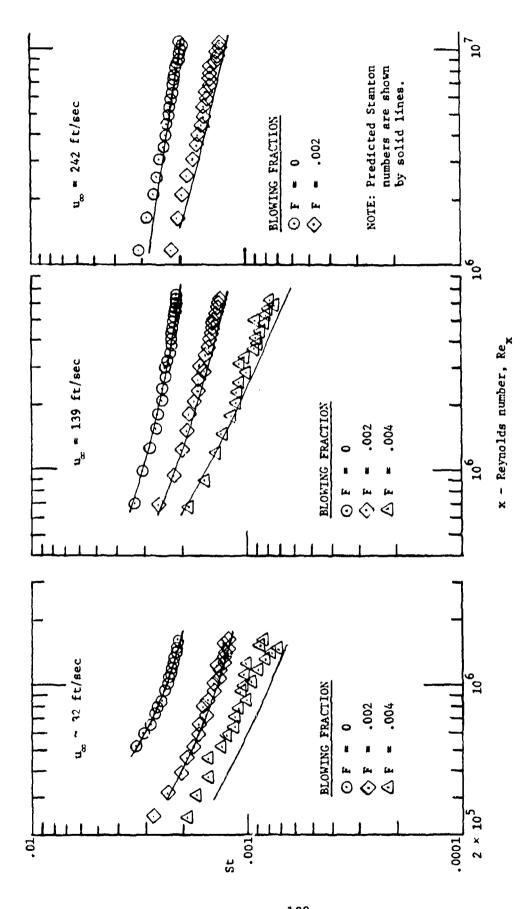


Fig. 4.13 Prediction of rough surface Stanton number data with blowing

CHAPTER V

SUMMARY AND CONCLUSIONS

The principal results and conclusions of the present rough surface experiments can be summarized as follows:

- 1. The present study has used a deterministic roughness, 0.050 inch diameter spherical elements arranged into a most-dense array. The thermal and hydrodynamic performance of the boundary layer on this surface was quite likely affected by its highly regular roughness. In particular, the marked absence of any 'transitional' behavior may be due, in part, to this effect.
- 2. The range of roughness Reynolds numbers $u_{\tau}k_{g}/v$, for the noblowing tests was from 24 to 200, covering the transitional roughness range and extending well into the fully rough range by conventional definitions. The fully rough state seems to have been attained: it certainly was for the heat transfer data and probably was for the friction factor data. The measured Stanton numbers are not dependent on free-stream velocity, but are functions only of the enthalpy thickness of the boundary layer and the blowing fraction, $\rho_{o}v_{o}/\rho_{o}u_{o}$. The unblown Stanton number data for all five velocities tested lie within a few percent of each other when compared at the same enthalpy thickness. With blowing, the uncertainty in the measurements increases but the tendency to be independent of velocity remains, even at very high blowing.

These comments seem also to apply to the friction factor behavior, but with less certainty, since the data are more sensitive to errors and show more scatter.

3. The behavior of the boundary layer at low roughness Reynolds numbers was investigated by an additional test with the boundary layer thickened by blowing through the first two feet of the eight foot test section. This action reduced the roughness Reynolds number to 14 at the downstream end of the plate. Even at this low roughness Reynolds numbers the heat transfer data continued to exhibit 'fully rough' behavior: Stanton number was the same function of enthalpy thickness as observed

at higher roughness Reynolds numbers. It remained independent of stream velocity.

The expected approach to smooth plate behavior was not observed. If smooth plate heat transfer represents the asymptotic behavior of the rough surface as the roughness Reynolds number is reduced, it was not apparent from these experiments.

4. Blowing through a rough surface diminishes the Stanton number. The effect of blowing is predictable using the same relationship found valid in smooth plate studies:

$$\frac{\operatorname{St}}{\operatorname{St}_{o}} = \left[\frac{\ln(1+B)}{B}\right]^{1.25} (1+B)^{.25} , \qquad (5-1)$$

where St and St are evaluated at the same enthalpy thickness, Δ (for a smooth plate, the Stanton numbers are to be evaluated at the same enthalpy thickness Reynolds number).

- 5. Transition on the rough surface began at approximately the same momentum thickness Reynolds number (350-450) for all of the conditions tested. For the surface tested here, transition occurs at about the same momentum thickness Reynolds number that would be expected for transition on a smooth plate. The transition moves upstream with increasing free-stream velocity and with blowing.
- 6. Blowing in a region which would otherwise have remained laminar may greatly increase the local heat load. The turbulent heat transfer coefficient, even with moderately strong blowing, is much larger than the laminar, unblown value.
- 7. The effects of roughness on skin friction can be accounted for in a finite-difference boundary layer prediction scheme, using a mixing-length model for mean field closure, by modifying the mixing-length variation very near the wall. Successful prediction of rough surface heat transfer requires modification of the turbulent Prandtl number distribution very near the wall.

In conclusion, it is safe to say that the work reported here is just the beginning of the rough surface investigations at Stanford. Already under way is an extensive program to make detailed hydrodynamic measurements in the rough surface boundary layer. When complete, these should provide a much better basis for formulating realistic prediction models.

REFERENCES

- 1. Moffat, R.J., and Kays, W.M., "The Turbulent Boundary Layer on a Porous Plate: Experimental Heat Transfer with Uniform Blowing and Suction," Report No. HMT-1, Thermosciences Division, Dept. of Mech. Engrg., Stanford Univ., 1967.
- 2. Simpson, R.L., Kays, W.M., and Moffat, R.J., "The Turbulent Boundary Layer on a Porous Plate: An Experimental Study of the Fluid Dynamics with Injection and Suction," Report No. HMT-2, Thermosciences Division, Dept. of Mech. Engrg., Stanford Univ., 1967.
- 3. Whitten, D.G., Kays, W.M., and Moffat, R.J., "The Turbulent Boundary Layer on a Porous Plate: Experimental Heat Transfer with Variable Suction, Blowing and Surface Temperature," Report No. HMT-3, Thermosciences Division, Dept. of Mech. Engrg., Stanford Univ., 1967.
- 4. Julien, H.L., Kays, W.M., and Moffat, R.J., "The Turbulent Boundary Layer on a Porous Plate: Experimental Study of the Effects of a Favorable Pressure Gradient," Report No. HMT-4, Thermosciences Division, Dept. of Mech. Engrg., Stanford Univ., 1969.
- 5. Thielbahr, W.H., Kays, W.M., and Moffat, R.J., "The Turbulent Boundary Layer: Experimental Heat Transfer with Blowing, Suction, and Favorable Pressure Gradient," Report No. HMT-5, Thermosciences Division, Dept. of Mech. Engrg., Stanford Univ., 1969.
- 6. Kearney, D.W., Moffat, R.J., and Kays, W.M., "The Turbulent Boundary Layer: Experimental Keat Transfer with Strong Favorable Pressure Gradients and Blowing," Report No. HMT-12, Thermosciences Division, Dept. of Mech. Engrg., Stanford Univ., 1970.
- 7. Loyd, R.J., Moffat, R.J., and Kays, W.M., "The Turbulent Boundary Layer on a Porous Plate: An Experimental Study of the Fluid Dynamics with Strong Favorable Pressure Gradients and Blowing," Report No. HMT-13, Thermosciences Division, Dept. of Mech. Engrg., Stanford Univ., 1970.
- 8. Andersen, P.S., Kays, W.M., and Moffat, R.J., "The Turbulent Boundary Layer on a Porous Plate: An Experimental Study of the Fluid Mechanics for Adverse Free-Stream Pressure Gradients," Report No. HMT-15, Thermosciences Division, Dept. of Mech. Engrg., Stanford Univ., 1972.
- 9. Nikuradse, J. "Stromungsgestze in rauhen Rohren," VDI Forschungsheft, No. 361, 1950, English Translation, NACA TM 1292.
- 10. Prandtl, L., and Schlichting, H. "Das Widerstandsgesetz rauher Platten," Werft, Reederei, und Hafen, 1934, p. 1.

- 11. Karman, Th. von, "Turbulence and Skin Friction," Jn. Aero. Sciences, Vol. 1, No. 1, January 1934.
- 12. Moore, W.L., "An Experimental Investigation of Boundary Layer Development Along A Rough Surface," PhD Dissertation, State University of Iowa, August 1951.
- 13. Hama, F.R., "Boundary Layer Characteristics for Smooth and Rough Surfaces," Trans. SNAME, Vol. 62, pp. 333-354, 1954.
- 14. Perry, A.E., Schofield, W.H., and Joubert, P.H., "Rough Wall Turbulent Boundary Layers," Jn. Fl. Mech., Vol. 37, pp. 383-413, 1969.
- 15. Perry, A.E. and Joubert, P.H., "Rough-wall Boundary Layers in Adverse Pressure Gradients," Jn. Fl. Mech., Vol. 17, pp. 193-211, 1963.
- 16. Liu, C.K., Kline, S.J., and Johnston, J.P., "An Experimental Study of Turbulent Boundary Layer on Rough Walls," Report No. MD-15, Thermosciences Division, Dept. of Mech. Engrg., Stanford Univ., July 1966.
- 17. Grass, A.J., "Structural Features of Turbulent Flow Over Smooth and Rough Boundaries," Jn. Fl. Mech., Vol. 50, pp. 233-255, 1971.
- 18. Wu, J., "Flow in Turbulent Wall Layer Over Uniform Roughness," ASME paper 73-APM-U, to be published in Jn. Appl. Mech., Trans. ASME.
- 19. Tsuji, Yutaka, and Iida, Shusuke, "Velocity Distributions of Rough Wall Turbulent Boundary Layers Without Pressure Gradient," <u>Trans.</u>
 <u>Japan Soc. Aero. Space Sci.</u>, Vol. 16, No. 31, pp. 60-70, 1973.
- 20. Antonio, R.A. and Luxton, R.T., "The Response of a Turbulent Boundary Layer to a Step Change in Surface Roughness, Part I: Smooth to Rough," Ju Fl. Mech., Vol. 48, Part 4, pp. 721-761, 1971.
- 21. Townes, H.W., Gow, J.L., Powe, R.E. and Weber, N., "Turbulent Flow in Smooth and Rough Pipes," Jn. Basic Eng., Trans. ASME, Series D, Vol. 94, No. 2, pp. 353-362, June 1972.
- 22. Townes, H.W. and Powe, R.E., "Turbulence Structure for Fully Developed Flow in Rough Pipes," Jn. Fl. Engrg., Trans. ASME, Series I, Vol. 95, No. 2, pp. 255-261, July 1973.
- 23. Nunner, W., "Heat Transfer and Pressure Drop in Rough Tubes," VDI Forschungsheft, No. 455, Series B, Vol. 22, pp. 5-39, 1956, English Translation, A.E.R.E. Library/Transactions 786, 1958.

- 24. Dipprey, D.F., and Sabersky, R.H., "Heat and Momentum Transfer in Smooth and Rough Tubes at Various Prandtl Numbers," <u>Intnl. Jn.</u> of Heat Mass Transfer, Vol. 6, pp. 329-353, 1963.
- 25. Owen, P.R. and Thomson, W.R., "Heat Transfer Across Rough Surfaces,"

 Jn. Fl. Mech., Vol. 15, pp. 321-334, 1963.
- 26. Gowen, R.A. and Smith, J.W., "Turbulent Heat Transfer from Smooth and Rough Surfaces," <u>Intnl. Jn. of Heat Mass Transfer</u>, V- 11 pp. 1657-1673, 1968.
- 27. Reshotko, M., "Roughness Effects on Heat Transfer in the Supersonic Region of a Conical Nozzle," Jn. Spacecraft and Rockets, Vol. 8, No. 10, pp. 1097-1099, October 1971.
- 28. Boldman, D.R. and Graham, R.W., "Heat Transfer and Boundary Layer in Conical Nozzles," NASA TN D-6594, February 1972.
- 29. Sood, N.S. and Jonsson, V.K., "Some Correlations for Resistances to Heat and Momentum Transfer in the Viscous Sublayer at Rough Walls," <u>Jn. Heat Transfer</u>, Trans. ASME, pp. 488-494, November 1969.
- 30. Norris, R.H., "Some Simple Approximate Heat Transfer Correlations for Turbulent Flow in Ducts with Surface Roughness," <u>Augmentation of Convective Heat and Mass Transfer</u>, published by the ASME, 1971.
- 31. Chen, Karl, K., "Compressible Turbulent Boundary-Layer Heat Transfer to Rough Surfaces in Pressure Gradient, "AIAA Journal, Vol. 10, No. 5, pp. 623-629, May 1972.
- 32. Dvorak, F. A., "Calculation of Turbulent Boundary Layers on Rough Surfaces in Pressure Gradient, "AIAA Journal, Vol. 7, No. 9, pp. 1752-1759, Sept. 1969.
- 33. Dvorak, R. A., "Calculations of Compressible Turbulent Boundary Layers with Roughness and Heat Transfer, " ALAA Journal, Vol. 10, No. 11, pp. 1447-1451, November 1972.
- 34. Nestler, D. E., "Compressible Turbulent Boundary-Layer Heat Transfer to Rough Surfaces," AIAA Journal, Vol. 9, No. 9, Sept. 1971.
- 35. Lumsdaine, E., Wen, H. W., and King, F. K., "Influence of Surface Roughness and Mass Transfer on Boundary Layer and Friction Coefficient," <u>Developments in Mechanics</u>, Vol. 6, Proceedings of the 12th Midwestern Mechanics Conference, pp. 305-318.
- 36. McDonald, H., and Fish, R. W., "Practical Calculations of Transitional Boundary Layers," <u>International Journal of Heat Mass Transfer</u>, Vol. 16, pp. 1729-1744, 1973.

- 37. Spalding, D. B. and Pantanker, S. V., <u>Heat and Mass Transfer in Boundary Layers</u>, Morgan-Granpian, London, 1967.
- 38. Kays, W. M., "Heat Transfer to the Transpired Turbulent Boundary Layer," Report No. HMT-14, Thermosciences Div., Dept. of Mech. Engrg, Stanford University, Stanford, California, June 1971.
- 39. van Driest, E. F., "On Turbulent Flow Near a Wall," Heat Transfer and Fluid Mechanics Institute, 1955.
- 40. Wolf, S., "Flow Losses for Heat Exchangers with Oblique Flow Headers," Technical Report No. 60, Department of Mechanical Engrg., Stanford Univ., June 1964.
- 41. Schubauer, G. E., Spangenberg, W. G., and Klebanoff, P. S., "Aerodynamic Characteristics of Damping Screens," NACA TN 2001, January 1950.
- 42. Rause, H., and Hassan, M. M., "Cavitation Free Inlets and Contractions." Mechanical Engineering, 71, 3, 213, 216 1949.
- 43. Cocharan, D. L. and Kline, S. J., "The Use of Short Flat Vanes for Producing Efficient Wide-Angle Two-Dimensional Subsonic Diffusers," NACA TN 4309, September 1958.
- 44. Bradshaw, P., "The effect of Wind-Tunnel Screens in Nominally Two-Dimensional Boundary Layers," <u>Journal of Fluid Mechanics</u>, Vol. 22, Part 4, pp. 679-687, 1965.
- 45. Morgan, P. G., "The Stability of Flow Through Porous Screens," Journal of the Royal Aeronautical Society, Vol. 64, pp. 359-363, 1960.
- 46. Scarborough, James R., <u>Numerical Mathematical Analysis</u>, 6th Edition, Johns Hopkins Press, pp. 541-551, 1964.
- 47. Kline, S. J. and McClintock, F. A., "Describing Uncertainties in Single Sample Experiments," Mechanical Engineering, Jan. 1953.
- 48. Clauser, F. H., "The Turbulent Boundary Layer, " Advances in Applied Mechanics, Vol. 1V, pp. 1-51, Academic Press, New York, 1956.
- 49. Cebeci, T., "Behavior of Turbulent Flow near a Porous Wall with Pressure Gradient," <u>AIAA Journal</u>, Vol. 8, No. 12, pp. 2152-2156, Dec. 1970.
- 50. Pletcher, R. H., "Prediction of Transpired Turbulent Boundary Layers," <u>Journal of Heat Transfer</u>, Trans. ASME, Series C, pp. 89-94, Feb. 1974.

- 51. Schlichting, H., <u>Boundary Layer Theory</u>, 6th Edition, McGraw-Hill Book Co., Inc., New York, 1968.
- 52. Lakshman, C. and Jayatilleke, V., "The Influence of Prandtl Number and Surface Roughness on the Resistance of the Laminar Sublayer to Momentum and Heat Transfer", <u>Progress in Heat and Mass Transfer</u>, Vol. 1, Pergamon Press, pp. 193-330, 1969.
- 53. Reynolds, W. C., Kays, W. M., and Kline, S. J., "Heat Transfer in the Turbulent Incompressible Boundary Layer, Part IV", NASA Memo 12-4-58W, 1958.

APPENDIX A

MAIN AIR SUPPLY SYSTEM COMPONENTS

The following is a brief description of the construction details of the major components in the main air supply system.

A.1 Inlet Header and Screen Box

The heat exchanger inlet header and screen box are shown with their overall dimension in Fig. A.1. One side of the box is removable to allow removal of the screens and screen spacers. The screen material is #40 mesh, .0065 stainless steel wire stretched on 1 x 4 inch clear pine frames. Four such frames are used with 2 x 4 inch clear pine spacer frames between each screen. Both sides of each screen and spacer frame are lined with felt gasket material to avoid leakage around the screen pack. Alignment pins are used to insure that the screen and spacer frames are properly aligned in the box. A photograph of the header, heat exchanger and screen box assembly before the nozzle was installed is shown in Fig. A.2.

A.2 Wind Tunnel Nozzle

The Roughness Rig nozzle shape is shown in Fig. A.3. The nozzle was fabricated in two halves from fiberglass and then hand fitted together. The nozzle wall shapes were first layed out on 1/16 inch aluminum templates from which the mold for the fiberglass build-up was made using high density, closed cell polyurethane foam. The approximate nozzle shape was carved from a block of polyurethane foam with a hot chromel wire using the aluminum templates as guides. The final shape was obtained by hand sanding the foam block and then reinforcing its surface with a thin layer of polyester resin which was sanded to conform to the final nozzle surface shape and surface smoothness.

Before the fiberglass layup was made, the mold was waxed and sprayed with a release agent. The first coat of resin was a white gel coat, sprayed on the mold. This was followed by two layers of fiberglass mat,

a layer of fiberglass cloth and two additional layers of fiberglass mat. An additional layer of mat was added to the nozzle flange for reinforcement.

Reinforcing ribs made from one-eighth inch thick masonite were attached to provide added strength and the nozzle halves were bolted together along the horizontal midplane. The final shape of the nozzle surface conformed to the design shape within .020 inches total run out. Figure 4.4 is a photograph of the two nozzle halves. The male nozzle mold is shown in the background.

A.3 Multistage Diffuser

The Roughness Rig diffuser is a multistage vaned diffuser unlike any previously built at Stanford in that several two dimensional diffusers are arranged in series. Figure A.5 shows a sketch of the diffuser with its overall dimensions. A photograph of the diffuser before it was installed in the rig is shown in Fig. A.6. The first stage of the diffuser has a movable top which can be adjusted to align itself with the test section top and form a smooth transition at the test section exit. Following the adjustable inlet section is a two-dimensional expansion which employs five vertical vanes to expand the diffuser width from 20 to 24 inches. Next are two separate two-dimensional expansions employing five and then eight vanes to expand the diffuser height to 24 inches. This empties into a plenum box connected to the blower inlet with a flexible connection. As can be noted from the photograph, the diffuser variable area inlet section and first vaned expansion has been constructed from plexiglass with aluminum vanes. At the bottom of the plenum box is a flange for connection to the small charging blower used to control the tunnel static pressure level. The estimated efficiency of this diffuser, had each of the stages performed in the assembly as it would have individually, would have been in excess of 50%. The actual diffuser efficiency is closer to 40% and even less at lower test section velocities. This performance is however acceptable and the wind tunnel does achieve its test section design velocities. Part of the difficulty with the diffuser performance is that the inlet flow into the diffuser is almost all boundary layer fluid, particularly when there is blowing in the test section. This contributes to its reduced performance.

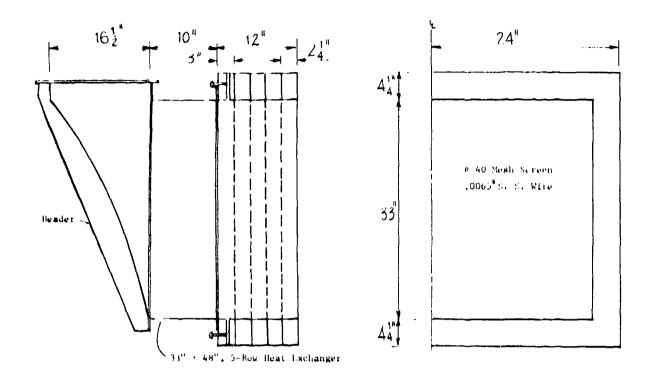


Fig. A.1 Heat exchanger oblique header and screen box

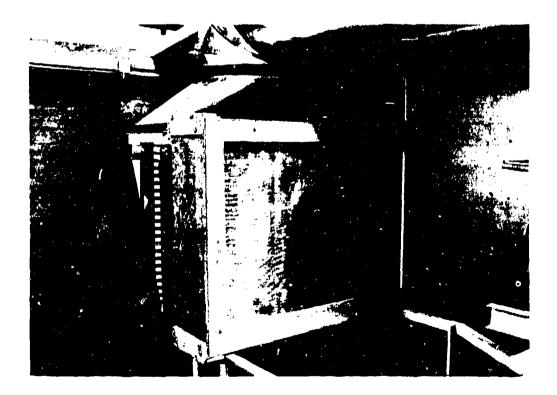


Fig. A.2 Photograph of the screen box, heat exchanger and header assembly

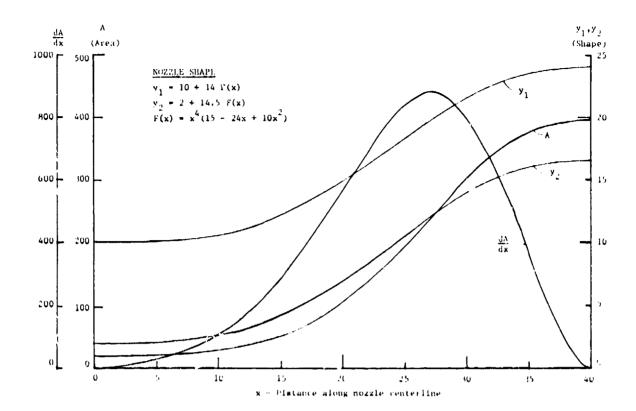


Fig. A.3 Nozzle wall shape



Fig. A.4 Photograph of the nozzle halves and mold before assembly

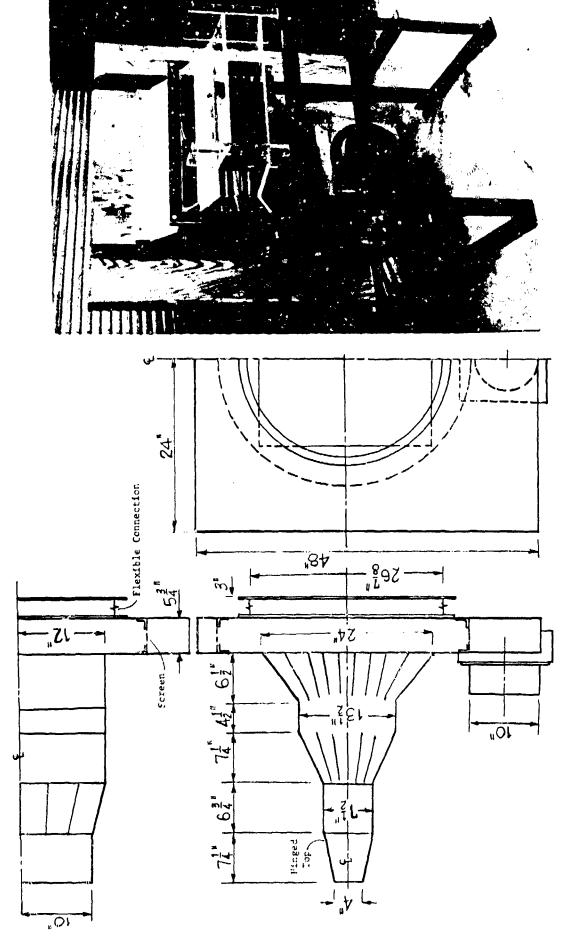


Fig. A.6 Photograph of the diffuser before installation in the wind tunnel Fig. 4.5 Multistage, vaned diffuser

APPENDIX B

TRANSPIRATION FLOW METERING SYSTEM

Constant current, hot wire type flow meters are used to measure flow rate in each of the individual plate transpiration air supply lines. These flow meters consist of two side-by-side probes installed in the center of a 3 foot long tube used as a riser between the transpiration header box and the control valves. Each unit is individually calibrated. The probes are small tubes which span the delivery pipe and contain the junctions of a differential iron-constantan thermocouple. One junction is in a thin wall, 0.025 inch diameter glass tube which is a passive probe. The other junction is in the center of a heated rod. The current in the heater is held constant during flow meter operation and the differential thermocouple output is the signal used to measure flow rate in the tube. A schematic of the circuit is shown in Fig. B.1. The heated probe consists of a #2-56 nylon screw, center-drilled with a #50 drill. Each end of the screw is threaded into a short piece of #8 brass screw stock, which serves as an end terminal for the heater power connections. Nichrome heater wire is wound onto the nylon screw using the threads to space the heater wire and is soldered at each end to the brass terminals. The heater wire is then potted in place with Permabond cement. A photograph showing the soldering of the heater wire to the brass end terminals on a flow meter probe is shown in Fig. B.2. Typical resistance of a probe heater is 5 ohms, although this value varies a few percent from heater to heater.

Each flow meter unit, consisting of the riser tube with differential thermocouple and heater, was individually calibrated on the flow bench in the Thermosciences Measurement Center. The meters were calibrated in both upward and downward flow, for blowing and suction, using a bank of Meriam Laminar Flow Meters as secondary standards. These laminar flow meters had been previously calibrated against standard ASME flow orifices. In the controlled environment of the Measurement center, it is felt that the flow measurements were accurate to within ± 1%, the accuracy normally claimed for the ASME orifice.

At each flow calibration point, the heater current was set and the differential thermocouple signal allowed to stabilize. A Hewlett-Packard 2401 IDVM using an external clock to allow integration of the flow meter signal over a 10-second interval was used to read the differential thermocouples. A typical flow meter calibration curve is shown in Fig. B.3.

The flow meters installed in the rig can operate in two modes. In the first mode, the power and differential thermocouple connections are made through a series of double pole switches such that flow meter operation is restricted to a single channel at a time. Activation of each successive flow meter heater circuit also activates its read-out circuit. The flow meters are operated in this mode while setting up the rigrunning conditions. In the second mode of operation, power is supplied to the meters in blocks of 6 at a time. This mode is used during datataking. A short computer program is used, called FLOMET, which calculates the flow meter signal requir d for the desired transpiration flow for each individual meter. The flow control valves must be adjusted, one at a time, such that each flo / meter produces its desired signal. The same flow meter calibration tables in FLOMET are also contained in the datareduction program, and the flow meter readings are re-recorded when the actual heat transfer deta are taken, so accurate flow measurement is assured.

The flow calibration tables contain flow and flow meter response in terms of SCFM and EMFO -- that is, cubic feet per minute of air at standard conditions and the emf reading the flow meter would have produced if air at standard conditions had been flowing through the transpiration system. To convert from actual to standard conditions, the following arguments must be used. The flow meter differential thermocouple signal is proportional to the temperature difference between the heated and the passive probes of the meter.

 $emf \propto \Delta T$ (B-1)

The temperature difference can be related to the heat transfer coefficient between the flow and the heated probe, and finally to the Nusselt number.

emf
$$\alpha$$
 $\Delta T = \frac{q}{Ah} = \left(\frac{qd}{Ak}\right)\frac{k}{hd} = \frac{qd}{Ak}\left(\frac{1}{Nu}\right)$ (B-2)

A similar expression can be written for a flow meter operating at standard conditions, denoted by the subscript 'c'.

$$\operatorname{emf}_{o} \propto \frac{\operatorname{q}_{o}^{d}_{o}}{\operatorname{A}_{o}^{k}_{o}} \left(\frac{1}{\operatorname{Nu}_{o}}\right)$$
 (B-3)

Taking the ratio of the two and noting that the area, diameter and heat flux will remain the same between the two conditions,

$$\frac{\text{emf}_{o}}{\text{emf}} = \frac{\text{Nu}}{\text{Nu}_{o}} \frac{k}{k_{o}}$$
 (B-4)

From the flow meter calibration data, the signal is seen to vary with the 0.45 power of the flow Reynolds number. Using the normal 0.3 dependence on Prandtl number, we can write

$$Nu \propto Re^{0.45} Pr^{0.3} = Re^{0.45} (k/\mu c)^{0.3}$$
 (B-5)

Also,

Re
$$\propto$$
 SCFM $\left(\mu^{-1}\right)$ (B-6)

With this, holding SCFM constant yields

$$\frac{\text{emf}}{\text{emf}} = \left(\frac{\text{SCFM}}{\text{SCFM}}\right)^{0.45} \left(\frac{\mu_o}{\mu}\right)^{0.45} \left(\frac{k_o}{\mu} \frac{\mu}{\mu_o} \frac{c}{c_o}\right)^{0.3} \left(\frac{k_o}{k_o}\right) = \left(\frac{\mu_o}{\mu}\right)^{0.15} \left(\frac{c}{c_o}\right)^{0.3} \left(\frac{k}{k_o}\right)^{0.7}$$
(B-7)

The following temperature and humidity dependence is assumed for the air properties:

$$k/k_0 = (T/T_0)^{0.74}$$

$$\mu/\mu_{Q} = (T/T_{Q})^{0.725}$$

 ω = absolute humidity of the air

(k/k₀) = not humidity dependent in the range of the test data

$$(\mu/\mu_0) = (1 - 0.7 \omega)$$

$$(c/c_0) = (1 + 0.9 \omega)$$
.

Substituting these into the expression for emf gives

$$\frac{\text{emf}_{o}}{\text{emf}} = \left(\frac{\text{T}}{\text{T}_{o}}\right)^{0.41} \frac{(1+0.9 \ \omega)^{0.3}}{(1-0.7 \ \omega)^{0.15}}$$
(B-8)

Using binomial expansions for the humidity terms, this can be simplified to

$$emf_o = emf(T/T_o)^{0.41} (1 + 0.38 \omega)$$
 (B-9)

All flow meter calibration data have been converted into standard conditions with this expression, for use in the data reduction program. Flow meter readings taken at test conditions are then converted to standard conditions before entering the calibration tables.

Experience with this metering system has been highly satisfactory. At relatively high flows, 10 cfm and over, adjustment of the valves to give the desired flow meter signal is relatively easy. At lower flow rates, flow valve adjustment is a bit more tedious because of the slower reciponse of the meter and the coarse adjustment that the ball valves afford in their nearly closed positions. The meters themselves seem stable and do not show any appreciable drift during a data run. On at least two occasions during the rig shakedown, flow meters were removed from the rig and recalibrated to be sure that the heaters and differential thermocouples were not 'aging' and changing with use. In only two cases

were changes in the meter calibrations noted, and each of those occurred because of overheating of the heater probe. The two probes had been inadvertently activated when there was no flow in the transpiration system. In all other cases, there had been virtually no change in the flow meter calibrations.

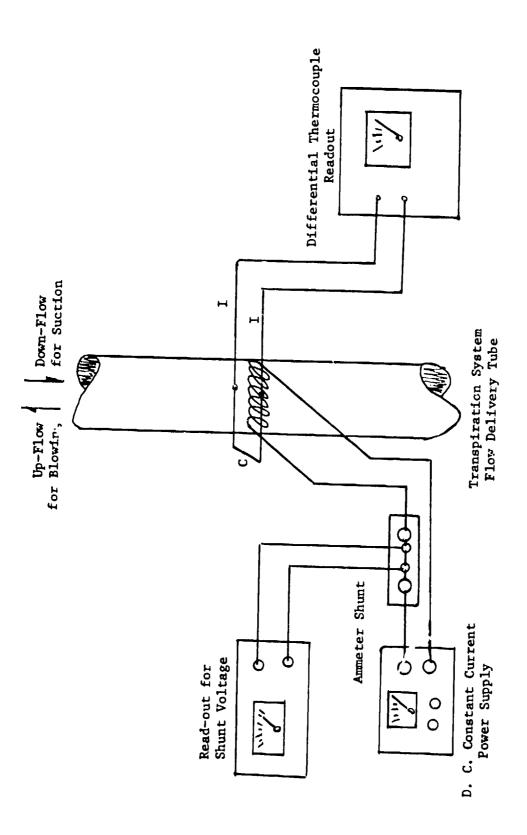


Fig. B.1 Schematic of the constant current, hot wire type flow meter

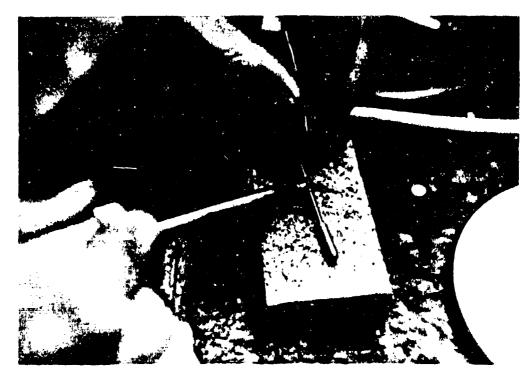


Fig. B.2 Photograph of the flow meter heater element

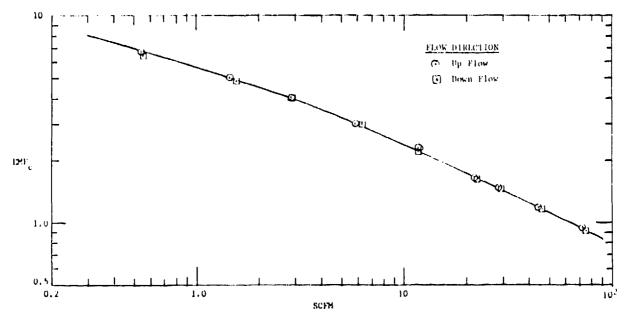


Fig. B.3 Typical flow meter calibration curve

APPENDIX C

PLATE POWER CONTROL CIRCUIT

Plate power control is achieved by controlling heater voltage using a solid-state control circuit for each plate, as shown in Fig. C.1. The circuit was designed by T. R. Mogel, of the Civil Engineering Department at Stanford, who also did the printed circuit board layout and acted as general consultant in shaking down the Roughness Rig power control system. The following description is paraphrased from a discussion of the control circuit written by T. R. Mogel for the seventh quarterly progress report on the research contract which funded the Roughness Rig project.

The power control circuit consists of three main parts: an operational amplifier, a high-power output amplifier, and a feedback network. The heater voltage is set by a 10 turn pot with $\pm 0.2\%$ linearity, which sends a proportional amount of the -10 volt reference to the heater controller. The voltage appearing across the heater is determined by the feedback network. The whole circuit can be thought of as an operational amplifier whose D.C. transfer characteristic will be:

$$V_{out} = -V_{in} R_1/R_3 \qquad (R_1 = R_2, R_3 = R_4)$$
 (C-1)

The output voltage will be

$$v_{\text{heater}} = -v_{\text{ref}} \alpha R_1/R_3$$
 (C-2)

For the resistances used in the circuit,

$$V_{\text{heater}} = 22.1 \alpha$$
 (C-3)

where α is 0 when the pot resistance is at the counterclockwise stop and when 1 is at the clockwise stop.

The capacitor is used to stabilize the amplifier by rolling off the high frequencies. A diode is used to prevent the output of the operational

amplifier from going negative. A second diode is used to allow the highest of the heater voltages from all the channels to be known.

A feature of this system is that the heater voltage stability is a function of the stability of the -10 volt reference and not the heater supply. The voltage drop of the lead wires of the heaters is compensated for, since the amplifier 'looks at' the voltage directly across the heater.

A shunt is placed in series with each heater element to measure the current. The heater voltage is also measured so that heater power can be determined independently from the control circuit. The accuracy of the data is not, therefore, in any way related to the operation of the control circuit.

The circuitry in the dotted box is on a printed circuit card (6 channels per board).

A protection system has also been designed which provides for shutdown of the heater power supply in any of the following cases: overvoltage on the bus bar, over-voltage on any individual plate or overcurrent on the bus. The system guards against exceeding a total bus current of 750 amps or bus voltage of 22 volts using fixed-level detectors. A variable-level detector circuit is used as a check against individual plate over-voltage. One potentiometer setting guards all of the plate circuits against the same maximum plate voltage. If any of the individual plate voltages exceeds the set value, the D.C. power supply is shut down. Circuit diagrams for these protection circuits are shown in Fig. C.2. They have been wired on a single printed circuit board and are included with the control circuits in a common card box behind the circuit control panel. A photograph of the back of the plate power control panel is shown in Fig. C.3. One of the four printed circuit boards is shown on an extender card above the card box. The bundle of wires at the bottom of the box include both the control leads to the power transistors mounted on the water-cooled bus bar and the heater voltage sense leads. A photograph of the bus bar box is shown in Fig. C.4. The heavy cables at the rear of the box are the power leads from the overhead bus bars.

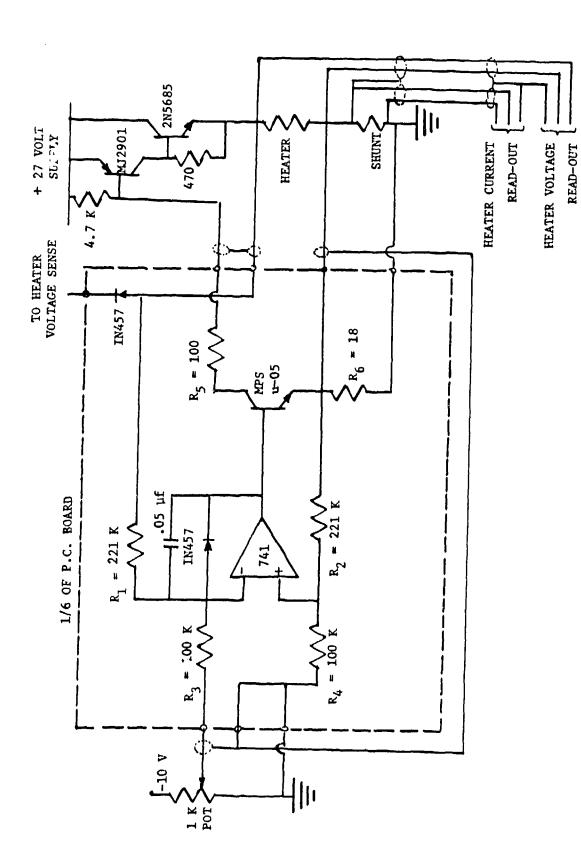
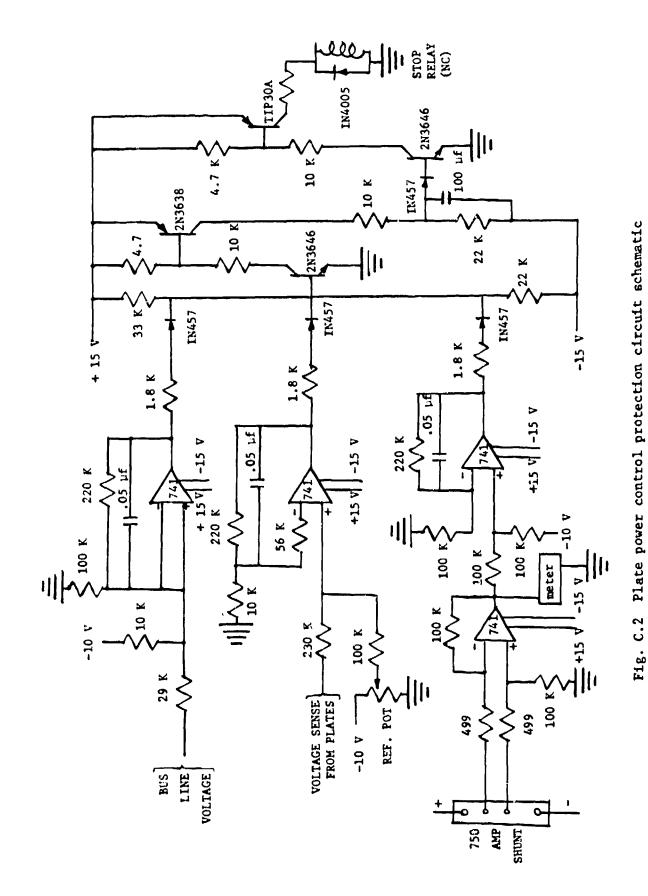


Fig. C.1 Plate power control circuit schematic



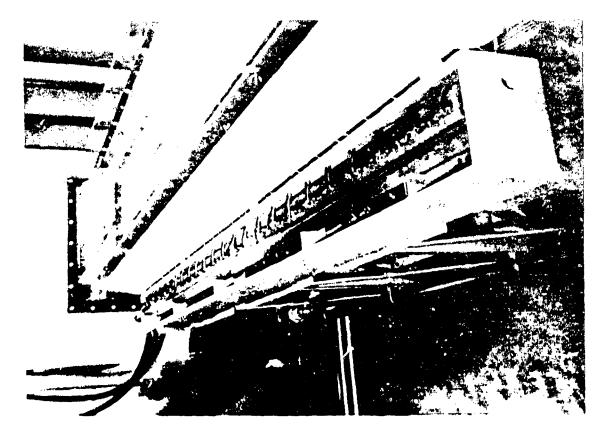


Fig. C.4 Photograph of the bus bar box

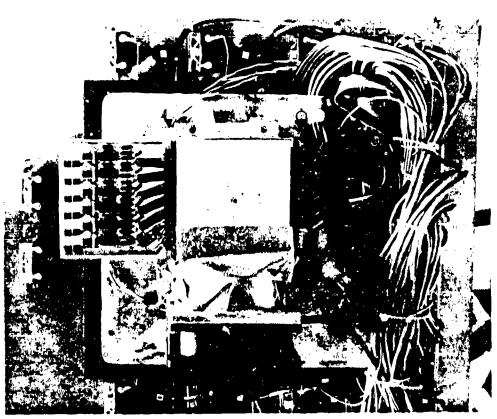


Fig. C 3 Photograph of the rear of the plate power control panel

APPENDIX D

POROUS PLATE FABRICATION

In order to evaluate the practicality of brazing the porous plates, the brazing of 'logpiles' made of segments of copper wire was attempted first. This was done by plating (with electroless nickel) several 1/2 inch long pieces of 0.050 inch dismeter copper wire which were then placed into a stainless steel mold and heated for seven minutes at 1700°F in a hydrogen atmosphere furnace. The result was a success: a brazed 'logpile' of copper wires. With this established, a search was begun for a supplier of copper balls for the porous plates. A 5 lb. sample of copper balls was obtained from the Pioneer Steel Ball Company of Unionville, Connecticut. Finding a plating contractor to plate the balls proved more difficult than had been anticipated. A part of the first sample of copper balls was sent to the Naval Weapons Lab Plating Shop at China Lake. The remainder were distributed among several local plating shops who were asked to bid on the job of plating 100 lbs. of the balls. As each sample was returned, a small plug was brazed in the Hansen Labs furnace using a quartz tube as a mold to test the quality of the braze. The ball sample sent to China Lake was used to assemble a 4 inch square plate for the first large-scale demonstration of the brazing technique. The mold for this sample was constructed with OFHC side rails and a sheet of 1/8 inch quartz for the bottom. The advantage of the transparent bottom was that the ball arrangement could be examined to ensure the ball array on the bottom of the mold had not been disturbed as successive layers were added. There was some hope initially that if the balls were carefully shaken or jarred they would arrange themselves in their most dense array; but after considerable experimentation this hope was finally abandoned and the balls were arranged by hand, one layer at a time, until the desired thickness was achieved. The 4-inch mold was successfully fired in the Hansen Lab furnace.

Work on four full-size mold cavities was then begun. High cost and slow delivery for the quartz bottoms mitigated against their use for

these molds and OFHC copper was used instead. An order for the full 100 lbs. of balls was placed, and Electro-Coatings, Inc. of Emeryville, California, was selected as the plating contractor. They had developed a small stainless steel basket to tumble the balls to improve the plating uniformity. Samples were then obtained with various plating thicknesses and used to determine the plating thickness which would provide the strongest bond without excessive buildup of the braze material at the ball connections. A small mold which produced I inch square samples was used for these test braze samples. A plating thickness of 0.0005 inches was finally selected. The plating contractor agreed to plate the balls in 9 lb. batches, which was just sufficient for a single 4 × 18 inch plate 1/2 inch thick. It was found that difficulty was encountered if a plate was assembled with balls from different plating batches. Small differences in plating thicknesses resulted in small differences in the ball sizes which didn't show up until several hundred of them had been arranged in rows in the mold.

The firing of these large molds was not possible in the Hansen Lab tube shop, but space and furnace time were provided to us at the Stanford Linear Accelerator (SLAC) tube shop. The SLAC tube shop has a 24" hydrogen atmosphere brazing furnace, one of the few of this size on the West Coast. The tube shop supervisor, Mr. Bob Bosenburg, and his people spent considerable time educating us in 'brazing technique'. Keeping the molds and balls clean and use of the proper release agant to keep the plates from sticking in the molds were only part of this help. The balls were assembled in the wolds by hand, filling four molds each time before firing. To assist in this phase of the project, two temporary employees were hired. The ball pack used consisted of 11 layers of balls, each layer 91 rows wide at the test surface with 357 balls in each row. The ball layer behind the surface layer contained 90 rows and all other layers used 89 rows. This construction provided a lip at the top surface of the plate so that when two plates were laid side by side they would form a continuous surface but there would be space between them beneath the top layer for the phenolic support strips. As already noted, the final layer of balls was arranged so that a ball row was omitted between every ten

rows to form grooves for the heater wires. A photograph of a ball assembly nearly completed in a copper mold before firing is shown in Fig. D.1. Fig. D.2 shows a photograph of the molds being filled. As shown in this phorograph, most of the ball arrangement was done with small brushes. With some practice it became possible to literally paint the ball rows into place. Each plate contained on the order of 350,000 balls. Assembly of the ball pack took between three and four days for an experienced worker.

Another problem that the SLAC tube shop personnel worked out for us was the correct 'heating pattern' for the brazing of the full-vize plates. The heating pattern finally adopted was to heat the plates to 840°C and allow the molds and plates to equalize at this temperature. This appears to be just under the melting temperature of the phornickel plating on the copper balls. After equalizing all the furnace thermocouples to within a few degrees, the plate and mold temperature was increased to 910°C and held for ten minutes. During this finel heating, all mold thermocouples were held within a ten degree band to avoid uneven heating. At the end of the ten minute cycle the furnace heating element was removed and the steel retort was air-quenched. Typical cooling times for this process was usually in excess of twelve hours before the plates could be removed from their moids. This particular heating pattern seemed to optimize the brazing process. At lower temperature, the braze fillets formed between balls were smaller and often would not hold. At higher temperatures, the plating seemed to evaporate or be absorbed into the balls and there was less available for the braze.

Figure D.3 is a photograph of the arrangement used for the molds on the pedestal of the 24 inch SLAC furnace. A photograph of a tube shop technician removing the heating element from the steel retort after firing is shown in Fig. D.4. A photograph of the finished plate, after the copper molds were removed is shown in Fig. D.5. The black material on the mold is the release agent used to avoid having the balls braze themselves to the mold. It was found after several firings the copper molds became extremely soft and lost all temper. Inspection was needed after each firing, before reassembly, to ensure that the side walls were

straight and square. Fig. D.6 is a close-up of the plate edge after firing. The ball layer lip is clearly shown. As described earlier, this lip made it possible to construct the rough surface test section without interruptions in the roughness pattern at plate joints.

The successful fabrication of the Roughness Rig plates was possible thanks to the help of several people. Particularly, the author is thankful to Nick Andrews, foremen at the Hansen Lab Tube Shop who first helped me successfully braze the trial samples. To Manny Gill of Electro-Coatings, Inc., who worked out the batch process which allowed the successful plating of the 8 million balls that were used. To Bob Bosenburg, Supervisor of the SLAC Tube Shop, and his people, who gave us space, furnace time and the benefit of their many years of experience in brazing the full size plates. Finally, we're thankful to the two temporary employees who helped assemble the plates, Jim Burlison and George Zanetti, whose ball-stacking ability was surpassed only by their patience and good humor.

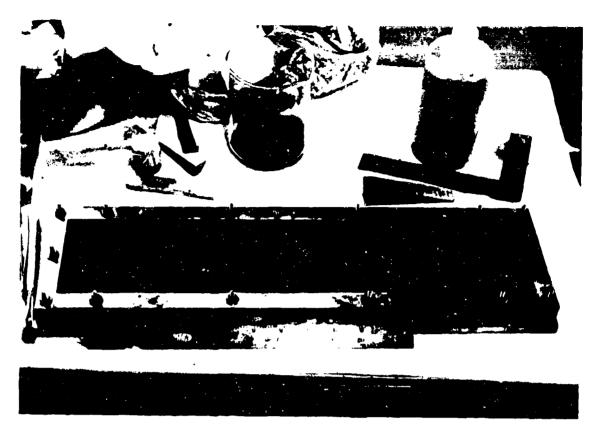


Fig. D.1 Photograph of ball assembly in mold, almost complete



Fig. D.2 Photograph of balls being arranged in the molds

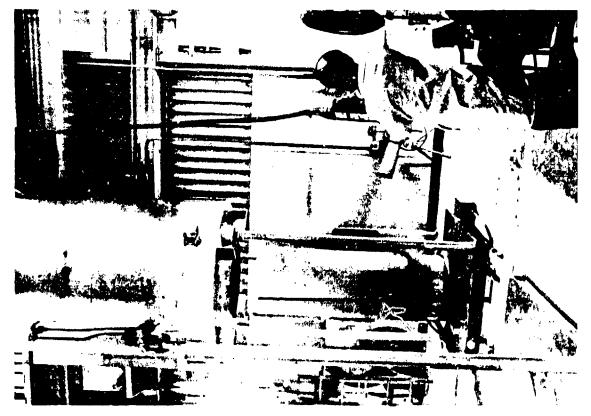


Fig. D.4 Photograph of furnace being raised after brazing the ball assembly

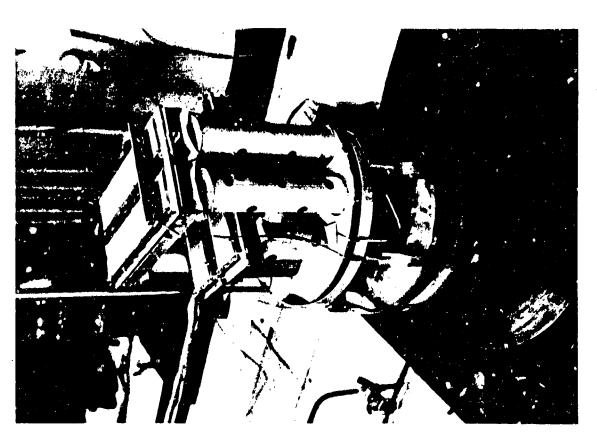


Fig. D.3 Photograph of the molds on the furnace pedestal befor firing

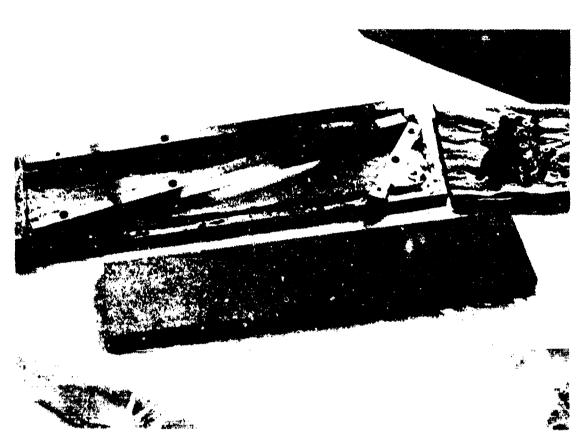


Fig. D.5 Photograph of the brazed ball plate with mold after firing

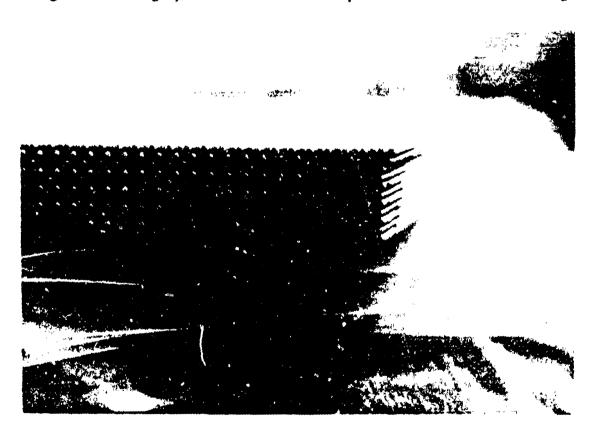


Fig. D.6 Photograph of an end view of the plate after firing

APPENDIX E

TABULATION OF EXPERIMENTAL RESULTS

This appendix contains a summary of the Stanton number data taken in these tests, along with a limited number of velocity profiles. Included are Stanton data for each blowing fraction and for each free-stream velocity tested. For each of the test conditions, an initial velocity profile has been included. This is typically the first profile measured once the boundary layer was fully turbulent. The purpose of including this profile is to provide a 'starting point' for prediction schemes, for those wishing to use these data to check a predictive method. The following is a summary of the abbreviations used in the data listings.

E.1 Stanton Number Data Listings

VEL	Free-stream velocity for the test.
T, TO	Free-stream static and total temperature.
RHO	Free-stream density.
P	Free-stream static pressure.
TDB, TWB	Dry and wet bulb temperatures.
PAMB	Ambient pressure.
VIR ORIG	Turbulent boundary layer virtual origin, determined from the curve fit to momentum thicknesses. Measured from the test section inlet, positive values are downstream.
PL NO	Plate number.
x	Distance along test section, from inlet.
STN NO	Stanton number.
MOM REY NO	Momentum thickness Reynolds number.
ENTH REY NO	Enthalpy thickness Reynolds number.
X REY NO	x-Reynolds number, measured from test section inlet.
(X-XO) REY NO	x-Reynolds number measured from the virtual origin.
F	Blowing fraction.
В	F/St.

TPL, TAIR Plate surface and transpiration air thermocouple readings.

MOMTH, ENTTH Momentum and enthalpy thickness for each plate.

Momentum thicknesses were obtained by plotting the 6 or 7 measured values of momentum thicknesses and interpolating to find values for each plate. Enthalpy thickness for each plate was obtained from the integral energy equation using the measured Stanton numbers.

DIS TH Displacement thickness from measured velocity profiles.

MOM TH Momentum thickness from measured velocity profiles.

SHAPE FACTOR Displacement thickness/momentum thickness.

X-XO Distance from virtual origin.

MOM TH (FIT) Momentum thickness from least-squares fit to momentum thickness data.

DIFF Difference between the measured momentum thickness and that from the curve fit.

CF2 (FIT) Skin friction based on momentum thickness fit.

CF2 SMOOTH Smooth surface skin friction, taken from

$$\frac{c_f}{2}$$
 = .0125 Re $_{\theta}^{-0.25} \left[\frac{\ln{(1+B)}}{B} \right]^{0.7}$ (E-1)

RATIO Rough surface skin friction/smooth surface skin friction.

ROUGH RE NO Roughness Reynolds number.

E.2 Velocity Profile Listings

F S VEL Free-stream velocity.

BLO FR Nominal blowing fraction for the run

STAT TEMP Free-stream static temperature.

STAT PRES Free-stream static pressure.

Plate over which the velocity profile was taken.

DIST Distance from test section inlet.

Y - YTOP Distance from crest of rough surface balls to center

of probe, with no corrections applied.

V/VINF Local-to-free-stream velocity ratio.

```
102:4 98:2 0:090 0:105
102:4 98:2 0:090 0:115
102:4 98:5 0:109 0:124
102:4 97:5 0:118 0:134
                                                                                                                                                                             2158 1006100 681199 0.
2519 1075487 750586 0.
2478 1144875 819972 0.
                                                                                                                                                                                   2158 1006100
                                   58 0.00236
                                                                                                                                  1891
2047
                                                                                                                                                                                                                                                                                                                                                                                                               v.
                                    65 0.00554
                                62 0.00229 2047 2319 1079487 770580 0.00229 2186 2478 1144873 819972 0.00200 0.00225 2324 2655 1214259 884958 0.74 0.00220 2403 2790 1285045 958744 0.20218 2000215 2723 3092 149248 1097517 0.86 0.00212 2845 3240 1491804 1164903 0.90 0.00211 2984 3387 1561190 1256289 0.94 0.00211 3105 3534 1650576 1305675 0.
                                                                                                                                                                                                                                                                                                                                                                                                               0.
                                                                                                                                                                                                                                                                                                                                                                                                                                                              102.5 96.9 0.126 0.143
                                                                                                                                                                                                                                                                                                                                                                                                                                                           102.5 97.3 0.136 0.145
102.5 97.5 0.134 0.152
102.5 97.9 0.142 0.161
102.5 98.1 0.150 0.170
102.4 98.1 0.157 0.178
102.4 98.1 0.164 0.187
102.4 98.0 0.172 0.125
102.4 98.5 0.179 0.204
                                                                                                                                                                                                                                                                                                                                                                                                              U.
        20
                                                                                                                                                                                                                                                                                                                                                                                                               v.
                                                                                                                                                                                                                                                                                                                                                                                                              0.
                                                                                                                                                                                                                                                                                                                                                                                                              U.
    Skin FRI ColfF and virtural origin from curve fit to mom the progressions curve fit there 0.0059*(X-X0)** 0.7957 *vir origin=2 18.725 from 5 FE X LIS 1H Mom TH SHAPE X=X0 MOM TH DIFF CF2 CF2 RATIO ROUGH NO IN INCH INCH FACTOR INCH (FIT) INCH (FIT) SMOOTH RE NO 2 6 0.00417 0.0175 2.4113 4 14 0.0050 0.0227 2.2891 6 22 0.0590 0.0227 2.2891 6 22 0.0590 0.02263 2.2469 8 30 0.0099 0.0314 1.0889 11.28 0.0406 -0.0008 0.00288 0.00242 1.18 29.09
                                                                                                                                                                                                                                  11.28 0.0406 -0.0008 0.00286 0.00242
19.28 0.0621 0.0015 0.00256 0.00220
51.28 0.0412 0.0003 0.00251 0.00148
45.28 0.1160 -0.0018 0.00210 0.00145
59.28 0.1515 0.0008 0.0020 0.00175
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                             1.18 29.09
1.16 27.55
1.17 26.18
1.17 25.32
                                    50 0.0905 0.0606 1.4894
50 0.1559 0.0909 1.4730
62 0.1767 0.1198 1.4749
**STARTOR ROLE 87 227/3 V=32 FT/SEC F=:001

**VEL= 35:3 FT/SEC T= 74:8 DEGF TO= 74:9 DEGF RHO=0:0746 EDS/FT5 P= 14:80 PM

**IONE 72:0 DEGF TWB= 62:0 DEGF PARGE 30:06 IP HG VIR ORIGEXO— 5:38 IN

**PL x STR MOM EURH X (x=x0) F B THE DEGF IDEA

**IPC THE STR MOM EURH RETHO RETHO HETHOUSE TO 0:0004 0:000 10:40 THE DEGF INCH IDEA

**IPC THE STR MOM EURH RETHO RETHOUSE TO 0:0004 0:000 10:40 THE DEGF INCH IDEA

**IPC THE STR MOM EURH RETHOUSE TO 0:0004 0:000 THE TOS TO 7:79 0:018 0:005

**IPC THE STR MOM EURH RETHOUSE TO 0:0004 0:000 THE TOS TO 7:79 0:018 0:005

**IPC THE STR MOM EURH RETHOUSE TO 0:0004 0:000 THE TOS TO 7:79 0:018 0:005

**IPC THE STR MOM EURH RETHOUSE TO 0:0004 0:000 THE TOS TO 7:79 0:018 0:005

**IPC THE STR MOM EURH RETHOUSE TO 0:0004 0:000 THE TOS TO 7:79 0:018 0:005

**IPC THE STR MOM EURH RETHOUSE TO 0:0004 0:000 THE TOS TO 0:005

**IPC THE STR MOM EURH RETHOUSE TO 0:0004 0:0004 0:000 THE TOS TO 0:005

**IPC THE STR MOM EURH RETHOUSE TO 0:0004 0:0004 0:0004 0:0004 0:0006

**IPC THE STR MOM EURH RETHOUSE TO 0:0004 0:0004 0:0004 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 0:0006 
       SKID FRE COLFF AND VIRTURAL ORIGIN FROM CURVE FIT TO MON THE MEASUREMENTS
    SKIEFRE COFFE AND VIRTURAL ORIGIN FROM CORMETED OMOM THE MEASUREMENTS
CURVE FIT THE SET OF OUR OF SET OF SE
                                                                                                                                                                                                                                     16.62 0.0785 -0.0010 0.00254 0.00172
24.62 0.1049 0.0012 0.00220 0.00159
32.62 0.1291 0.0011 0.00190 0.00150
44.62 0.1627 0.0000 0.00175 0.00159
56.62 0.1939 -0.0027 0.00159 0.00135
72.62 0.2330 0.0014 0.00143 0.00125
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                           1.38
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                       25.64
24.30
                                 38 0.1944 0.1280 1.5185
50 0.2409 0.1620 1.5181
62 0.2965 0.1960 1.5075
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                            1.20
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                          22.35
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                            1.21
```

```
STARTON RUL 8/21/73 V=32 F1/SEC F=.002

VEL= 34.4 F1/SEC T= 75.1 DEGF TO= 75.2 DEGF RRG=0.0744 EBS/FT3 P= 14.76 PSI

TUBE RL-D DEGF TWB= 64.0 DEGF PARSE 29.97 IN NG VIR ORLOXOS 1.16 IN

NO TRCH NO REYNO 0.0018 1.695
2 G 0.00166 556 315 103/T5 83.662 0.0017 1.044 99.5 76.7 0.051 0.018
3 10 0.00535 847 611 172855 15204 0.0018 0.551 99.5 76.7 0.051 0.018
4 14 0.00271 1100 948 241948 221946 0.0018 0.551 99.5 76.7 0.051 0.018
5 18 0.00271 1300 348 241948 221946 0.0018 0.551 99.5 76.6 0.064 0.055
5 18 0.00273 1.385 1252 51140 291089 0.0018 0.829 99.5 77.1 0.089 0.052
6 22 0.00202 1625 1530 380282 360231 0.0018 0.829 99.5 77.1 0.080 0.052
7 26 0.00191 1884 1702 449424 429373 0.0018 0.895 99.9 76.9 0.094 0.039
8 30 0.00178 2109 2046 516566 498515 0.0018 1.036 99.4 77.2 0.109 0.104
8 30 0.00170 2568 2292 587708 567657 0.0018 1.006 99.4 78.9 0.122 0.118
9 34 0.00170 2568 2292 587708 567657 0.0018 1.006 99.4 78.9 0.122 0.118
10 38 0.00170 2568 2573 72593 705942 0.0018 1.006 99.4 77.5 0.164 0.160
11 42 0.00155 2835 2773 725943 705942 0.0018 1.006 99.4 77.5 0.164 0.160
12 46 0.00161 3077 3003 795135 775089 0.0018 1.096 99.4 77.5 0.104 0.160
15 58 0.00142 3526 3457 933419 913508 0.0018 1.295 99.4 77.5 0.104 0.160
16 02 0.00132 3526 3457 933419 913508 0.0018 1.295 99.4 77.5 0.294 0.206
17 66 0.00139 4573 4531 120948 118937 0.0018 1.295 99.4 77.6 0.224 0.226
18 70 0.00139 4573 4531 120948 118937 0.0018 1.285 99.4 77.6 0.224 0.226
19 74 0.00129 4996 4960 1417814 1597355 0.0018 1.285 99.4 77.6 0.224 0.226
20 78 0.00121 5203 516 140057 1406050 0.0018 1.285 99.4 77.6 0.246 0.275
21 82 0.00129 4996 4960 1417814 1597355 0.0018 1.557 99.4 76.8 0.255 0.251
22 86 0.00121 5203 516 140057 1466050 0.0018 1.557 99.4 76.3 0.351 0.312
24 94 0.00120 5610 5609 1624841 1004790 0.0018 1.557 99.4 76.3 0.351 0.352
         SKIN FRI COLFF AND VIRTURAL ORIGIN FROM CHRVE FIT TO MOM THI MFASURE MENTS
CURVE FIT THE TATE 0.0077+(x-x0)++ 0.8204 + VIR ORICXO)= 1.157 INCHC5
10.15 IN MOM IN SHAPE x-x0 MOM TH DIFF (F) CF2 RATIO HOUGH
HO IN INCH INCH FACTOR INCH (FIT) INCH (FIT) SMOOTH RE NO
1 2 0.0580 0.0393 3.0075
2 6 0.0671 0.0313 2.1443
3 14 0.0936 0.0343 3.0075
                                                                                                                                                                                                                                           14 0.0985 0.0839 1.5428
22 0.1464 0.0941 1.5607
50 0.1905 0.1222 1.5588
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                             1.28 22.12
1.23 20.77
1.20 19.75
                                          38 0.2329 0.1500 1.5529
50 0.2964 0.1927 1.5585
62 0.3465 0.2293 1.5188
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                     1.15 18.56
    STANTON NUMBER ROLL BY 30773 V=32 FT/SEC F=.004
VLE 33.3 FT/SEC T= 68.0 DeGF TO= 68.2 DEGF ROD=0.0748 LINS/FT3 P= 14.60 Ps1
TDBE 77.0 DEGF TWEE 68.0 DEGF PAMSE 29.78 IN NO VIR ORIGINAS = 53.2 IN
NO TRCH NO HEYNO REYNO REYNO REYNO REYNO DEGF CO OEGF TO TAIR MOMTH ERTTH
NO TRCH NO HEYNO REYNO REYNO REYNO DEGF CO OEGF TO TAIR MOMTH ERTTH
1 2 0.00057 592 160 34067 90018 0.0041 1.421 94.9 47.5 0.023 0.003
3 10 0.00240 994 960 1024 170335 226866 0.0041 1.421 94.9 47.5 0.025 0.003
3 10 0.00245 1346 1024 170335 226866 0.0041 1.421 94.9 47.5 0.026 0.003
3 10 0.00245 1346 1024 170335 226866 0.0041 2.197 94.8 71.9 0.096 0.003
4 14 0.00192 1737 1454 238469 295120 0.0041 2.197 94.8 71.4 0.102 0.005
5 18 0.00112 2146 1859 306603 363155 0.0041 2.197 94.8 71.8 0.126 0.105
6 22 0.00150 2487 2242 374737 431289 0.0041 2.197 94.8 71.8 0.126 0.105
6 30 0.00125 3500 3561 579140 635691 0.0041 2.711 99.6 72.4 0.167 0.152
7 26 0.00116 3918 3724 647.74 703625 0.0041 3.544 94.8 72.4 0.167 0.152
11 40 0.00116 4275 4083 714604 5855 0.0041 3.544 94.8 72.0 0.230 0.211
12 46 0.00116 4275 4083 714604 771950 0.0041 3.544 94.8 72.0 0.230 0.211
13 50 0.00099 5003 4790 851676 90827 0.0041 3.549 94.8 72.0 0.230 0.211
14 59 0.00116 5566 5157 919610 976361 0.0042 3.787 94.7 73.1 0.272 0.201
15 58 0.00106 5723 5401 987044 1044499 0.0041 3.865 94.8 72.4 0.315 0.352
16 62 0.00095 6098 8539 1056078 112629 0.0041 3.585 94.8 72.4 0.355 0.356 0.352
16 62 0.00099 6098 5839 1056078 112629 0.0041 3.585 94.8 72.4 0.355 0.356 0.352
16 62 0.00009 6098 5839 1056078 112629 0.0041 3.585 94.8 72.4 0.355 0.356 0.352
16 62 0.00009 6098 5839 1056078 112629 0.0041 3.585 94.8 72.4 0.355 0.356 0.352
16 62 0.00009 6098 5839 1056078 112629 0.0041 3.585 94.8 72.4 0.355 0.356 0.352
16 62 0.00009 6098 5839 1056078 112629 0.0041 5.365 94.8 72.4 0.355 0.356 0.352
17 60 0.00009 6098 5839 1056078 112629 0.0041 5.375 94.8 72.4 0.355 0.356 0.352
18 70 0.000009 684 874 851076 0.0042 4.107 94.8 72.5 0.400 0.344
22 86 0.00077 887 7885 10560 1589560 0.0041 4.701 94.8 72.5 0.400 0.365
24 
               SKIN FRE COLFF AND VIRTURAL ORIGIN FROM CURVE FIT TO MOM THE MEASUREMENTS
            SKIN FRI COLFF AND VIKTURAL ORIGIN FROM CHRYE FIT IC MOM THE MEASUREMENTS
CURVE FIT THETA: 0.0066+(x-xu)+* 0.9550 .*VIR ORIEXO) = -5.519 INCHES
PL x u15 TH MOM TH SHAPE x-xu MOM TH DIFF CF2 CF2 RATIO ROUGH
NO IN INCH FACTOR THEH (FIT) THEH (FIT) SMOOTH RE NO
0 0 0.0255 0.0152 1.76645
1 2 0.0754 0.0255 3.1527
2 0 0.0927 0.0957 1.6645
3 10 0.41218 0.0775 1.5575 15.52 0.0784 0.0010 0.00091 0.00098 0.95 14.69
5 18 0.2076 0.1258 1.2530 21.42 0.1227 -0.0028 0.00179 0.00098
                                                                                                                                                                                                                                                13.32 0.0784 0.0010 0.00091 0.00098 21.32 0.1227 -0.0028 0.00079 0.00080 3.53.52 0.1279 -0.0001 0.00068 0.00069 45.52 0.2519 0.0068 0.00059 61.32 0.4400 0.0021 0.00846 0.00045
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                             0.99 15.08
1.02 12.67
1.02 11.94
                                       18 0.2076 0.1256 1.0530
                                               30 0+3159 0:1880 1:6700
42 0:4152 0:2452 1:6456
                                             58 0+5726 0.3431 1.6688
                                                                                                                              0.43/8 1.6705
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                        1.02
```

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SKIR FRI COLFF AND VIRIURAL DRIGHT FROM CURVE FIT TO MOM THI MEASUREMENTS CURVE FIT THE NATION OF SKIR FRI COLFF AND VIRIURAL DRIGHT FROM CURVE FIT TO MOM THI MEASUREMENTS CURVE FIT THE NATION OF SKIR FRI COLFF AND VIRIURAL STREET OF SKIR FRI COLFF AND VIRIURA
| STANTOR HUR | 5/307/3 | RO | 4L0W | V=90 | FT/31C | | | | | | | | | |
| VIE2 | 90.8 | F1/34C | 12 | 76.5 | RO | T05 | 77% | DEGE | ROOZD | 07/35 | CO5/F13 | DT | 14.63 | P/31 |
| 108- | 08-0 | DEGE | TWO | 67.5 | DEGE | PAMOZ | 20.71 | TR | HG | VIR | ORTO | XOZ | 3.53 | TR |
| EL | X | STR | MOM | RETHO | DEGE | DEGE
                                                                                                                                                                                                                                                                                                                                                                                 99-7 97-9 0-152 0-193

99-7 97-1 0-162 0-193

99-7 97-1 0-162 0-192

99-7 96-3 0-170 0-161

93-6 95-8 0-179 0-170
                                    54 0.00218
58 0.00218
                                                                                                              6852
7383
                                                                                                                                                     6467 2434254 2275126 0.
6860 2614569 2455441 0.
                                                                                                                                                                                                                                                                                                                                        u.
                                  58 0.00218 7503 6860 2614509 2455441 0.

60 0.00215 8069 7642 2975200 2816071 0.

70 0.00215 8450 8026 5155515 2996887 0.

78 0.00209 9196 8786 5515619 557552 0.

82 0.00209 9197 9159 569640 557552 0.

86 0.00208 9917 9552 5876775 5717647 0.

90 0.00204 0278 9905 4057090 5897462 0.
                                                                                                                                                                                                                                                                                                                                                                                  99.6 99.9 0.187 0.178
99.7 98.7 0.195 0.197
99.7 98.4 0.204 0.195
99.7 98.5 0.212 0.203
              18
                                                                                                                                                                                                                                                                                                                                      υ.
            21
                                                                                                                                                                                                                                                                                                                                      u.
                                                                                                                                                                                                                                                                                                                                                                                  99.6 97.2 0.220 0.211
99.6 97.0 0.228 0.220
                                                                                                                                                                                                                                                                                                                                        U.
                                    94 0.00204 10639 10271 4237405 4078277 0.
                                                                                                                                                                                                                                                                                                                                                                                   99.6 99.2 0.236 0.228
              SKIR FRE COLFF AND VIRTURAL ORIGIN FROM CURVE FIT TO MON THE MEASUREMENTS
           PL X 515
NO IN IN
                                                FIT THETA BLODYGECK-XD)** U.7597 FUR ORICXD:= 3.538 INCHES

DIS TH MOM IN SHAPE X-XD NOM IN DIFF CF2 (F2 RATIO ROUGH

R INCH INCH FACTOR INCH (FIT) INCH (FIT) SMOOTH RE NO
                                           0 0.0190 0.0090 2.1178
2 0.0237 0.0095 2.5047
5 0.0307 0.0182 1.6921
                                   10 0.0611 0.0560 1.6946
18 0.0945 0.0596 1.5842
30 0.1454 0.0929 1.5660
                                                                                                                                                                                               14.47 0.0593 -0.0003 0.00309 0.00174
                                                                                                                                                                                                                                                                                                                                                                                                                                                         1./8
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                            78.71
                                                                                                                                                                                             26:47 0:0935 0:0007 0:00267 0:00155
36:47 0:1240 -0:0001 0:00243 0:00145
54:47 0:1012 -0:0006 0:00223 0:00135
                                    42 0.1908 0.1241 1.5372
58 0.2426 0.1618 1.4996
                                                                                                                                                                                                                                                                                                                                                                                                                                                         1.68
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                             69.81
                                      78 0.3030 0.2039 1.4863
                                                                                                                                                                                                 74.47 0.2041 0.0003 0.00207 0.00128
```

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STANTON HUI4 8/24/73 V=90 FT/SEC F2:001

VEL= 90.5 F1/SEC T2 77:0 DEGF TO= 77:7 DEGF RHOZD:0736 LG5/FT3 P= 14:05 PS1

IDBE 80.0 DEGF TWE= 64:0 DEGF PAMIE 21:79 IN HG VIR ORIGENCE 6:5R IN

REVIN REVINO REVINO REVINO REVINO REVINO DEGF RHOZD:0736 LG5/FT3 P= 14:05 PS1

NO INCH NO REVINO REVINO REVINO REVINO DEGF RHOZD:0736 LG5/FT3 P= 14:05 PS1

NO INCH NO REVINO REVINO REVINO DEGF PAMIE 21:79 IN HG VIR ORIGENCE 6:5R IN

NO INCH NO REVINO REVINO REVINO DEGF PAMIE 21:79 IN HG VIR ORIGENCE 6:5R IN

NO INCH NO REVINO REVINO REVINO DEGF PAMIE 21:79 IN HG VIR ORIGENCE 6:5R IN

1 2 0:00159 324 849 241 84935 65783 0:0011 0:682 100:7 77:4 0:028 0:0015

3 10 0:00315 2060 1828 449177 423125 0:0011 0:549 100:6 77:1 0:046 0:051

4 14 0:00271 2695 2524 6:8848 602/96 0:0011 0:409 100:6 77:1 0:046 0:051

5 18 0:00249 308 3840 9A8189 96:137 0:0011 0:402 100:7 77:8 0:060 0:057

7 26 0:00229 308 3840 9A8189 96:137 0:0011 0:402 100:8 77:8 0:087 0:085

8 30 0:00207 5121 5020 13:4/531 13:1879 0:0011 0:553 100:8 77:8 0:087 0:085

8 30 0:00207 5121 5020 13:4/531 13:1879 0:0011 0:553 100:8 77:9 0:114 0:114 0:08 0:0011 0:553 100:8 77:9 0:114 0:114 0:08 0:0011 0:553 100:8 77:9 0:152 0:124 0:104 0:0011 0:553 100:8 77:9 0:152 0:124 0:104 0:0011 0:553 100:8 77:9 0:152 0:154 0:154 0:001 0:553 100:8 77:9 0:152 0:154 0:154 0:001 0:553 100:8 77:9 0:152 0:154 0:154 0:001 0:553 100:8 77:9 0:152 0:154 0:154 0:001 0:553 100:8 77:9 0:152 0:154 0:154 0:001 0:553 100:8 77:9 0:152 0:154 0:154 0:001 0:553 100:8 77:9 0:152 0:154 0:154 0:001 0:553 100:8 77:9 0:152 0:154 0:154 0:001 0:553 100:8 77:9 0:152 0:154 0:154 0:001 0:553 100:8 77:9 0:152 0:152 0:154 0:001 0:544 0:00 8 77:9 0:152 0:154 0:154 0:001 0:544 0:00 8 77:9 0:152 0:154 0:154 0:001 0:544 0:00 8 77:9 0:152 0:154 0:154 0:001 0:544 0:00 8 77:9 0:245 0:221 0:214 0:001 0:544 0:00 8 77:9 0:245 0:221 0:214 0:001 0:544 0:00 8 77:9 0:245 0:221 0:214 0:001 0:544 0:00 8 77:9 0:245 0:221 0:221 0:221 0:221 0:221 0:221 0:221 0:221 0:221 0:221 0:221 0:221 0:221 0:221 0:221 0:221 0:221 0:221 0:221 0:221 0:221 0
               SKILLERI COLFF AND VIRIONAL ORIGIN FROM CURVE FIT TO MON THE MEASUREMENTS
         SMIT FRI CORP AND VIRIORAL CRIDIN FROM CHRVE FIT TO MOR THE MENTAL COURSE FIT THE COURSE COURSE FIT THE COURSE COURSE FIT TO SHOOTH COURSE COURSE FIT TO SHOOTH COURSE COU
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                               RATIO ROUSH
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                         0.42 0.0450 -0.0015 0.00241 0.00150 1.42 0.0862 0.0033 0.00237 0.00125 55-42 0.1254 -0.0000 0.00211 0.00110 45-42 0.1632 -0.0007 0.00194 0.00102 61442 0.2084 -0.0035 0.00171 0.00045 77.42 0.2525 0.0045 0.0017
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                              77.90
70.20
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                        1.90 70.70
1.92 56.32
1.91 61.97
                                                               22 0:1326 0:0849 1:5624
34 0:1976 0:1270 1:5662
                                                  40 0.2512 0.1639 1.5329
62 0.3259 0.2125 1.5340
78 0.3766 0.2500 1.5060
STARTON, NOR BZ2373 V=90 FTZSEC F=002

VEC= 89.7 FTZSEC | 12 75.0 DEGF | TOT 75.7 DEGF | PRO20.0745 | UGZFT3 | PZ 14.75 P21

TDBZ 77.0 DEGF | TWHE 64.0 DEGF | PAMEE 29.96 In RG | VIP ORIGINATE | O.40 IR

PC | X | STH | MOM | ENTH | X | X = A0) | F | B | TOL | TAIR MOME ENTH | X | X = A0) | F | B | TOL | TAIR MOME ENTH | X | X = A0) | F | B | TOL | TAIR MOME ENTH | TAIR MOME ENTH | X | X = A0) | F | B | TOL | TAIR MOME ENTH | TOL | TOL | TAIR MOME ENTH | X | X = A0) | F | B | TOL | TAIR MOME ENTH | TOL | TOL | TAIR MOME ENTH | TOL | TOL | TAIR MOME ENTH | TOL | TOL
   SKIN FRI COLFF AND VIRIURAL ORIGIN FROM CURVE FIT TO MOM THI MLASUREMENTS CURVE FIT THETAS U.OUTI*(X-X0)** U.ROT7 *VIR ORIGINS U.QUU INCHES U.QUU IN
```

```
94 0.00066 22974 21949 4185115 4166416 0.0041 6.291 185.7 77.6 8.516 0.495
        Skin FRE COEFF AND VIRTURAL ORIGIN FROM CURVE FIT TO MOM THE MEASUREMENTS
      Skir FFI CupF AND VIRIORAL ORIGINA FROM CINCE FIT TO MOM THE MISSING MILES COMMERTED FOR CORRESPOND ORIGINAL OR
                                                                                                                                                                                                      9.58 0.0075 -0.0017 0.00728 0.00074 (1.58 0.1376 0.0025 0.00169 0.00055 0.0055 0.00069 0.00046 0.00046 0.0051 0.0051 0.00046 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0051 0.0
                                 34 0.3395 0.2011 1.0883
46 0.4494 0.2637 1.0489
67 9.5826 0.3535 1.0478
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                       3.02 52.81
2.95 48.97
2.87 45.03
                                                     0.6824 0.41/9 1.6329
SKI). FRE COLFF AND VIRTURAL ORIGIN FROM CURVE FIT TO MOM THE MEASUREMENTS

    4.75
    0.0019
    -0.0008
    0.00524
    0.000003.404
    77.76

    8.75
    0.1049
    0.0004
    0.00234
    0.0000038.706
    60.05

    20.75
    0.209
    0.0016
    0.00117
    0.0000024.801
    46.97

    32.75
    0.5275
    0.0000
    0.00000
    0.0000024.801
    22.14

    44.75
    0.4287
    -0.0000
    0.0000076.908
    0.000076.908

    76.75
    0.6827
    -0.0003-0.00035
    0.0000027.277
    0.

                              6 0.1049 0.0627 1.6774
10 0.1893 0.1049 1.7735
22 0.4338 0.2193 1.9785
54 0.6567 0.5268 2.0046
46 0.8001 0.4327 1.9877
02 1.1075 0.5550 1.9933
                                  78 1.3833 0.6829 2.0256
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SKIN FRE COLFF AND VIRTURAL ORIGIN FROM CURVE FET TO MOM THE MEASUREMENTS
               SKIR FR: COLFF AND VIRIORAL ORIGIN FROM CURVE FIT TO MOM THE MEASURE ME CURVE FIT THE LAT 0.0005+(7-XU)++ 0.7055 +VEP ORIGINS = -0.701 IRCHE CURVE FIT THE LAT 0.0005+(7-XU)++ 0.7055 +VEP ORIGINS = -0.701 IRCHE CURVE FIT THE LAT 0.0005+(7-XU)++ 0.7055 +VEP ORIGINS = -0.701 IRCHE CURVE FIT THE LAT 0.0005+(7-XU)++ 0.000
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                        RALLO ROUSH
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                 1.90 123.37
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                   1.89 111. 15
            | STANTON RULE | A7, 773 | V=139 | FT/Set | F=+001 | | | | | | | | | | |
| VEL=137-0 | FT/Set | T= 75.9 | Def | TO= 77.5 | DEF | RHO=0.0739 | E65/FT5 | P= 14.70 | PS1 |
| COLE | 70.0 | Def | TWE | 65.5 | Def | PAMSE | 29.85 | In | Ho | VIR | ORIGINACE | TE | TE | TE |
| RO | HICH | NO | REVINO | REVIN
                                                 94 0.00157 21174 20032 0402330 0541391 0.0010 0.610 101.6 78.6 0.308 0.291
                      SKIN FRA CULFF AND VIRTURAL ORIGIN FROM CURVE FIT TO MOM THE MEASUREMENTS
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94 0.00121 24546 23560 6409166 6503258 0.0019 1.568 99.4 78.0 0.360 0.346
     SKIA FRI COLFF AND VIRTURAL ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS CURVE FIT THETA= U.DUBU+(X-XO)++ 0.8995 *VIR ORI(XO)= -1.383 INCHES
PL X DIS TH MOM TH SHAPE X-XO MOM TH DIFF CF2 CF2 RATIO ROUGH
NO IN INCH INCH FACTOR INCH (FIT) INCH (FIT) SMOOTH RE NO
1 2 U-U295 U-U135 2-1785
                           2 0.0295 0.0135 2.1785 6 0.0651 0.0384 1.6944 10 0.0870 0.0524 1.6611 11.38 0.0537 0.0013 0.00225 0.00103 2.18 101.00 18 0.1462 0.0898 1.6281 19.38 0.0867 -0.0031 0.00202 0.00085 2.39 95.90 30 0.2114 0.1311 1.6127 31.38 0.0867 -0.0031 0.00202 0.00085 2.39 95.90 40.3072 0.1959 1.5886 47.38 0.1958 -0.0021 0.00168 0.00074 2.50 91.28 62 0.3872 0.2496 1.5511 63.38 0.2517 0.0021 0.00168 0.00062 2.67 63.38 0.2517 0.0021 0.00167 0.00057 2.75 63.38 0.4741 0.3091 1.5535 79.38 0.3082 -0.0009 0.00149 0.00052 2.84 79.38
  STANTON RUIL B/c7/73 V=139 FT/SLC F=.004

VEL=135.3 F1/SLC I= 77.3 DLGF TO= 78.8 DEGF RHO=U.0758 LBS/FT3 P= 14.71 P51

TOBE 77.0 DEGF TWD= 65.0 DEGF PAMB= 29.87 In HG VIR ORIGOXD= 1-68 IN

FL X 51N MOM ENTH X (X=X0) F B TPL TAIR MOMIN ENTTH

HO INCH NO KEYNO REYNO REYNO REYNO DEGF DEGF DEGF DEGF DEGF LNCH INCH

1 2 0.00246 1212 855 134622 247705 0.0037 1.522 101.6 78.4 0.018 0.012

2 5 0.00266 3029 2547 403667 516950 0.0037 1.522 101.6 78.0 0.0045 0.038

3 10 0.00185 4712 4185 673112 786195 0.0038 2.052 101.4 77.8 0.070 0.082

4 14 0.00157 6260 5662 942356 1055439 0.0038 2.384 101.5 77.9 0.095 0.184

5 18 0.00140 7808 7095 121601 1254884 0.00359 2.709 101.4 78.2 0.0116 0.105

6 22 0.00130 9222 8514 1480846 1593929 0.0038 2.384 101.5 77.9 0.095 0.184

7 26 0.00118 10635 9897 1750091 1863173 0.0039 3.285 101.5 77.9 0.158 0.147

8 30 0.00111 12049 11248 2019355 2132418 0.0035 3.285 101.5 77.9 0.158 0.147

9 34 0.00101 14074 13868 2557825 2670408 0.0037 3.326 101.2 77.3 0.199 0.187

10 38 9.00108 14674 13868 2557825 2670408 0.0037 3.921 101.2 77.5 0.218 0.205

12 46 0.00108 17299 16424 3096314 3209397 0.0038 3.037 101.2 77.9 0.238 0.225

12 46 0.00108 17299 16424 3096314 3209397 0.0038 3.037 101.2 77.9 0.238 0.225

12 46 0.00108 17299 16424 3096314 3209397 0.0038 3.037 101.2 77.9 0.238 0.225

15 58 0.00088 21338 20255 3908048 4017131 0.0035 4.289 101.4 77.9 0.237 0.262

15 58 0.00088 21338 20255 3908048 4017131 0.0036 4.289 101.4 77.9 0.357 0.350

16 62 0.00089 22684 21519 4173293 4286376 0.0038 4.329 101.4 77.9 0.357 0.350

17 4 0.00089 24730 22801 4482538 4555620 0.0039 4.849 101.4 77.9 0.357 0.359

18 70 0.00081 28001 28029 520272 5363540 0.0034 4.849 101.4 77.9 0.357 0.359

19 74 0.00078 29340 27878 5519516 5632599 0.0038 4.840 101.4 77.9 0.357 0.539

10 74 0.00078 29340 27878 5519516 5632599 0.0038 4.840 101.4 77.9 0.357 0.539

10 74 0.00078 29340 27878 5519516 5632599 0.0038 4.840 101.4 77.9 0.357 0.539

20 78 0.00078 29340 27878 5519516 5632599 0.0038 5.594 101.4 77.9 0.436 0.445

21 0.00077 32040 30379 6058
     SKIN FRI CGLFF AND VIRTURAL ORIGIN FROM CURVE FIT TO MOM TIL MEASUREMENTS CORVE FIT THETAT 0.00/4+(X-X0)+0.09208 *VIR ORI(X0) = -1.677 INCHES PL X UIS TH MOM TH SHAPE X-X0 MOM TH DIFF CF2 CF2 RATIO ROUGH 1 2 0.00405 0.01/6 2.2968 2 0 0.0744 0.0047 1.6651 5 10 0.1176 0.0049 1.6814 11.68 0.0709 0.0010 0.00184 0.00064 2.89 92.01 5 18 0.1995 0.1159 1.7207 19.68 0.1147 -0.0013 0.00162 0.00042 3.41 80.18 8 30 0.3044 0.1788 1.7028 31.68 0.1778 -0.0010 0.00142 0.00042 3.41 80.71 2.46 0.4298 0.2574 1.6698 47.68 0.2591 0.0017 0.00125 0.00035 5.59 75.88 16 62 0.5584 0.3381 1.6517 63.68 0.4157 -0.0005 9.00105 0.000015 4.000015 4.18 89.59
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STANTON RUN 6/1/75 NO BLOW V=19h F7/SEC

VEL=191.0 FT/SEC T= 78.1 DEGF TO= 81.1 DEGF RHO=0.0732 LB5/FT3 P= 14.62 P5/I

TDB= 80.5 DEGF TWB= 65.0 DEGF PAMUE 29.69 IN HG VIA ORIGINAL = -2.09 IN

PL X STN MOM ENTH X (x-x0) F B 1PL TAIN MOMTH ENTTH

NO INCH NO REYNO REYNO REYNO REYNO REYNO DEGF PAMUE 29.69 IN HG VIA ORIGINAL = -2.09 IN

1 2 0.00421 1414 793 188548 442145 0. 0. 96.5 100.1 0.031 0.024

3 10 0.00319 3960 3613 942740 1196337 0. 0. 96.5 100.1 0.051 0.024

4 14 0.00293 5574 4766 1319036 1575433 0. 0. 96.5 100.1 0.051 0.024

5 18 0.00276 6505 5857 1690952 1950529 0. 0. 96.4 96.2 0.069 0.062

6 22 0.00268 7542 6861 2074028 2327626 0. 0. 96.3 94.6 0.080 0.075

7 26 0.00256 8673 7848 2451124 2704772 0. 0. 96.4 96.2 0.102 0.093

9 34 0.00247 9616 8797 2828221 3081818 0. 0. 96.4 96.2 0.102 0.093

9 34 0.00248 10559 9723 3205317 3458914 0. 0. 96.4 96.2 0.102 0.093

10 38 0.00236 11501 10629 3582413 3836010 0. 0. 96.5 95.7 0.132 0.122

12 46 0.00220 12444 11509 3595914 313106 0. 0. 96.5 95.7 0.132 0.122

12 46 0.00222 13387 12380 4J36605 4590202 0. 0. 96.5 95.7 0.132 0.122

14 50 0.00222 15178 14089 5090797 5344394 0. 0. 96.5 95.7 0.132 0.122

15 58 0.00219 16121 14923 5467093 5721440 0. 0. 96.5 96.5 96.6 0.142 0.171

16 60 0.00212 17064 17575 58449489 0.04858 0. 0. 96.5 96.5 96.6 0.180 0.176

17 60 0.00212 17064 17575 58449489 0.04858 0. 0. 96.5 96.5 96.6 0.181 0.167

17 60 0.00212 17064 17575 58449489 0.09858 0. 0. 96.5 96.5 96.1 0.199 0.176

18 70 0.00201 18761 17378 6599181 0852778 0. 0. 96.5 96.5 96.6 0.246 0.225

20 78 0.00202 20457 18977 7353373 7606970 0. 0. 96.5 96.6 96.6 0.246 0.226

21 80 0.00204 23003 21318 8484661 8738259 0. 0. 96.6 98.6 0.246 0.226

22 90 0.00204 23003 21318 8484661 8738259 0. 0. 96.6 98.6 0.246 0.226

24 94 0.0021 23 2385 22105 8861757 9115350 0. 0. 96.6 98.6 0.246 0.226
                         94 0.00213 23851 22105 8861757 9115355 0.
   SKIN FRI COEFF AND VIRTURAL ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS CURVE FIT THETA= 0.0053*(X-X0)** 0.8456 *VIR ORI(X0)= -24694 INCHES PL X DIS TH MOM TH SHAPE X-X0 MOM TH DIFF CF2 CF2 RATIO NO IN INCH INCH FACTOR INCH (FIT) INCH (FIT) SMOOTH B 0 0.0107 0.0059 1.8188 1 2 0.0270 0.0189 1.8188 2 0.0270 0.0189 1.8138
                      6 0.0521 0.0306 1.7029
14 0.0930 0.0574 1.6207
22 0.1246 0.0802 1.5536
                                                                                                                                         16.69 0.0575 0.0001 0.00291 0.00146
24.69 0.0800 -0.0002 0.00274 0.00134
36.69 0.1119 -0.0001 0.00254 0.00123
48.69 0.1421 0.0004 0.00254 0.00116
64.64 0.1807 -0.0003 0.00256 0.00116
80.69 0.2179 0.0001 0.00228 0.00104
                                                                                                                                                                                                                                                                                                                             2.00 159.89
2.05 155.13
2.09 150.46
2.12 147.21
2.16 144.02
                      34 U-1699 U-1120 1-5173
40 U-2143 U-1417 1-5117
62 U-2691 U-1811 1-4862
                        40 0.2143 0.1417 1.5117
62 0.2691 0.1811 1.4862
78 0.3198 0.2178 1.4688
                                                                                                                                                                                                                                                                                                                                2.19 141.58
  94 0.00150 28860 26978 8923784 9132638 0.0009 0.619
                                                                                                                                                                                                                                                                             44.1
                                                                                                                                                                                                                                                                                                         80.3 0.304 0.284
        SMIN FRI COLFF AND VIRTURAL ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS
      SATE FAIL CORPT AND VIRTURAL ORIGIN FROM CHRYL FIT TO MOM THE MEASUREMENTS
CUMPYL FIT THETAT 0.0060+(X=X0)+0 U.8601 *VIR ORIGXD = 2*203 InCHES
PL X DIS TH MOM TH SHAPE X=X0 MOM TH DIFF CF2 CF2 RATIO HOUGH
NO IN INCH INCH FACTOR INCH (FIT) INCH (FIT) SMOOTH RE NO
0 0 0.0101 0.0050 1.8165
1 2 0.0297 0.0167 1.7841
2 0 0.0044 0.0355 1.7275
4 14 0.1002 0.0062 1.0246 16:20 0.0060 =0.0002 0.00254 0.00108 2:36 155555
4 20 0.1004 0.0050 1.5904 0.0035 0.00037 2.02 10007
                                                                                                                                         10:20 0:0000 -0:0002 0:00254 0:00108
24:20 0:0932 0:0003 0:00235 0:00097
30:20 0:1317 -0:0000 0:00217 0:00087
48:20 0:1685 0:0000 0:00295 0:00096
64:20 0:2156 -0:0007 0:00195 0:00074
                                                                                                                                                                                                                                                                                                                              2.42 147.67
2.50 141.86
2.56 137.78
                        22 0.1486 0.0929 1.5995
                        34 0-2056 0-1317 1-5606
46 0-2568 0-1682 1-5264
                                      0.3285 0.2163 1.5180
                                                                                                                                                                                                                                                                                                                                2.62
                                                                                                                                                                                                                                                                                                                                                        133.75
                                                                                                                                                                                                        0.0003 0.00184 0.00069
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STARTOR RUN 8/28/73 V=190 FT/SEC F=:002
VEL=189.7 FT/SEC T= 77.3 DEGF TOF 8U-3 DEGF RHOEU-0736 LB5/FT3 P= 14.67 P51
TDB= 86.0 DEGF TWB= 67.0 DEGF PAMU= 29.79 IN HG VIR ORIG-XO= -3.55 IN
PL X SIN MOM ENTH X (X=XO) F B TPL TAIR MOMTH ENTTH
NO INCH NO REYNO REYNO REYNO REYNO DEGF PAMU= 29.70 IN HG VIR ORIG-XO= -3.55 IN
1 2 0.00314 1601 966 188361 522701 0.0020 0.534 98.6 79.6 0.017 0.010
2 6 0.00291 35/9 2866 56.0082 899422 0.0020 0.744 98.7 78.9 0.038 0.050
3 10 0.00232 5274 4618 941603 1276143 0.0020 0.873 98.8 78.4 0.056 0.049
4 14 0.00206 68/5 6205 1318524 1652864 0.0020 0.873 98.8 78.4 0.056 0.049
6 22 0.00180 9983 9214 2071967 2406307 0.0021 1.114 98.8 79.1 0.090 0.042
6 22 0.00180 9983 9214 2071967 2406307 0.0021 1.159 98.8 79.2 0.106 0.008
7 26 0.00181 11396 10671 2448688 278328 0.0020 1.330 98.8 79.5 0.121 0.112
8 30 0.00164 12903 12098 2825409 3159749 0.0021 1.222 98.8 79.5 0.121 0.113
8 30 0.00164 12921 13091 3202130 3536470 0.0020 1.232 98.9 78.5 0.121 0.113
10 38 0.00164 17047 16215 3955573 4889913 0.0021 1.359 98.9 78.7 0.16/ 0.158
11 42 0.00148 17047 16215 3955573 4889913 0.0021 1.359 98.9 78.6 79.2 0.181 0.172
12 46 0.00152 18459 1/578 4332294 4666534 0.0021 1.359 98.9 78.6 0.210 0.201
15 58 0.00164 19778 18914 4709015 5043355 0.0021 1.350 98.6 79.2 0.181 0.172
16 62 0.00142 23522 22856 5839179 6173519 0.0021 1.450 98.6 79.2 0.181 0.122
17 66 0.00139 2653 2545 5692621 692691 0.0021 1.450 98.6 79.8 0.220 0.229
18 70 0.00152 18459 17578 4535294 96063 760002 1.553 98.7 79.4 0.298 0.299
18 70 0.00131 28060 26738 6969342 7303682 0.0020 1.553 98.7 79.4 0.298 0.294
19 74 0.00131 28060 26738 6969342 7303682 0.0020 1.553 98.8 79.4 0.298 0.294
20 78 0.00132 29384 2/995 7346063 7680403 0.0020 1.550 98.8 78.9 0.524 0.298
20 78 0.00128 33528 31732 8476227 8810567 0.0020 1.551 98.8 79.4 0.298 0.294
21 82 0.00128 35283 31732 8476227 8810567 0.0020 1.551 98.8 79.9 0.351 0.350
22 90 0.00128 35928 31752 8476227 8810567 0.0020 1.561 98.8 79.9 0.351 0.350
      SKIN FRI COLFF AND VIRIURAL ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS CURVE FIT THETA= 0.0050*(X=X0)** 0.9390 *VIR ORIGXO= *2.5488 INCHES VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT THETA= 0.0050*(X=X0)** 0.9390 *VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS VIR ORIGIN FROM CURVE FIT TO MOM THE ORIGIN FROM
   STANTON RUN 8/29//3 V=190 FT/SLC F=004

VEL=190.4 FT/SLC I= 78.0 DEGF TO= 81.0 DEGF RHO=0.0734 LBS/FT3 P= 14.60 PSI

TDB= 84.0 DEGF TWB= 68.0 DEGF PAMB= 29.76 IN HG VIR ORIG/XU= -2.70 IN

HO INCH NG REYNO REYNO REYNO REYNO PFYNO DEGF RHO=0.0734 LBS/FT3 P= 14.60 PSI

2 0.00274 1696 1224 168488 442853 0.0038 1.370
2 0.00274 1696 1224 168488 442853 0.0038 1.371
3 10 0.00170 6596 57/3 942240 1196.45 0.0038 2.594
3 10 0.00170 6596 57/3 942240 1196.45 0.0038 2.594
4 14 0.00146 8669 7/87 1319136 1573541 0.0038 2.594
5 18 0.00125 12532 11685 2072924 2527353 0.0038 3.001 97.6 79.5 0.112 0.104
6 22 0.00125 12532 11685 2072924 2527353 0.0038 3.001 97.6 79.5 0.112 0.104
7 26 0.00117 14416 13575 2449625 2704229 0.0038 3.277 97.8 79.6 0.155 0.124
9 34 0.00106 18091 17269 3203617 3456022 0.0038 3.514 97.7 79.6 0.155 0.124
9 34 0.00107 19881 19097 3580513 3634918 0.0038 3.514 97.7 79.0 0.112 0.164
9 38 0.00107 19881 19097 3580513 3634918 0.0038 3.514 97.7 79.6 0.155 0.124
10 2 0.00092 21766 20905 3957409 4211414 0.0037 3.785 97.7 79.6 0.201 0.203
11 02 0.00090 27231 26251 5058097 5382502 0.0038 3.594 97.7 79.6 0.251 0.221
12 46 0.00107 28660 22046 4334305 4588710 0.0037 3.785 97.7 79.6 0.201 0.203
15 50 0.00098 3.5940 24475 9711201 4965606 0.0037 3.869 97.9 79.9 0.220 0.260
16 50 0.00098 38515 36590 31627 0218785 6475190 0.0039 4.003
17 60 0.00098 38515 36587 35168 6972577 729482 0.0039 4.003
18 70 0.00085 36580 33395 6596810 0.0039 4.003
19 70 0.00085 36580 33395 6596810 0.0039 4.003
10 0.00099 4.0008 38255 36945 7349473 7603478 0.0039 4.004
10 70 0.00085 36580 33395 6597671 0.0008 4.078 97.8 79.9 0.347 0.356
10 78 0.00085 36580 33395 6597671 0.0008 4.078 97.8 79.9 0.347 0.356
10 79 0.00085 36580 33395 6597671 0.0008 4.078 97.8 79.9 0.349 0.450
10 78 0.00085 38580 43348 8857058 9111463 0.0038 5.873 97.7 79.8 0.406 0.349
12 80 0.00084 45404 43948 8857058 9111463 0.0038 5.883 97.7 79.8 0.406 0.449
                                                                                       94 U.00064 45404 45488 8857058 9111463 U.0038 5.873 97./ 80.0 U.484 D.467
         SKIN FHI COLFF AND WINDHAL ORIGIN FROM CURVE FIT TO MOM THI MEASUREMENTS
CURVE FIT THE LET 0.0000 N. MADE N. 9370 SVIR ORIEXD = 2.507 INCRE.
NO IN INCR. 2504 FACTOR INCH 1110 INCR. (FIT) SMOOTH
1 2 0.0302 0.0183 1.0331
2 0.0183 0.0042 1.6987
4 10 0.1575 0.0040 1.6987
4 10 0.1575 0.0040 1.6987
5 2 0.2252 0.1328 1.0351
5 2 0.2252 0.1328 1.0350
6 22 0.2252 0.1328 1.0350
6 22 0.2252 0.1328 1.0350
7 0.01528 0.0000 0.0012 0.0004 2.80 104.79
9 34 0.3134 0.1920 1.6550
1 40 0.4116 (2514 1.055)
4 4 0.3154 0.0021 0.0551
5 6 2 0.5531 0.5271 1.6396
6 62 0.5531 0.5271 1.6396
7 7 0.3855 0.0009 0.0009 0.0000 947.87
9 1.540 0.6061 0.3657 1.6574
7 2.70 0.3855 0.0009 0.0009 0.00000 947.87
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94 0.00202 30461 265331097067211232101 0.
                                                                                                                                                                                                                                                                              Ű.
                                                                                                                                                                                                                                                                                                               105.8 106.8 0.261 0.227
                                                                                     AND VIRTURAL ORIGIN FROM CURVE FIT TO MOM THE MEASUREMENTS
        SKIN FRE COLFF
    SKIN FRI COLFF AND VIRIDARL ORIGIN FROM CURVE FIT TO MOM THE MEASUREM CURVE FIT THERE 0.0052*(X=XO)** 0.47050 *VIR ORIGINO!* 2.2255 INCHE'S PL X LIS IN MOM IN SHAPE X=XO MOM IN DIFF CF2 CF2 RA NO IN INCH INCH FACTOR INCH (FIT) INCH (FIT) SMOOTH 4 14 0.0902 0.00562 1.6057 INCH (FIT) INCH (FIT) SMOOTH 4 0.0902 0.00524 0.00139 2 8 30 0.1574 0.1026 1.5543 52.25 0.1016 0.0003 0.00270 0.00119 2 4 60 0.2096 0.1027 1.4757 44.23 0.1036 0.0012 0.00270 0.00119 2 16 62 0.2680 0.1037 1.4590 0.423 0.1035 0.0004 0.000244 0.00103 2 2 0 78 0.3240 0.2218 1.4608 80.23 0.2218 0.0001 0.00237 0.00099 2
                                                                                                                                                                                                                                                                                                                                                                    RE NO
P.14 200-04
                                                                                                                                                                                                                                                                                                                                                                           2.26 190.44
2.31 185.01
STANTON RUN B/29/75 V=242 FT/5EC F=+001
VEL=25B+9 FT/5EC 7= 77.5 (EGF TO= 82.2 (RGF RHO=0.0755 EB5/FT3 P= 14.66 P5)
TOB= 78.0 OLGF TWE GF. 0 DEGF PAMS= 29.78 TR HG VIR OR46+X0= -0.61 TR
PL X STR MOM EGTH X (X=K0) 1 B TPL TAIR MOMIN EATTH
MO INCH NO REYNO REYNO REYNO REYNO REYNO
1 2 0.00300 1658 1156 236874 389120 0.0010 0.261 94.5 70.9 0.014 0.010
2 6 0.00300 3908 3216 7710621 782867 0.0010 0.350 99.5 78.6 0.035 0.027
3 10 0.00264 5817 6670 1658116 1750362 0.0010 0.371 99.5 78.5 0.055 0.027
4 14 0.00238 7817 6670 1658116 1750362 0.0010 0.418 99.6 78.4 0.066 0.066
5 10 0.00230 993 8250 2131663 2204104 0.0010 0.418 99.6 78.4 0.066 0.066
5 20 0.00231 1251 9794 2605610 2677857 0.0010 0.492 99.6 79.0 0.081 0.070
6 22 0.00221 1251 9794 2605610 2677857 0.0010 0.492 99.6 79.0 0.081 0.070
8 30 0.00207 14449 12781 3595105 562551 0.0010 0.492 99.5 79.6 0.109 0.055
8 30 0.00207 14449 12781 3595105 562551 0.0010 0.492 99.6 79.7 0.122 0.108
9 34 0.00203 16107 14218 4026652 4099049 0.0010 0.492 99.6 78.9 0.136 0.120
10 38 0.00195 17647 15630 4500999 4572046 0.0010 0.492 99.6 78.9 0.136 0.120
11 42 0.00191 17187 17075 4970347 5046593 0.0010 0.518 99.6 78.9 0.136 0.120
12 46 0.00109 2029 19796 5921041 5994088 0.0010 0.525 99.6 80.9 0.174 0.155
13 50 0.00190 2029 19796 5921041 5994088 0.0010 0.525 99.6 80.9 0.174 0.155
14 54 0.00186 2569 2160 659589 647859 0.0010 0.544 99.6 79.8 0.190 0.179
15 58 0.00178 26293 23825 7343083 7415330 0.0010 0.571 99.6 79.8 0.190 0.179
16 62 0.00178 39135 26485 829078 8362674 0.0010 0.571 99.6 79.7 0.258 0.259
17 74 0.00178 39135 26485 829078 8362674 0.0010 0.575 99.5 79.8 0.190 0.254
18 70 0.00178 39135 26485 829078 8362674 0.0010 0.575 99.5 79.8 0.190 0.254
19 74 0.00178 3057 27804 0764525 8856571 0.0010 0.565 99.5 78.9 0.270 0.298
20 78 0.00178 38135 3291 1005931410731501 0.0010 0.615 99.5 78.9 0.270 0.298
21 80 0.00160 36123 3291 1005931410731501 0.0010 0.615 99.5 78.6 0.202 0.207
22 86 0.00160 36123 3291 1005931410731501 0.0010 0.615 99.5 78.6 0.305 0.278
24 90 0.00160 36123 3291 1005931410731
      SKIR FRE COLFF AND VIRTURAL ORIGIN FROM CURVE FIT TO MON THE MEASUREMENTS
    PROFIT (+11) TRCH (+11) 5M001H

6+01 0+0358 -0+0001 0+00326 0+00124

10+01 0+0300 0+0001 0+00260 0+00096

50+61 0+1223 0+0004 0+00232 0+00082

46+01 0+1740 -0+0010 0+00210 0+00073

62+61 0+229 0+0007 0+00195 0+0006

78+61 0+2697 -0+0002 0+00185 0+00063
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62 0.3336 0.2221 1.5018 0.4028 0.2699 1.4923 2.82 174.69 2.88 166.20 2.91 160.33

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STANTON, NUM: 87/3973 V=282 FT/SLC F=.007
VEL=236:2 FT/SLC T= 78.7 DeGF TO= 83.5 DeGF NHO=0.0733 LBS/FT3 P= 14.65 P91
TODE 85.0 DEGF TH=6 86.0 LEGF PME=29.76 IN HG VIR ORIGEXO= -5.65 IN
PL X SIH AOM ERTH X (X=X0) F B TPL TAIR MORTH ERTTH
NO INCH NO HCYNO NEYMO REYMO REY
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VELOCITY PROFILES WITH BLOWING AT 32 FT/SEC

BLO FR 0 STAT TLMP 7 STAT PRES 14. PORT DIST 1	.40EGF 530PSIA 3 HINCH DP VVVINE	0.00 74.6 14.60 6	BUEGF UPSIA	0.00	List 61	0.00	LULGE	0.00 75.8	SETZSLO SE VOEGE
STAT TLMP 7 STAT PRES 14. PORT DIST 1	3.4DEGF 53DPSIA 3 FINCH	74 - 8 14 - 8 6	BUEGF UPSIA	75+ 14+76	List 61	68.	LULGE	75.8	
STAT PRES 14. PORT DIST 1	SUDPSIA HINCH	14+800 6	JF51A	14.76					UE GF
PORT DIST 1	FINCH B	6			JI'5IA	14.000	THE LA		
1 Y-YT	FINCH		JCH	4),) Y W	14.780F51A	
1 Y - Y	_	1411	JC.H		4		.5		
	N NATUR		141NCH		10111011		1411/11		ICH
		Y-YTOU	ANTAK	V-VIAU	V/VINF	V=V tAII	WZVIII	V_V100	u zu tra
	0.	0.	0.47146	U.	U.	0.	O. A. TITE	0.	0.
0.0		0.011	0.399	0.011	U.411	0.011	0.560	0.011	0.312
0.0		0.015	0.415	0.013	0.433	0.012	0.370	0.016	0.541
0.0	-	0.015	0.432	0.015	0.457	0.013	0.581	1.020	
0.0			0.465	0.015	0.470	0.013	0.386		0 • 575
J. U			0.496					0.046	0.457
		0.050		0.020	0.497	0.015	0.400	0.066	(1.4.15
U • U			0.519	0.025	0.519	0.016	0.405	6.091	0.542
0.0			0.537	0.050	0.540	0.017	0.415	0.116	0.573
V.U			0.553	0.035	0.565	0.018	0.425	0.141	0.60%
0.0			0.502	0.040	0.579	0.019	0.428	0+166	Uefiall
0.0			0.568	0.045	0.593	0.020	0.433	0.216	U.F.HB
U.U			0.604	0.055	0.619	0.022	11.442	0.296	0.727
V•0		0.075	0.625	0.005	0.636	0.0.4	0.454	0.306	0.415
0.0		0.090	0.646	0.075	0.654	0.026	0.405	il a faqsés	0.830
0.0			0.678	0.090	0.679	0.058	1.475	0.66	0.931
0.0			0.701	0 - 115	0.719	0.032	0.495	veibb	0.900
V • U			0.733	0.140	0.748	0.038	0.400	0.760	0.991
0.0			0.757	0.165	0.777	0.048	0.544	0.000	1.000
U • U		0.215	0./77	0.190	0.604	0.058	0.505		
U • U		0.265	0.817	0.215	0.828	0.008	0.586		
V• U			0.861	りゃりいり	u.87.3	0.078	9.608		
V • U		0.365	0.895	0.515	0.915	3500.0	نے کؤوں ۔ (ا		
0.0			0.9.24	U • 3(1)	(1.945	0.098	(1.,,41		
U • ()		0.465	0.450	0.415	0.964	0.118	0.004		
0-1			0.981	0.405	0.481	0.143	0.696		
U + 1		0.005	0.996	りゃりりり	() • 44()	0.168	0.730		
0.1		0.165	1.000	11+6465	1.000	0.193	0.757		
V • 1						0+218	0.781		
0.2						0.245	0.306		
ہے ۔ ن						ひゃどいれ	1) • 4> 61		
0 • ~						0.293	0.554		
0.3						0.518	0.475		
0.3						0 (364	0.911		
0.4	27 0.448					0.418	0.938		
0 • 4						0.468	0.960		
Uals	2/ 1.000	1				0.578	0.078		
						ひょういん	0.986		
						0.618	0.097		
						0.718	1.000		

VELOCITY PROFILES WITH BLOWING AT 90 FT/SEC

Y-YTOP V/VINF Y-YTOP V/VINF Y-YTOP V/VINF Y-YTOP V/VINF Y-YTOP V/VINF Y-YTOP V/VINF O. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0. 0.	F S VEL BLO FR STAT TEMP STAT PRES PORT DIST			0+001 DEGF 77+0DEGF PSIA 14+650PSIA 3		89./1FT/SEC 0.002 75.0DLGF 14.750PSIA 5 10THCH		88.82FT/SEC 0.004 76.4UEGF 14.76UPSIA 3 10THCH		86.22FT/SEC 0.008 76.3DEGF 14.770PSIA 2 10INCH	
0.0 0.0 0.446 0.011 0.458 0.011 0.429 0.011 0.375 0.011 0.374 0.013 0.495 0.013 0.496 0.013 0.497 0.013 0.496 0.013 0.497 0.013 0.496 0.013 0.497 0.015 0.492 0.015 0.492 0.015 0.490 0.015 0.497 0.015 0.496 0.015 0.497 0.017 0.468 0.017 0.494 0.017 0.464 0.017 0.411 0.017 0.417 0.020 0.482 0.019 0.505 0.019 0.474 0.019 0.418 0.19 0.424 0.025 0.501 0.021 0.517 0.021 0.484 0.021 0.432 0.024 0.449 0.035 0.555 0.026 0.538 0.025 0.496 0.026 0.449 0.029 0.473 0.035 0.555 0.026 0.538 0.025 0.496 0.026 0.449 0.029 0.473 0.035 0.557 0.031 0.500 0.027 0.514 0.036 0.487 0.034 0.487 0.095 0.558 0.564 0.041 0.593 0.032 0.550 0.031 0.472 0.034 0.487 0.055 0.564 0.564 0.041 0.593 0.032 0.550 0.044 0.519 0.551 0.055 0.588 0.046 0.605 0.037 0.550 0.044 0.521 0.049 0.551 0.055 0.600 0.605 0.606 0.566 0.566 0.566 0.566 0.566 0.566 0.566 0.566 0.566 0.566 0.566 0.566 0.566 0.566 0.566 0.566 0.566 0.566 0.567 0.029 0.576 0.029 0.051 0.482 0.059 0.576 0.042 0.591 0.070 0.598 0.150 0.071 0.669 0.057 0.600 0.066 0.567 0.029 0.551 0.094 0.552 0.095		V V+		V. VIAN DANIE		w w.c.		V V*611		U UTAL	
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0.013 0.450 0.013 0.468 0.013 0.468 0.013 0.487 0.013 0.308 0.013 0.302 0.015 0.462 0.015 0.462 0.015 0.4007 0.015 0.462 0.015 0.400 0.015 0.407 0.0020 0.482 0.019 0.505 0.019 0.474 0.019 0.418 0.017 0.444 0.017 0.464 0.017 0.441 0.017 0.441 0.017 0.424 0.025 0.501 0.021 0.517 0.021 0.484 0.021 0.432 0.024 0.429 0.035 0.553 0.026 0.558 0.023 0.496 0.026 0.449 0.029 0.473 0.035 0.557 0.031 0.500 0.025 0.500 0.031 0.496 0.027 0.514 0.036 0.487 0.039 0.507 0.045 0.558 0.004 0.552 0.056 0.576 0.027 0.514 0.036 0.487 0.039 0.507 0.045 0.588 0.046 0.605 0.037 0.534 0.046 0.521 0.049 0.551 0.005 0.605 0.600 0.056 0.566 0.056 0.566 0.056 0.566 0.056 0.556 0.057 0.659 0.051 0.462 0.659 0.651 0.602 0.604 0.551 0.602 0.604 0.551 0.602 0.604 0.551 0.602 0.604 0.551 0.604 0.551 0.602 0.604 0.551 0.604 0.605 0.007 0.606 0.566 0.056 0.05											
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0.075											
0.090											
0.115 0.697 0.081 0.692 0.097 0.606 0.086 0.010 0.094 0.622 0.140 0.734 0.091 0.710 0.100 0.067 0.631 0.096 0.627 0.119 0.663 0.160 0.100 0.645 0.144 0.706 0.190 0.215 0.835 0.161 0.830 0.097 0.691 0.151 0.665 0.169 0.742 0.215 0.835 0.161 0.830 0.097 0.690 0.156 0.720 0.194 0.782 0.265 0.888 0.160 0.866 0.107 0.707 0.260 0.787 0.249 0.815 0.335 0.971 0.211 0.896 0.117 0.727 0.206 0.787 0.269 0.886 0.461 0.850 0.925 0.127 0.743 0.231 0.813 0.314 0.314 0.326 0.925 0.127 0.745 0.231 0.815 0.369 0.369 0.369 0.372 0.385 0.996 0.445 0.990 0.172 0.885 0.331 0.911 0.419 0.986 0.426 0.386 0.996 0.127 0.885 0.331 0.911 0.419 0.989 0.386 0.996 0.127 0.885 0.331 0.911 0.419 0.989 0.396 0.386 0.996 0.127 0.885 0.381 0.946 0.409 0.996 0.436 0.347 0.922 0.876 0.431 0.997 0.409 0.996 0.397 0.397 0.397 0.397 0.397 0.998											
0.140 0.734 0.091 0.710 0.067 0.651 0.096 0.627 0.119 0.663 0.165 0.771 0.111 0.748 0.077 0.651 0.106 0.645 0.144 0.706 0.190 0.802 0.136 0.795 0.087 0.671 0.131 0.065 0.144 0.706 0.215 0.865 0.161 0.835 0.007 0.607 0.156 0.720 0.194 0.782 0.265 0.888 0.186 0.866 0.107 0.710 0.111 0.751 0.249 0.815 0.315 0.934 0.211 0.896 0.117 0.727 0.206 0.787 0.269 0.885 0.916 0.236 0.925 0.127 0.705 0.231 0.815 0.519 0.936 0.415 0.456 0.925 0.127 0.745 0.231 0.869 0.509 0.972 0.465 1.000 0.536 0.990 0.172 0.482 0.531 0.869 0.509 0.972 0.485 0.486 0.996 0.197 0.482 0.531 0.869 0.569 0.972 0.485 0.486 0.996 0.197 0.481 0.531 0.911 0.419 0.989 0.496 0.436 0.996 0.197 0.222 0.876 0.431 0.997 0.459 0.459 0.996 0.497 0.247 0.905 0.531 0.990 0.409 0.499 0.247 0.905 0.531 0.994 0.469 0.996 0.247 0.905 0.531 0.994 0.469 0.996 0.247 0.905 0.531 0.994 0.469 0.996 0.247 0.905 0.531 0.994 0.469 0.996 0.247 0.905 0.531 0.994 0.469 0.996 0.247 0.905 0.531 0.994 0.469 0.996 0.247 0.905 0.531 0.994 0.469 0.996 0.247 0.905 0.531 0.994 0.469 0.996 0.247 0.905 0.531 0.994 0.469 0.996 0.247 0.905 0.531 0.994 0.469 0.996 0.247 0.905 0.531 0.994 0.469 0.996 0.247 0.905 0.531 0.994 0.469 0.996 0.247 0.995 0.531 0.994 0.469 0.996 0.247 0.995 0.531 0.994 0.469 0.996 0.247 0.995 0.531 0.994 0.469 0.996 0.247 0.995 0.247 0.995 0.531 0.994 0.469 0.247 0.995											
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0.465 1.000 0.530 0.990 0.172 0.812 0.531 0.911 0.419 0.989 0.586 0.996 0.197 0.843 0.581 0.946 0.409 0.996 0.996 0.436 0.431 0.942 0.450 0.451 0.942 0.509 1.000 0.247 0.405 0.531 0.9494 0.551 0.9494 0.451 0.572 0.551 0.9494 0.572 0.573 0.573 0.998 0.008 0.0											
0.586 0.996 0.197 0.843 0.381 0.946 0.469 0.996 0.436 1.000 0.222 0.876 0.431 0.972 0.569 1.000 0.247 0.905 0.551 0.994 0.469 0.247 0.931 0.631 1.000 0.347 0.494 0.397 0.397 0.492											
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0.247 0.905 0.531 0.904 0.272 0.931 0.631 1.000 0.347 0.4979 0.347 0.492											
0.272 0.931 0.631 1.000 0.347 0.979 0.397 0.992											- · · ·
0.347 0.479 0.397 0.492											
0.447 1.000						0.397					
						0.447	1.000				

VELOCITY PROFILES WITH BLOWING AT 139 FT/SEC

F 5 VEL	139.71F1/5EC		137.03			FT/SEC		of 1/SEC
BLO FR			0.001		0.00		0.00	
STAT TEMP			75 • 90E GF		76.5bL6F		77 • 30EGF	
STAT PRES	14.6709518		14./UUP51A		14.700P5IA			
1 304	3		3		3		3	
0151	101468		101NCH		101000		10INCH	
	Y-YTON VAVINE		Y-Y10P	V/VIIIF	Y-YTOP	V/VIIIF	Y-YTOP	V/VIII
	0.	U.	0.	U.	0.	0.	U.	
	0.011	0.434	0.011	0.426	0.011	0.410	0.011	0.371
	0.013	0.454	0.013	0.444	0.013	0.419	0.015	0.582
	0.015	0.469	0.015	0.455	0.015	0.426	0.015	0.592
	0.017	0.480	0.017	0.466	0.017	0.445	0.000	0.417
	0.019	0.491	0.022	0.495	0.019	0.460	0.025	0.435
	0.021	0.500	0.027	0.516	0.021	0.467	0.030	0.458
	0.026	0.522	0.032	0.5.55	0.026	0.488	0.935	0.4/1
	0.031	0.544	0.037	0.552	0.031	0.507	0.045	0.501
	0.030	0.564	0.047	0.586	0.036	0.528	0.055	0.529
	0.041	0.578	0.057	0.612	0.041	0.547	0.065	0.548
	0.051	0.607	0.06/	0.636	0.051	0.577	0.065	0.590
	0.061	0.634	0.077	0.600	0.061	0.603	0.110	0.654
	0.086	0.693	0.092	0.692	0.071	0.629	0.135	0.675
	0.111	0.747	0.11/	0.739	0.081	0.651	0.160	0.709
	0.130	0.795	0.14	0.785		0.674	0.145	0.746
	0.161	0.839		0.826	0.111	0.714	0.210	0./7/
	0.211	0.917	0.192	0.861	0.130	0.759	0.235	908+0
	0.261	0.972			0.161	0.792	0.285	0.862
	0.311	0.996	0.267	0.952	0.180	0.424	0.355	0.910
	0.361	0.999	0.31/	0.985	0.211	0.855	0.385	0.950
	0.401	1.000	0.367			0.912	0.435	0.977
			0.417		0.511	0.962		0.998
					0.361	0.498	0.635	1.000
					0.411	1.000		

1	/FT0C11/	r profil	LES WITH	I HLOWIN	NLOWING AT 190 FT/SEC					
F S VEL		5F T/5EC		H TVSEC		IF I/SEC	190.4	SETISEC		
BLO FR	C •		0.00		0.00		0.00)4		
STAT TEMP	78.10EGF		76 + 6 D L G F		77 - 3ULGF		/H • ODE OF			
STAT PRES	14.620P5IA		14.6/0P5IA		14.670PSIA		14.660P5IA			
PORT	2		4		4		4			
DIST	TOTHCH		14 INCH		14 INCH		14114011			
			Y-Y TOP		Y-Y101'		Y-YTOP			
	U.	0.	0 •	0.	U.	U •	U •	υ.		
	0.011	0.494	0.011	0.410	0.011	0.392	0.011	0.332		
	0.010	0.529	0.013	0.420	0.013	0.402	0.013	0.354		
	0.021	0.557	0.015	0.432	0.015	0.411	0.615	0.349		
	0.026	0.580	0.017	0.438	0.017	U.419	0.017	0.357		
	0.051	0.603	0.022	0.463	0.019	0.428	0.019	0 • 366		
	0.036	0.621	0.027	0.480	0.021	0.457	0.021	0.571		
	0.046	0.658	0.032	0.495	0.023	0.445	0.023	0.380		
	0.056	0.692	0.037	0.511	0.025	0.450	0.025	0.387		
	0.066	0.724	0.042	0.523	0.027	0.458	0.027	0 • 595		
	0.076	0.755	0.047	0.535	0.032	0.471	0.032	0.407		
	0.086	0.783	0.057	0.558	0.037	0.487	0.037	0.424		
	U • 10t	0.857	0.067	0.578	0.042	U.498	0.042	0.435		
	0.131	0.896	0.077	0.597	0.047	0.511	0.047	0.448		
	0.156	9.944	0.087	0.616	0.052	0.522	0.052	0.457		
	0.181	0.9//	0.097	0.634	0.057	0.5.55	0.057	0.464		
	0.200	0.993	0.122	0.673	0.067	ひゅうちょ	0.067	0.487		
	0.231	0.999	0.147	0.711	0.077	0.572	0.077	0.506		
	0.250	1.000	0.172	0.746	0.037	0.588	0.087	0.521		
			0.197	0.779	0.097	0.603	0.097	0.538		
			0.222	0.812	0.107	V. 019	0.107	0.554		
			0.272	0.870	0.132	U+654	0.132	0.585		
			0.322	0.925	0.157	4.041	0.157	0.618		
			0.572	0.964	0.182	0.724	0.182	0.648		
			0.422	0.989	0.207	0.758	0.207	0.080		
			0.522	1.000	0.232	0.788	0.232	0.708		
					0.282	0.846	0.242	0.701		
					0.332	0.897	0.332	0.814		
					0.382	0.944	0.382	0.364		
					0.432	0.977	0.432	0.909		
					0.482	0.993	0.482	0.948		
					0.582	1.000	0.582	0.990		
							0.602	1.000		

+ 5 VLL	VELOCITY	Y PHOFII	LES WITE	4 BLOWIN		42 FT/SE	
BLO FR	0.	EF 17 3EC				PET/SEC	
STAT TEMP		SPEGF	0.00	JULGF	0.002		
STAT PHLS		A1516		11'5IA	/B . /ULGF		
PORT	4	Jr 446	24.000	31 31×	14.650P5IA		
บไร้เ		uc H		1611	21NCH		
••••	SINCH		• • •		4 417611		
	401Y-Y	V/VIIN-	Y-YTOP	V/VINE	Y-Y10F	V/VINE	
	0.	0.	U•	0.	U.	u.	
		0.423		1).464		0.372	
		0.451				0.372	
		0.444		0.527	U.015	0.363	
		U.471	0.026			0.369	
		0,495				0.374	
	0.030		0.046	0.649	0.021	0.389	
	0.035		0+071	0.723			
	0.045	U.564	0.096			0.437	
	11.055	0.590	0.121	0.847		0.454	
	0.065	0.612	0.146	0.900		0.471	
	0.080	0.642	0 - 196			0.500	
	0.105	0.687			0.061	0.524	
	0.130	0.724			0.086	0.572	
	0.155	0.166	0.340	1.000	0.111	0.015	
	0.180	0.802			0 - 130	0.649	
	りっぱひり	0.835			0.161	186.0	
	0.255	0.896			0.211	0.742	
		U.946			0.201	0.605	
		0.980			0.511	U.HSO	
		0.995			0.361	0.906	
		0.999			0.4411	0.975	
	ひ・ちちゃ	1.000			0.561	0.447	
					9.061	0.999	
					1.761	1.000	

APPENDIX F

LISTING OF DATA REDUCTION PROGRAM

The following pages contain a source listing of the data reduction program used to reduce the Stanton number data. The program is written in Fortran IV and employs fixed-field input. Output includes the reduced data and a listing of the raw input data as well. Input requirements and card formats are documented with comment cards in the listing of the DATIN subroutine of the code.

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| A1A REDUCTION PROGRAM FOR BULGHNESS RIG STANION NO & ENERGY BAL PUSS DIMENSION CONNICADA RI(24) TCAST(24) TTIN(24) BITN(24) CLLST(15) CONNON/DAI/TAIR PSTATISG PBOXITPL(24) TCS(12) STR(24) CFM(24) CFM(24) FIAR (24) TCFM(24) TTINF TEXT TO THE RESULUTION OF THE RES
  12.984.2.990.2.984.2.986.2.969.2.965.2.967.3.015.2.980.2.957.2.984.
  22.966/2.956/3.023/3.020/
     CONSTANTS FOR CORRECTION OF AIR INLEY TEMP FOR CASTING TEMP MISMATCH
LATA ATIN/.30.,21.,16.,14.,30.,26.,36.,20.,29.,24.,38.,16.,47.,36.
1.27.,48.,25.,17.,36.,26.,18.,32.,20.,24.

LATA BTIN/.097.,150.,064.,194.,055.,637.,025.,058.,040.,060.,121.
1.076.,076.,040.,029.,078.,038.,097.,030.,000.,026.,343.,050.,020/
     IHROR=-1
CALL DATAIN
     RE IS ABS HUMIDITY: LBS WATER/LB AIR
RELIN 0196919-.000679526+TWB+8.91667L-n6+TWB+*2-.000232+(1DB-TWB)
     PHES-PAMB+850.74/12.

DENCOR: --0075-(TDB-60.)

IF:(K1.g1.5).AND.(K1.L1.6))PSTAT=0.

PST=PSTAT/12.+62.3+DENCOR
     PS:=PS:ATF 22.*B2:3*DENCOR
[H:((K1.EQ.3).OR.(K1.EQ.6))PBOX=0.
PBX=PBOX/12.*B2.3*DENCOR

CP=0.241*(1.+.9*RH)

LOI 15 A CO.STANT USED IN CALC OF ELECT POWER

COI=3.415.3500.*(144./148.)
      PH4=0.71**0.4
     DU 9 J=1,4
      1 = 6 (U-1) + 1
IF ((K) = NE - 3) = AND = (K1 = LT - 6) ) CALL EMFT (TAMB(U) = TANBT (U) + 513)
      1: ((K1.EG.3).OR.(K1.EG.6))TAMBT(U)=TOB
CAS1=(TCS(N+1)+TCS(N))/2.
     CASETICS(N+1)+TCS(N+2))/2.

CASETICS(N)+TTCS(N)-CASI)/2.

CALL EMPT(CASS, TCAST(1):$63)

CASETICS(N)+CASI)/2.

CALL EMPT(CASH, TCAST(1+1):$63)
      CASS=(CASI+1CS(N+1))/2.
CALL EMFT(CAS5,TCAST(I+2):$63)
CASG=(TCS(N+1)+CAS2)/2.
     CASETICS(N+1)+CAS2//2.

CALL EMFT(CAS6, TCAST(1+3)+363)

CAS7=T(CAS2+TCS(11+2))/2.

CALL EMFT(CAS7, TCAST(1+4)+363)

CAS6=TCS(N+2)+(TCS(11+2)+CAS21/2.
     CASE (510+2)*(165(17*2)*CASE (72)*CASE (72)*CASE (72)*S63)
CALC OF FREE STREAM TEMP FROM 1C REANING
CALL EMFT(TAIR, TIMF, $65)
KHOIHF=(PRES+PST)/(53, 33*(TIHF, 459, 6))*(1.**, 20*RH)
       IF ((K1.GT.3).AND.(K1.L1.6)) GO TO 12
       ALFE.86
       RCF: 86
VAIH=5GRT(VAR=5G=32.17=62.467(L.=RHOINF))
VELHO=VAIR=7AIR/(778.664.349CP)
CURRECTION OF FREE STREAM TEMP FOR VIL HEAD
TZENO=TINF+(1.=RCF)=VELHO
       TIM =TINF-RCF+VELHD
HHUGLD=RHOINF
       HINGTH = (PRES+PST)/(53.33+(TINF+459.6))+(1.-.26+RH)
       PHILINE (PRESPYSI)/(33.33811119+45946)/
PSI=(PRESPYSI)/(44.
VISINF=(11.0+.0175+TINF)/(1.66+RHOINF)
CURSIANT FOR REYNOLDS NO CALC
VAIH=VAIH+SORT(RHOOLD/HHOINF)
CUNS=VAIH/VISINF
       CUNSTANT FOR STANTON NO CALC

LUZZO-SOCPORHOIMFOVAIR

IF((IOUT.EG.O),AND.(IRNOR.LE.O))CALL HEAD

IF((IOUT.GT.O),AND.(IRNOR.LE.O))CALL MEAD2(IOUT.X22:(0:512)
       SUM=0.
FL=1./3.
       + LG=1./6.
+ LGO=XZERO/12.
        DO 5 N=1+24
PLATE POWER CALCULATION
       PHOLOGOUSES (N) EH(N)/HI(N)
CALL EMPT(TAR(N))T, $63)
CALL EMPT(TPL(N))TPLATE:$63)
         IF (K1.E0.3150 TO 501
```

```
CALC OF TRANS FLOW FROM FLO METER EMERKICI BLUM, KIES SUCK
                     CALL EMFT(TCFM(N), TT, $63)

RHO=(PRES+PBX)/(53.33+(TT+459+6))+(1.-.26+RH)
                    HNO=TPRES+PRY/(53.33*(TT+459.6))*(1.-.26*RH)

CALL FLOM(PAMB,TT,RH,RHO,K1:N,CFM(N),ACFM,PLFLO)

CONRECTION OF AIR TEMP FOR CAST TO AIR TEMP MISMATCH

1P(-1.LE.9)TTAM=.75*TAMBT(1)+.25*TAMBT(2)

1P((N.GT.9).AND.(N.LT.16))TTAM=(TAMB1(2)+IAMBI(3))/2.

1P(N.GE.16)TTAM=.75*TAMBT(3)+.25*TAM-T(4)

TH=TT+(TTAM-TT)*(1.-EXP(-1.175/ACFM))

TH=T+(ATIN(1)+RTIN(1))*(1:N-TCAST(N))
                    T=1+(ATIN(H)-BTIN(H)+(TIN-TCAST(N)))
RAD FROM FRONT OF PLATES
                 TH((TENBAL, EG. 1), AND. (KI.NE.3)) GO TO 41
RCONTU. 163
THANF4=((TZERO+459,6)/100.)**4
GO TO 42
     41 HC0020.199
                    TRADE4=((TDU+459.6)/100.)**4
 TRADE 4=((TDD+459.6)/100.) ***

WERNITERCON**171**(((TPLATE+459.6)/100.) ***4-TRADE4)/3600.

CONDCTION LOSSES FROM PLATES. COND(I) FROM EXPLR.

GUDGE=COND(I)*(TPLATE-ICAST(II))/1000.

TLAS#=((TCAST(I)*459.6)/100.) ***

TAJN4=((T+459.6)/100.) ***

IF(x1*LT.3)*00 TO 21

K1=5 NO TRADE

TPLN=TPLATE+3600.**PPDEN/(40.*1..)

TPLN=TPLATE+3600.**PPDEN/(40.*1..)

TPLN=TPLATE+459.6)/100.) ***

TRADE TRADE

TRADE TRADE

TRADE TRADE

TRADE TRADE

TRADE TRADE TRADE TRADE TRADE TRADE TRADE

TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE TRADE
                 JUACK=GHACK/3600.
                   ELDIAVED.
                  ACFMED.
                   AAA=U.
                   PLILO-U.
                 TIEU.
BLUF = U.
                   BLL (N) =0.
                   60 TO 23
  21 FEXPEPEFEC+CP+3680.7(40.+12.)
                 11 (KI+LT.2)00 TO 22
                   KIEL SUCKING
               TPLO=IPLATE +PRIEN/(PLFLO+CF)+(1.-ExP(=FEXF))
TPLO=E(TPLO+GF,0)/100.)++4
HAD FHOM BACK FOR SHOCK HC AT HIATE ALK JENN
UMACK=.1714+(.1815+TPLB4-.1455+TCA54-.0360+TAIN4)/3366.
ECGN=CP+PLFLO+(TPLB+TINF)
                60 16 23
K1=1 BLOwING

22 TFLO=TPLATE=PROJEN/(PLFLOWCF)+(1.-ExP(FLXP))
               TPELISC (TPEJSENSO) / 100-15-64

FAU FROM BACK FOR BLOWS HC AT AIR TE P

404CF=1714-(-00076FPLSES-1358-TAIR4--1455-TCA24)/3000-
EUO:VECHAPLELOA (TELATE -T)
25 IF (APACKALTAGA) UBA AEGA
               ULUSEUFR'IT+ BACK+GCOND
                WILL AT SPHULLI-OLUS
              GREATHERDITARDLES

GREATHERDITARDLES

IF (1.6T.1) STOLESSTR(N-1)+BEOF

STATION 'NO LALC+ BASED ON STAG TO MALE TEMP DIFF

STREETERS TO MALE TEMP DIFF

STREETERS

FOR THE STREETERS

FOR
                IF (KI-EU.3) 60 TO 16
HEOF=ACEMONHOV(60.0.50VAIRONHOIH)
            HLE (1.) = BLOF / STY(N)

15 (1.) = BLOF / STY(N)

15 (1.) = STY(N) + HLOF)

SUMECONS + ENTTH/12.
60 To 3
10 Esta=(51'((2)+BEOF+510Eu)+4.
              Ent THEE INTTHACE THIZ.
                SUN = SUM+CONSTEL THY24.
     FLUZFLG+FL
3 FLGT=FLG=FLGO
            FLGT=FLG=FLGU

1+ (FLGT=LGT=CONS

1+ (1PHOH+GT=0)GU T( 6

IF ((H1-NE-3)-OH-(HEX(N)-CU-U-))GO TO 69
              TIAGELIABEL
STOLITABLESTNOWN
                HEXU(ITAB)=HEX(1)
69 THE 40=F MONT (11) +CONS/12.+0.5
                INEXATIONS + LG+U+5
               IHEAN=HEX (11) +0.5
                INCLINESUMOU.S
```

```
1+60=+60+12.+0.5
       IF(REX(N)+_0+0+7 GO TO 74
51_=+02957(PR4+REX(N)++0+2)
       IF(K1.HE.3)STS=HLOF/(EXP(HEOF/STS)=1.)
STR(N)=STS/((TPLATE+459.0)/(TZLRO+459.0))+*.4
        60 10 15
   74 51,20.
        STRUMEN.
   75 CONTINUE
        IF (10U1.GT.0) will te (G.4) h. Ir Lo. Sh(ii) . HeMO. IRELN. IRE XX. IREXI. BLOF.
      LIBEL (N) . TPLATE + T . F MUMT (II) + LIST TH
    # FURMAT (13+14+F3+7+211+217+4+F0+3+2F0+1+2F6+3)

IF (1001+Lu, 0) and IL (6+5) N+5 Tr(0) + F1 Tr++1REF1+F40MT(N) + IREMO+IFLU+IR

IE XX+FLUF+IREXII+BLUF+BEL(G)+TT+TCAST(G)+TPLATF+TPLDEF1+ECOAV+RLUS+R
     5 ( OKMA1 (13+) 9.5+6.3+16.1+6.3+16.14+18+6.2+18+6.5+6.5+6.3+46.6+1.3(8.5
      1 . 14)
        60 10 u
   19 PERCEGATIVAN AT *100*
STAREGATIV(25.*.5*CP*RHOIDF*VAIR)*10000*
DECMEGATIV(CF*(TPLATE=T)**075/60*)
      TECHROR-GERIDO TO G
TECHROR-GERIDO TO G
TECHDUTALO, DEMITECTO TO TO THE CONVAGATIAPERCAXXXASCEMADECMATTATP
TLATEATATOASTUDA OF RULA ORACKA OCONPATINASTUR
       FURNAT (15+3F10+6+8+3+3F7+3+4F7-1+3F8+5+F6+1+F6+2)
- 1F(100T+GT+0) WRITE (6+2)N+FMDER+ECONV+GEOS+GRET+PERC+STREK+ACFM+TPE
        TELLIOUT.ct. 1 PPUNCE 2 NO PROFILE CONVENIENCE OF TERC & STREET ACEMS TO LA
     . 1 United (130) H. S. &t 9. S. F. B. S. F. 7. 20 F. B. S. 21 /. 10 14)
     6 COLLI TRUE
    US IF ((IUUT-E 1-U) - AND - (IRROK-LF - 0)) CALL DATAOT
       CHICK TO SEE IF ERROR ANALYSIS IS ASK FORFIER=0+HG TER=1+YES IF (TER-E0+HGO) TO 104
    13 CULTINUE
        1F (1RRQR+E,,+035111#5TN(1,)
        IF (IRRUR.GT.0) DELST (IRROF) = 5131 - 514(12)
        IF ((K1.E3.3) - AND - (IRROR-OL - 8)) GO TO 100
        IF ((K1.11.3).And.114808.31.12))00 TO 100
  60 10 24
100 50ml K-0+
       110 101 1=1.1KHOH
  101 30 ACH= 30ME. +DECS (11) +CEC3 (1)
        SULLE HESIGRE (SUMER)
        was It (me II)
        IF CIRRORALGARIWREIT (6.4102)STITE (DELISECT) + 1=1.4 IRROREASUMER
        IF (IRROR-E. . 12) WRITE (6.103) STITE (ULL'ST(I) + I=1+IRPOR) + SUMER
  18 FURMATIVE UNCERTAINTY ANALYST'S SUMMART FOR PLATE 1217
                    10x+*PLAIL (TANTOL NUMBER = *++20+5+/
10x+* (0.5 b) (0.1(i) = *++20+5+/
10x+* (-0.1 MV Tin) = *++20+5+/
                    *1UX**UNLERTAINTY ALALYSI'S SUMMARY FOR PLATE 12*/
1UX**PLATE STATION NUMBER =**(2U*)*/
1UX** (0*) DEGF THE =***EU*)*/
1UX** (-U1 by TIN) =**EU*)*/
   103 FURMALL/
       .5
                                     (.0) by Tip) = "*F_0.50*;
(.0) #1920 VAI()="*E20.50*;
(.0) My TPLATE) = "*E20.50*;
(.0) My TAIR) = "*E20.50*;
(.0) VOL1 EPD) = "*E70.50*;
                     lux.
                     10X+*
                     1 u X + 1
                     164.
                                     164.
                     lux."
                     10.4
                                     TOPS TO FLORIKIET . E 20.5 . /
                     lux.
                                     (u.1 And Prox) = **F20-5*/
                                     LAUL AL TAHILL
                     Luxet
                     Lux.
                              STANTON NUMBER ERROR= ** EZU-5)
   104 IRRORE-1
        CHECK FOR HORE CASES IN SEGUENCE FIREXT-1 GO BACK & READ SOME MORE
         IF INHEAT of Galligo To 24
         WK11E(0+11)
    11 FORMATTINES
         4104
        لياس
        IN SUBROUTINE DATAIN FOR ST DATA REDUCTION PROGRAM SUBROUTINE PATAIN THIS SUBROUTINE READS IN DATA FROM THE ROUGHRESS RIG
C DATAIN
        CUMMUNUAT/TATE PSTATE SCEPHUNETPL (24) - 105(12) - 518(24) - CF4(24)
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```
1.TAK(24).TCFM(24).EH(24).ES(24).KNEXT.MOMEC.XZERU.NST.FMOMT(24).
2510(24).REXO(24).ITAB.TAMB(4).IOUT,IFR.IRNOR
COMMON/AMB/PAMG.TOB.TWB.DENT(18).K1.IENBAL.VAK
                 IHROR=IRROK+1
IF(IRROR-EG-0)GO TO 100
                  VARY INPUTS WITHIN UNCERTAINTY INTERVALS FOR ENHOR ANALYSIS IF (IPROR.GT.1)GO TO 1
         TUB=TUB+0.5
GO TO 54
1 IF (IPHOH.GT.2)GO TO 2
         TULETUB-0.5

TAIRETAIR+.01

GU TO 54

2 IF (IRHOR.GT.3)GU .0 3

TAIRETAIR+.01

VAREVAR+.002
         00 TO 54
3 IF (IRHOR-GT.4160 TO 4
               VAREVAR-.002
TrL(12)=:PL(12)+.01
                60 10 54
       GU TO 54

4 IF (1RHON.GT.5)GD TO 5

TPL(12)=TPL(12)=.01

TAH(12)=TAH(12)+.01

GU TO 54

5 IF (1RHON.GT.6)GD TO 6
               TAR(12)=TAR(12)-.01
EH(12)=EH(12)+.001
        GO TO 54
6 IF (INROH.GT.7)GO TO 7
               En(12)=EH(12)-.001
E5(12)=E5(12)-.01
                GO TO SE
      GU TO 54

1 F (1 F HOR.GT.8) GU TO 8

E 5 (12) = E 5 (12) + 01

T C 5 (6) = T C 5 (6) + .01

GU TO 54

8 IF (1 R HOH.GT.9) GO TO 9

T C 5 (6) = T C 5 (6) - .01

T C 5 (12) = T C 5 M (12) + .01

GU TO 54
       GU 10 54

9 I) (IRHOH.GT.10)GO TO 10

TUFH(12)=TUFH(12)=.01
              CFM(12)=CFM(12)+.025
   CFM(12)=CFM(12)+.025

GO TO 54

10 Ir-(1RHON.GT.11)GO TO 11

CFM(12)=CFM(12)-.025

TAMB(3)=TAMB(3)+.01
   PUUX=PBOX+.1
              60 TO 54
TOO CONTINUE
            HEGIN READING INPUT HERE, ALL IMPUT LIXED FIELD FOR AT AS INCLEATED MELOW READ IMPUT, FIRST CARD HAS 72 COL OF TITLE AND COL 30 IS AL IMPLICATOR IF ALS STANTON RUN WITH BLUWING IF KISS STANTON RUN WITH SUCTION IF KISS STANTON RUN WITHOUT TRANSPIRATION IF KISS STANTON RUN WITHOUT TRANSPIRATION IF KISS STANTON RUN WITHOUT TRANSPIRATION IF KISS SUCTION ENERGY BALANCE RUN IF KISS SUCTION ENERGY BALANCE RUN IF KISS NO WLOW EMERGY HALANCE RUN
              HEAD (5,25) UENT . KI
  25 FURMAT(1844,7X,11)
IF(K1-LT,3) GO TO 61
11AH=U
             DU 57 1=1.24
51H(I)=0.
 510(1)=0.
57 HEXU(1)=0.
          SIGNIFICAL

REQUITION

IF KIET OR 2, READ TOR, TWB, PAME, TAIP, PSTAT, VAR, SG, PEOX, UFTO FORMAT

IF KIET OR 2, READ TOB, TWB, PAME, TAIR, PSTAT, VAR, SG, PEOX

IF KIET OR 6 READ TOB, TWB, PAME, TAIR, PSTAT, VAR, SG, PEOX

IF KIET OR 7 READ TOB, TWB, PAME, TAIR, PSTAT, VAR, SG

IF KIET OR 5 READ TOB, TWB, PAME, TAIR, PAME, TAIR, PSTAT, VAR, SG

IF (KI, GT, 3), AND, (KI, LT, 6) THEAD(5, 29) TOH, TWB, PAME, TAIR, PBOX, VAR

TWB.ET BULD TEMP IN DEGF

TOUBLAY BULB TEMP IN DESF

PAMETAMB PRISSURE IN INCHES OF HG

TAIRETHEE STREAM TEMP IN MY FOR STANTON KUNS AND AND TEMP IN MY FUR

ENERGY BALANCE RUNS

PSTATETUNNEL-TO-AMB PRESSURE DIFFERE, CE IN INCHES OF FLUID OF SPECIFIC

GRAVITY SG FOR STANTON RUNS, AND VEL TO BE USED IN STANTON NO ERROR
```

```
EVAL FOR ENERGY HAL RUPP, IN FIZSEC 504 MANO FEBILD FOR VAR READING
          PROXETRAN'S HEADER BOX PRESSURE IN INCHES OF HED
          FURMAL (MF 10.0)
          DU 26 121.4
    FOR KI 3 OR 6 READ TPL+TAR+EH+E5+ 4F, U FORMAT

IF (KI+IU-S)+OR+(KI+I',6) HEAD TO +271TPL(I)+TAR(I)+E+(I)+E5(I)

FOR KI-I+244+5 HEAD TO +3TAR+IC+M+CFM+EM+ES+ 6+ID FORMAT

20 IF (KI+HL+D)+AND+(KI+L)+HIAD(5+27)TPL(I)+TAR(I)+TC+M+(I)+C+M+TI+
         11111111115611
          TPL (I) PLATE TEMP IN MY
          TARCIFETHALIS ALE ILMPS IN MY
          THE MICHAEL METER TEMPS IN MY CMF (1) = TOWN METER SIGNAL IN MY EREIT= METER SIGNAL IN MY EREIT= METER VOLTAGE IN VOLTS
           SCITETHURE VOLTAGE IN MILLIVOLTS
         FURMAT BELU-01
READ IN TUSCED CASEING TEMPS IN MV+ GEID FURMAT
          HUAD(5:28) (TCS(1):1=1:12)
          READ DATAMECTALISTINES IN MY FOR BLOWING RUNG OFFIA FAS FORMAT
TECHNICLES AND (KILLIST) HEAD DERROTTED AND TILLED
          FURMATION TO U.U.
         10 56 121124
FMOMI(1)=0.
         FMUMITIES.

READ IN KNEXT-MOMER-XZ/ROSN'ST-FOUTSTER- 2(TE-9X)-FID-3(9X-11) FORMAT KNEXT-FOR MORE CASES TO FOLLOW-O IF LAST CASE

MUS-CK-FI IP MOM THICK/ESSES ARE TO BE READ IN TO CALC MOM THI REY NO MAI ROSK EDUATION OF THREE RE VIRTUAL ORIGIN FOR (X-X0) REY 100

NSI-FI IF NO BLOW STANION NUMBERS ARE TO BE USED TO CALC STATION FETHIS OPTION TO USED A NO BLOW STANION NO ROW MUST THE EXECUTED FIRST TORREST FOR DITAILED OUTPUTS A FOR SHORT OUT WACARDS IN STANION NO UBLICITATION FOR STANION ROWSELFORD FOR SASTREES PLATE IS ONLY
          READ CO. STIKNERT MOMCK . AZERO . MIT FOUT . TER
          FORMAT ([] + 0X + 11 + 9X + F10 + 0 + 9X + 11 + 9X + 11 + 9X + 11 )
          HEAD IN FMOMISE MOMENTUM THE FOR LACH PLATE, 1206 FORMAT
          REDICTIONS (FROM) (1) +. 21 - 241
     TO FURMATTIZE GOLD
     ba Courthus
          RETURN
L MLAU
             A BRIDGE
                                HEAD FOR ST DATA REDUCTION PROGRAM
          THIS SUBMOUTING PRIMITS OUT OBJUTUT HEADING COMMOD/HD/TAMET COJ-VAIR+TIH+*TZERO+RHOINE+PSI
          CUMMUNIAMBINAMO TOD THE DELIT (18) . KI . I MAL . VAR
       WRITE (078) DENT
B FURNELLHISTAM
           IF ( (K1.G1.3) - AHD - (K1.L1.6)) GO TO 12
          TE (K1+LQ+0)G0 TO 45
WHITE (G+45) VAIR+T19E+T2ERO
     40 FURMATCH + STANDON NUMBER AT FREE STREAM VICE**FG-2*

1 FI/SEC TINF=*SFG-1** DEGF TOINF=**FG-1** DEGF
                                                                          TOTAL TIPE TOTAL
           ILIJIAL TU
           60 10 46
     45 WRITE CO. 441 VAIR+TIRE+12CRO
     44 FORMATCHE . NO TRANS ENERGY BAL RUN AT FREE STREAM VEL: 1.Fb.2.
         11 F1/5LC
                                11MF= 1+6.1+1 DEGF
                                                                          1011#="+| 6.1. DEGF")
          ILIJBAL = 1
           K1=3
     46 WRITE COMMING THE STATE OF THE PARK
4 FORMATCH REGINE = ** FOR 4 ** LECKETS PRES
1 ** TUDE = ** G.S.* DEGE TWG = ** FC.S.* DEGE
                                                                                 PRESINCETARGORAL DSI
                                                                                             PAMESTIFE 211 FIGT
          60 10 14
   12 WRITE COVERTING A DELICATION OF PARTY
       2 FORMATCHE * PREPSY HALANCE RUN TIME ** FEB. 1 ** DEGF*/
1* TOUT** FB. 1 * * DEGF PAMET** FB. 1 ** DEGF PAMET** FB.
                                                                                           PAMBETIFE.211 ANG!)
          ILIMAL#1
           1F(K1.LG.4) K1=1
           II (K1 a) 0 a 5 1 K1 5 2
     14 WPITE(0+47)(TAMBT(1)+T=1+4)
47 FORMAT(* TAMBT=+F6=1+* 04+6
        # WPITE(0)47)(TAMBT17); E194)

FORMAT(* TAMBET*): DLGF TAMBET*): DLGF TAMBET*): TOLGF TAMBET*): DLGF TAMBET*): DLGF TAMBET*): TOLGF*)

IF(IENDAL: 0:0)WRITE(0:40)

FORMAT(* PL STANTON LITTE L* H MOMTH MOM X X X=X0
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SUBROUTINE FLOW FOR ST DATA REDUCTION PROGRAM
    SUBROUTINE FEOREFIRMS TO PROBLEMS OF A COMOPLE CONDITION (5-2-24) - 565-2-24)
    CUMMONVAFEU/XXXX+SEM
FEGN METER CALIBRATIONS+INDATED B/24//3
    [JATA(CO (1-J-K)+)=1+5)+J=1+7+K=1+12)

76-514+4-769+2-866+1-846+ -721+6-519+4-764+2-866+1-846+ -721 8775- 1

+5-400+4-362+2-598+1-759+ -759+5-490+4-362+2-598+1-759+ -759 8775- 2
            +5-544+4-842+2-826+1.700+ .577+5-544+4.842+2-842+5-790+ .577 8775- 5
+7-194+5-8000+5-055+2-117+ .955+4-965+5-965+2-986+1-115+ .887 8775- 4
            +0.500+5.520+5.188+2.154+1.112+6.131+4.850+1.023+1.424+1.170 HZ75+5
+5.875-3.675+2.933+1.990+ .871+5.700+4.715+5.105+1.674+ .860 HZ75+ 0
            *4**U0***B97*1.477*1.573*1.180*4.420*2.988*1.990*1.577*1.204 4775- 7
*5.794*4.687*2.765*1.853* .813*5.781*4.687*3.033*1.000* .872 8775- 8
*6.044*4.893*2.942*2.026* .831*0.044*4.84*2.980*1.542* .867 8775- 9
            +6.738+5.112+2.946+1.+75+1.004+6.002+4.719.3.810+2+638+ .982 8/75=10
+6.324+5.046+3.107+2.151+1.075+6.125+4.850+3.119+1+677+ .980 8/75=11
             .4.820.3.042.2.05/1.501.1.105.4.826.3.141.2.084.1.661.1.756 8//3-12
            /5.190/3.073/2.334/1.505/.054/4.531/3.204/2.491/1.059/.741 8/73-13
/5.130/2.090/1.440/1.125/.838/3.022/2.044/1.400/1.095/.484 8/73-14
            +3.700+2.740+1.772+1.206+ .957+3.700+2.670+1.757+1.170+ .955 8/75-15
+3.896+2.765+1.915+1.527+1.194+5.950+2.675+1.355+1.498+1.155 8/75-16
             .4.766.3.253.2.263.1.789.1.404.4.740.3.214.2.229.1.782.1.366 H/73-17
            +6.948+4.700+5.080+2.190+1.056+6.000+4.700+5.080+2.190+1.050-8/75-1P
             .5.900+2.673+1.495+1.292+1.076+4.096+2.783+1.880+1.464+1.105 8/75-14
            .5.044.4.030.2.521.1.748. .806.5.069.4.112.2.684.1.442. .863 8//3~24
            1.02 0.3/14.95.23.97.41.90 1.83 6.10.15.26.24.58.41.64 8/73- 7 . .54 1.40 0.01.14.59.70.93 .54 1.57 6.10.23.91.80.52 8/73- 8 .62 1.53 6.17.24.56.73.39 8/73- 6
                                                       .60. 1.59. 3.16. 6.38.47.55 8//3-16
.62. 1.56. 5.99.24.04.64.22 8//3-11
                .59: 1.56: 6.20:13.90:46.86:
.58: 1.64: 6.17:14.90:64.05:
              1./4, 6.35:14./5:24.13:41.82: 1.74: 6.02:15:25:24.12:41:08 8//3-1:
     JAIA((($(1+J+K)+1=1+5)+J=1+2)+K=13+24)
            * .70* 1.89* 6.42*15.06*23.67* .70* 1.90* (.30*15.15*24*08 8/73-15* 1.88* 5.93*14.85*23.99*41.05* 1.91* 0.40*15*21*23.88*41*55 8/75-16*
             2.22, 6.24,14.87,25.00,41.82, 2.22, 6.15,15.32,24.55,41.49 8//3-20
                                                       .60. 1.65. 6.30.24.27./7.16 8//3-21
                ***** 1 42 * 6.07 * 18 * 00 * 77 * 51 *
             1.97 6 3 14.77.20.53.41.59 1.98 6.17.15.31.24.12.41.22 8/73-22 1.67 6.13.14.98.25.89.41.42 1.89 6.11.15.25.24.52.42.07 8/75-25
                .58: 1.60: 6.28:14.40:64.50: .58: 1.58: 5.98:25.92:65:11 H/75-24
               \(\T+459.61/529.6) ++.41+(1.+.38+RH)
     XXX=EF.
     1F (XXX.GT.L(1.K.N))GO TO 60
1F (XXX.LT.L(5.K.N))GO TO 30
     00 20 1=2.5
     IF(XXX.LT.E(1/K.N))GO TO 20
SCFMEEAP(ALOG(5(1-1/K.N))+(ALOG(XXX)-ALOG(E(1-1/K.N)))+(ALOG(5(
    11 + K + N 1 ) - ALUG(S(I-1 + K + N ) ) / (ALOG(L(1 + K + N ) ) - ALOG(L(I-1 + K + N ) ) ) )
     60 10 500
     CULTINUE
 OU SCHMEERPIALOGIS(1+K+N))+(ALOGIXXX)-ALOGIE(1+K+N)))+(ALOGISI
    .2.K.N1)-ALUG(5(1.K.N1))/(ALUG(E(2.K.N1)-ALUG(E(1.K.N1)))
     60 10 500
 30 5CF H=LXP(ALOG(5(4+K+N))+(ALOG(XXX)=ALOG(E(4+K+N)))+(ALOG(5(
15.K.N) -ALOG(5(4.K.N)))/(ALOG(E(5.K.N))-ALOG(E(4.K.N))))
50. ACFM=5CFM+.075/HHO
     PLFL0=5CFM+.075760+ ....
     METERST
     Liai
     SUBROUTING EMP FOR ST DATA REDUCTION PROGRAM SUBROUTING EMPT(LMF)TI(*)
DIMENSION E(40)+T(40)
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DATA £/0.50+0.60+0.70+0.80+0.85+0.90+0.95+1.00+1.05+1.10+1.19+1.20
                                                              +1.2:+1.30+1.35+1.40+1.45+1.50+1.55+1.60+1.65+1.40+1.76+1.76+1.80
                                                              +1.85+1.90+1.95+2.00+2.05+2.10+2.15+2.20+2.25+2.30+2.55+2.40
                                                              12.5012.6012.7012.80/
                           110.64+112-34+114.06+115-76+119-18-122-59+125-9/+129-5//
                            IF ( (EMF-61-0-5) - ANU- (EMF-LL-2-6) IGU TO 1
                            micate (meatil) EMP
            IN FORMATOTICE OUT OF EMP TABLES IN SUBROUTING EMPT. EMPT. F10.3)
                            RETURN I
                  1 IF (EMF.GT.1.0) GO TO 2
                           141 = 1
                            142-11
                             60 10 5
                 2 IF (EMF. 61.1.5) GO TO 3
                            14.2.44
                            112-18
                             60 10 5
                 3 IF (LMF. 61. _. U) GO TO 4
                            1.1 - 18
                             NELZH
                             60 TO 5
                 4 141-28
                            Ne -40
                 5 UU o I=Nleid
                              THE LEWIS LIVE LATER TO TO T
                 o CONTINUE
                           [1-1(1-1)+(1(1)-1(1-1))+(ENF-E(1-1))/(E(1)-E(1-1))
                              RETURN.
                            JUDROUTINE FOR ST DATA REDUCTION PROGRAM
SUDROUTINE HEADS(10UT*XZERU**)
COMMON/HDY(AMDT(4)*VAIR*TINE*TZEPO*HHOINE*PSI
COMMON/AMB/PAMB*TNB*TWB*DENT(18)*KI*IEHBAL*VAR
                               IF (K1-ML-6)GO TO 1
                              WHITE (0+10)
                TO FORMATO///*THE OUTPUT ROUTING REQUESTED CAN BE USED FOR STANTON RU
                         INS OR ENERGY BAL HERS UNLY ! I TOUT # 0
                                HETURN 1
                      1 IF (KI-GT-5)GO TO 3
WRITE(G-8)DENT-VAIR-TINF-12: RO-RHOIN-P51-TDR-TWB-PAMB-XZERO
                                  IF (1001.61.1)PUNCH 5.DEM. VAIR. TIME. TZERO, RHOTHE . PSI. TUB. TWB. PAMB
                   1.4.LRO

B FORMAT(IHI:18A4//* VEL=*,F5.1;* FT/SEC 1=*,F5.1;* DEGF TO=*,F5.1;

1.** OEGF RIGO=*,F6.4;* LB5/F13 P=*,F6.2;* P51*/* TDR=*,F5.1;* DEGF 2.* TWB=*,F5.1;* DEGF PAMH=*,F6.2;* IN HG VIR ORIG;*XO=*,F6.2;* IN*

3//* PL X 5TN MOM ENTH X (X-XO) F B TPL
4 TAIR MOMTH ENTTH*/* NO INCH NO REYNO REYNO REYNO REYNO REYNO*, F6.2;* IN*
5x;* ULGF DLGF INCH INCH*)
5 FORMAT(IH:18A4//* VEL=*,F5.1;* FT/SEC T=*,F5.1;* DEGF TO=*,F5.1;* ULGF RIGO=*,F6.4;* LB5/FT3 P=*,F6.2;* P51*/* TDR=*,F5.1;* ULGF RIGO=*,F6.4;* LB5/FT3 P=*,F6.2;* P51*/* TDR=*,F5.1;* ULGF RIGO=*,F5.1;* ULGF RIGO=*,F5.1;* ULGF RIGO=*,F5.1;* ULGF RIGO=*,F6.4;* LB5/FT3 P=*,F6.2;* P51*/* TDR=*,F5.1;* ULGF RIGO=*,F5.1;* ULGF RIGO=*,F6.2;* IN*
4 TAIR MOMTH FNTH*/* NO INCH NO REYNO REYNO REYNO REYNO*,F6.2;* TGRESSORE R
                           1 . ALLRO
                            5x++0EGF DEGF INCH INCH+)
GO TO 4
                       3 CONTINUL
                                  1LLHAL#1
                                 1+ (K1.L0.4)K1=1
1+ (K1.L0.5)K1=2
                                   VALR=VAH
                       WRITE CO.71DENT.TDB.TWB.PAMB.VAIR
7 FORMAT(1HI.1844//* TDB=**F5.1** DEGF TWB=**F5.1** DEGF PAMB=**F6
1.2** IN HG VEL(5IN ERR)=**F6.1** FT/SEC*//* PL PL POWR E CONV
2 LOUSES DIFF PERCHT 5TN FLOW TPL TAIR PL*/* NO UT
3U/SEC BTU/SEC BTU/SEC BTU/SEC**IDX**ERROR ACFM DEGF DEGF
4 NO**
1F (TOBLEGT TERRORS DEGENTATION SWEETER ACFM)
1F (TOBLEGT TERRORS DEGENTATION SWEETER ACFM)
                        # NO*/
IF (IDUT.GT.1)PUNCH 9:DENT:IDB:TWB:PAMB:VAIR
9 FURMAT(IH :IBA47/* IDB:**F5:1:* DEGF IMB:**F5:1:* DEGF PAMB=*:F6
1.2:* IN He VEL(5TH LRR)=**F6:1:* FY/SEC*//* PL PL POWR E CONV
2 LOSSES DIF! PERCHT SIN FLOW THE TAIR PL*/* NO DT
SU/SEC BIU/SEC BIU/SEC BIU/SEC*/10%:*ERROR ACFM DEGF DEGF
# NO*/
# DEGF DEGF
                         4 HL TURN
                                  LINU
                                THE THE PARTY SOME STATES OF THE PROPERTY OF T
L DATADI
                           COMMON DATA AND SOME STATEOUR NO MATEOUR AND PROTECT OF THE PROTEC
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2510(24) - REXO(24) + ITAG + TAMB(4) + 10UT + IFR + IRMUR
      COMMON/DAO/ (EX (24) . STN (24) . BFE (24)
     COMMON/AMB/PAMB. TOB. THE DELLT(18) AK1 - TENBAL - VAR
      IF ( (IENBAL.EQ.0) . OR. (K1.EQ.3) ) 40 40
IF ((IENBAL, EG.O) - OR (RI, EG.S) / O TO 40
ANTTE(6, 17) THE 10B - PAMB - POBLY TAIR (TAMB) () - I = 1 + 4)

17 FUR MAT(IM ** RAW IMPUT DATA FOR ENERGY DALANCE RULL**

1* TWB=**F6.1** DEGF TUB=**F6.1** DEGF PAMBE**, PB. 2*** AND PDCAE**,

2FU.2***HACO TAIR=**F6.3** MY'Y TAMBE***, PB. 2*** MY TAMBE***, PB. 2***

3MY TAMB3=**F5.3** MY VAMBUE***FB. 2*** MY'I

1*(KI = EG.3) ARITE(6.8)

A FOLMAT(** PI AND TAIMN) TAMBE** FMC/JOLYS** EC(6V) CAST
                                           TILIMUS TANIMUS EMITOLYSI
  B FURMAT( PL NO
                                                                                                      ECHAVI CAST
  I TC TCS(MV) *)

IF(KI+NE-3)#RITE(6-9)

9 FURMAT(* PL NO TPL(MV) TAR(MV
FC(MV) CAST TC TCS(MV) *)
                                                              TAR(MUI TERMINA) TELUIMUI LILLVOL
    17) EC (MV) CAST
100 2 N=1:24
1F(14.6T.12) GO TO 3
     IF (KI . EQ. 3) ARITE (6.32) H. TPL (N) . TAR (N) . EH (N) . ES (N) . (1, 765 (N) . IF (KI . NE. 3) ARITE (6.15) (A TPL (N) . TAR (N) . TCFM(N) . CFM(N) . EI (N) . ES (N) . (1
    LITCS (N)
      GU TO 2
  3 IF (K1.EQ.3) ARTTE (6.39) NOTPL (N) OTAR (N) ORHITO (ES (N)
      IF (K1.NE.3) WRITE (6.40) N. TPL(N) . TAR(P) . TCFM(N) . CFM(N) . CFM(N) . CSA(N)
  2 CUNTINUL
      HETURN
50 IF (K1-NE.3) WRITE(6.51) TWB. TCB. PAME . TANE(1) . TAME(2) . TAME(3) . TAME(4
1 /' TATR='+F6.3', MV VARE-+F7.3', **H20 P5:AT=**F6.2+*L2H20
2PUDX=*F6.2+*2H20 SPGR=*F6.3+ XZED0=*+F6.2+*L2+*
53 If (K1.EQ.3) WHITE (6.33)
33 FORMAT(* PL NO TPL(MV) TAR(MV) LH(VOLTS) LC(MV) CASI
1 IC TCS(MV) STS ST/STS *)
1F(K1.NE.3) WRITE (6.34)
34 FORMAT(* PL NO TPL(MV) TAR(MV) ICFM(MV) IFLO(MV) LH(VOL
1T) EC(MV) CAST TC TCS(MV) STS ST/STS ST/STO (N
     1(1+11)/81)
      UU 13 N#1.24
      1 (STR(N).GT.O.)STRAO=STN(N)/STR(N)
      IF (STR(N).EQ.O.)STRAC=U.
IF (REX(N).LT.REXO(1))GO TO 70
       IF ((NST.EG.0).OR. (K1.EG.3))GO 10 58
      DU 72 I=2, ITAB
IF (REX(N).LT. REXO(I)) GU TO 71
 72 CONTINUE
       STOO=STO(ITAB)
       60 TO 73
 71 RATOO=(REX(N)-REXO(I-1))/(REXO(I)-RFXO(I-1))
$[00=$TO(I-1)+($TO(I)-$TO(I-1))*RATOO
       60 TO 73
 70 $100=$TO(1)
73 $15TO=$TN(N)/$TOO
16 (BEE(N).GT.O.)FLNB=ALOG(BEE('.)+1.)/BEE(N)
       IF (ME(N).LE.O.)FLMB=0.0
IF (N.GT.12)GO TO B9
IF (KI.EQ.3)#RITE (A.*S.N. (PL(N).TAR(N).EH(N).ES(M).N.TCS(M).STR(M).
     151HA0
      15|RAU

FORMAT(110 2F10.3,F10.4,F10.3,110,F10.3,F10.5,5F10.2)

1F(K1.NF,3) WRITE(6,1B) H. TPL(N), TAR(N), TCFM(N), CHM(N), LH(N), E5(N), N

1:TCS(N',5TK(N),STRAO,STSTO,FLNB
       FORMA' (110.4F10.3.F10.4.F10.3.110.F10.3.F10.5.3F10.5)
       66 TO 13
 59 1F(K..E0.3) WRITE(6.39)N.TPL(N).TAR(N).EH(N).E5(N).5TR(N).5TRAO
39 FOR AT((10.2F10.3,F10.4,F10.3,20X,F10.5,3F10.3)
1F(K1.HE.3)WRITE(6.MO)N.TPL(N).TAR(N).TCFM(N).CFM(N).E5(N).E5(N).S
      171-(N) . STRAD, STSTO, FLNB
        "JRMAT(110,4F10.3,F10.4,F10.3,20X,F10.5,3F10.3)
  SU TO 13
       IF (N.GT.12)GO TO 38
IF (N.EG.3) #RITE(6,32)N: TPL(N): TAR(N): EH(N): ES(N): N: TCS(N): STR(N):
      ISTHAO
       IF (K1. NE. 3) WRITE (6, 15) N. TPL(N) . TAR(N) . TCFM(N) . CFM(N) . EII(N) . E5 (N) . N
      LITCH (N) STRIN) STRAD
  GG TO 13
38 IF (K1.80.3) WRITE (6.39) H. TPL (H) . TAR (N) . EH (N) . ES (H) . STR (H) . STR AD
IF (K1.NE.3) WRITE (6.40) H. TPL (N) . TAR (N) . TCPM (N) . CFM (H) . LICH) . ES (H) . G
      ATRILLIA STRAU
  13 COLLY INUE
       HLTUHN
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